

Cosmic Ray Detection

• Charles Timmermans



Outline

- Energy loss by a particle in a medium
- Creation of detectable radiation
- Air Showers
- The Pierre Auger Observatory
- Hands on experimentation

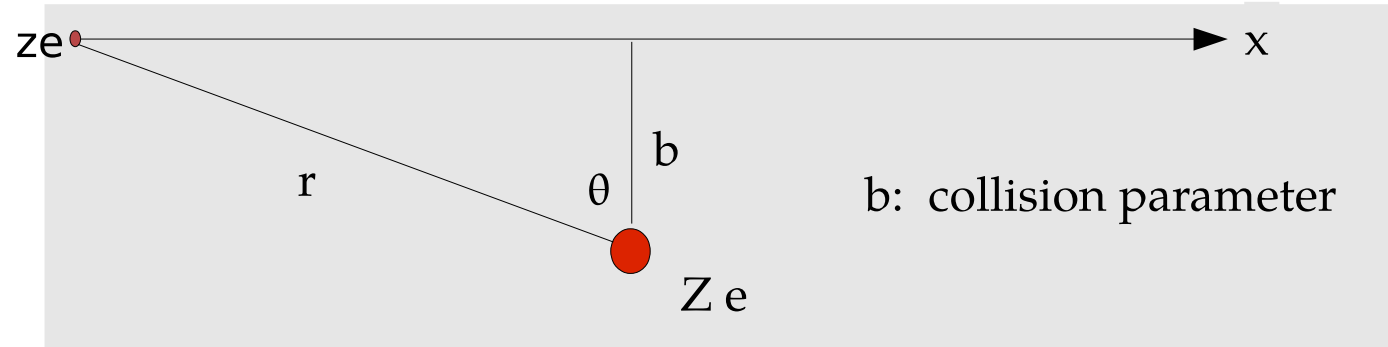
Basics

- All detection of particles rely on interaction of particles with an active medium, thereby losing energy and producing a measurable signal
- That measurable signal can be
 - Light
 - Radio waves
 - Sound
 - Heat
 - Bubble density

Energy loss: Bethe-Bloch

- Ionisation and excitation of the surrounding medium:
 - Air shower detection
 - Signal from water Cherenkov detectors
 - Signal from scintillator detectors
 - Tracking detectors
 - Calorimeters

Classical Energy-loss formula



b: collision parameter

Momentum of target particle changes when test particle travels through medium:

$$\Delta \vec{p} = \int_{-\infty}^{+\infty} \vec{F}_{\text{Coulomb}} dt$$

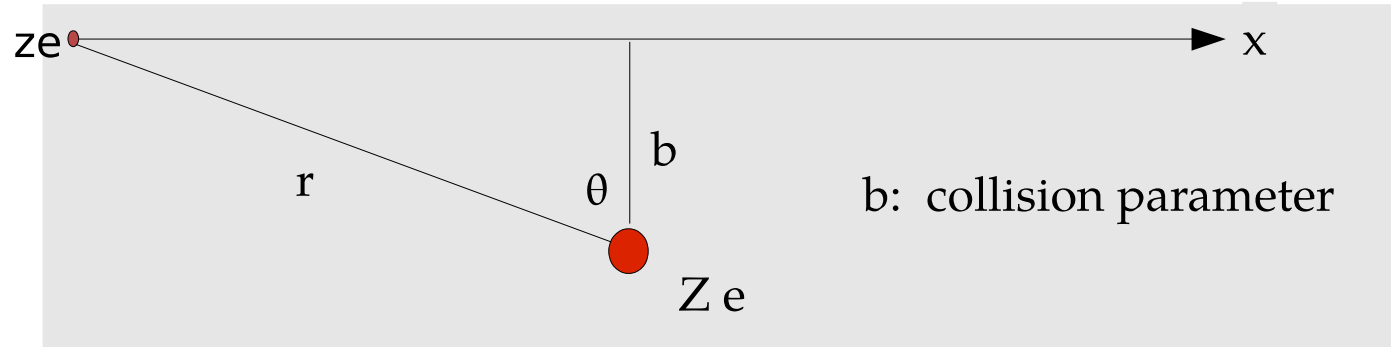
$$\vec{F}_{\text{Coulomb}} = \frac{1}{4\pi\epsilon_0} \frac{zZe^2 \vec{r}}{r^3}$$

Integrating from $-\infty$ to $+\infty$, only transverse force is important:

$$F_T = F_{\text{Coulomb}} \cdot \frac{b}{r(t)} = F_{\text{Coulomb}} \cdot \cos \theta(t)$$

$$F_T = \frac{1}{4\pi\epsilon_0} \frac{zZe^2}{b^2} \cos^3 \theta$$

Classical Energy-loss formula



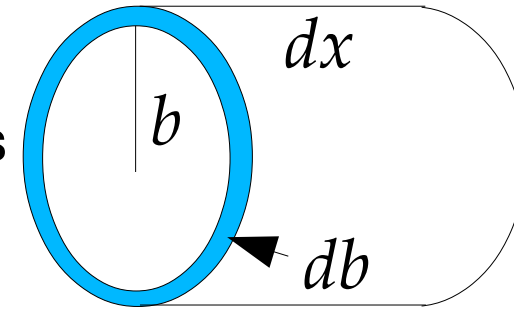
Therefore the energy transfer onto a target with mass m_t is given by

$$\Delta E(b) = \frac{\Delta p(b)^2}{2m_t} = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{2z^2 Z^2 e^4}{m_t c^2 \beta^2 b^2}$$

Only interactions with electrons are important!

Classical Energy-loss formula

Let the test particle move through a tube of electrons



Number of electrons: $N_e = 2\pi b \cdot db \cdot dx \cdot n_e$

Energy loss due to tube:

Electron density: $n_e = N_A \cdot (Z/A) \cdot \rho$

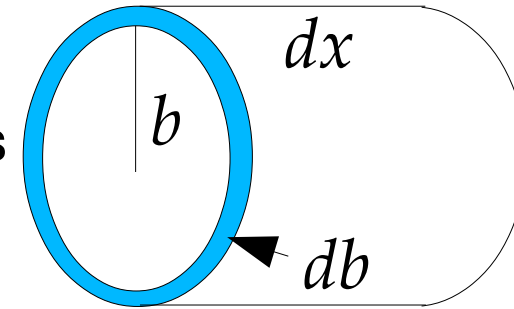
$$\begin{aligned} -dE(b) &= \Delta E(b) n_e dV \\ &= \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{4\pi z^2 e^4}{m_e c^2 \beta^2} n_e \frac{db}{b} dx \end{aligned}$$

Integrating over b:

$$-\frac{dE}{dx} = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{4\pi z^2 e^4}{m_e c^2 \beta^2} N_A \cdot (Z/A) \cdot \rho \ln \frac{b_{max}}{b_{min}}$$

Classical Energy-loss formula

Let the test particle move through a tube of electrons



$$-\frac{dE}{dx} = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{4\pi z^2 e^4}{m_e c^2 \beta^2} N_A \cdot (Z/A) \cdot \rho \ln \frac{b_{max}}{b_{min}}$$

Maximum energy transfer (central collision): $\Delta E_{max} = 2\gamma^2 m_e \beta^2 c^2$

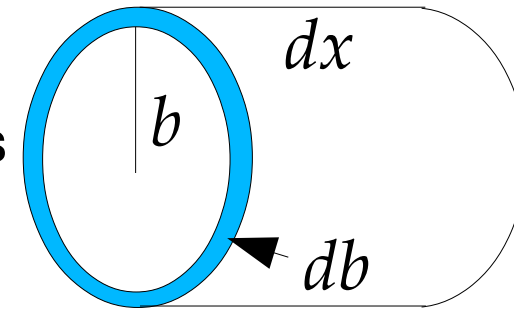
Minimal b:
$$b_{min} = \left(\frac{1}{4\pi\epsilon_0} \right) \frac{ze^2}{\gamma m_e c^2 \beta^2}$$

Minimal energy transfer: Mean excitation energy per electron /

Maximal b:
$$b_{max} = \left(\frac{1}{4\pi\epsilon_0} \right) \frac{ze^2}{c\beta} \sqrt{\frac{2}{m_e I}}$$

Classical Energy-loss formula

Let the test particle move through a tube of electrons



$$-\frac{dE}{dx} = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{4\pi z^2 e^4}{m_e c^2 \beta^2} \underbrace{N_A \cdot (Z/A) \cdot \rho}_{\text{Electron Density}} \cdot \frac{1}{2} \ln \frac{2c^2 \gamma^2 \beta^2 m_e}{\underbrace{I}_{\text{Average excitation potential}}}$$

$$W_{\max} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e / M + (m_e / M)^2}$$

$$\frac{\delta(\beta\gamma)}{2} \rightarrow \ln \left(\frac{\hbar \omega_{\text{plasma}}}{I} \right) + \ln(\beta\gamma) - \frac{1}{2}$$

Bethe-Bloch:

$$\left\langle -\frac{dE}{dx} \right\rangle = \rho K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} \right) - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

m_e mass of the electron,

c speed of light,

$K = 4\pi N_A r_e^2 m_e c^2$ numerical coefficient with value $K = 0.307075 \text{ MeV mol}^{-1} \text{ cm}^2$ (N_A is Avogadro's number and r_e is the classical radius of the electron),

z charge of the high energy particle traversing the matter,

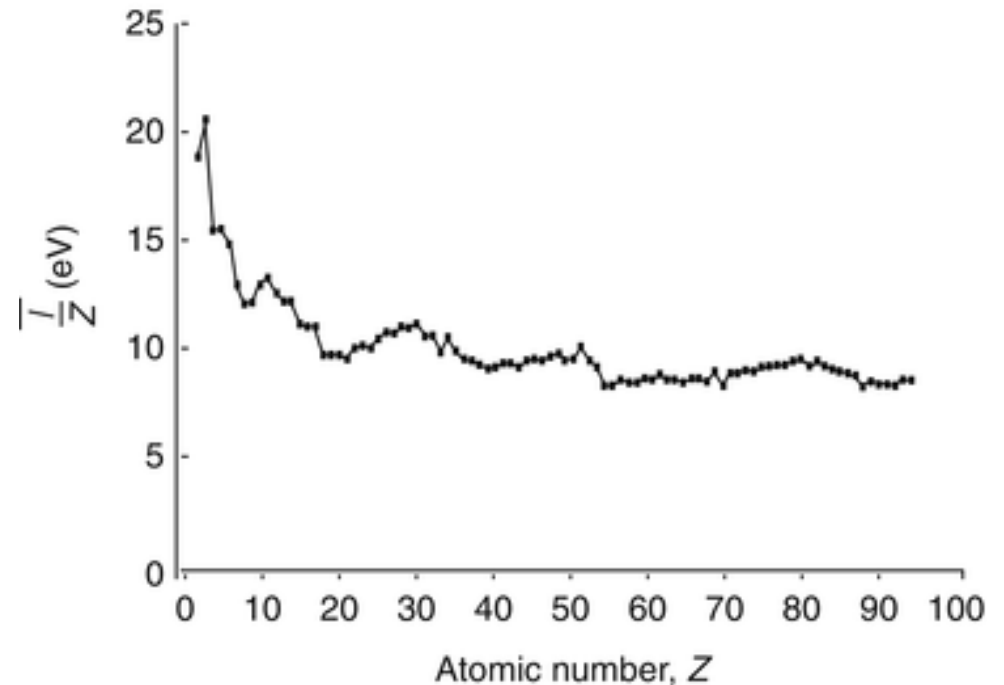
β, γ velocity as fraction of the light velocity and Lorentz contraction factor of the high energy particle traversing the matter,

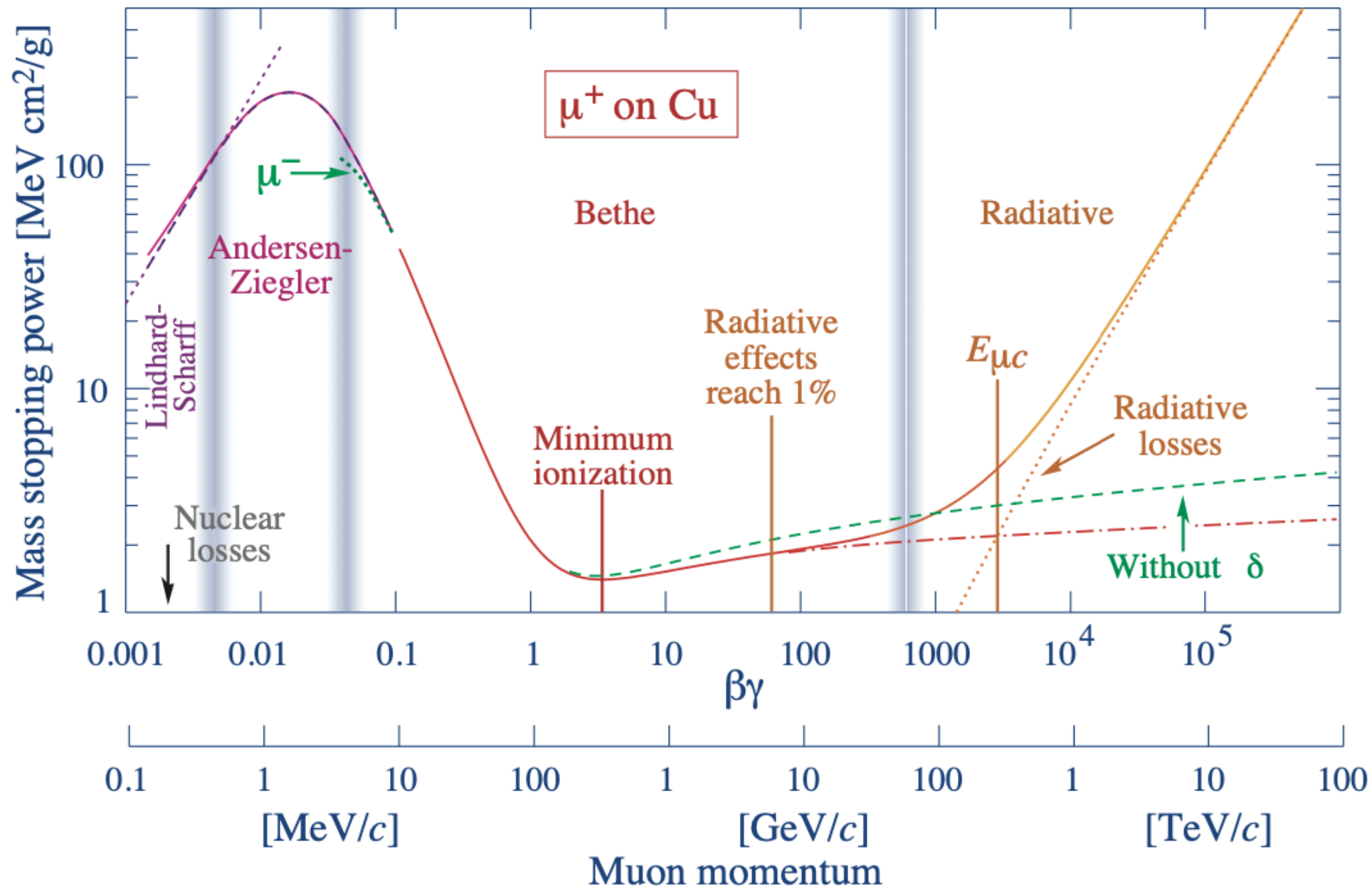
ρ the mass density of the matter,

Z, A charge and mass number of the atom the matter consists of,

I mean excitation energy for the electrons in the atom,

$\delta(\beta\gamma)$ correction for the polarisation of the matter for very-high-energy particles.





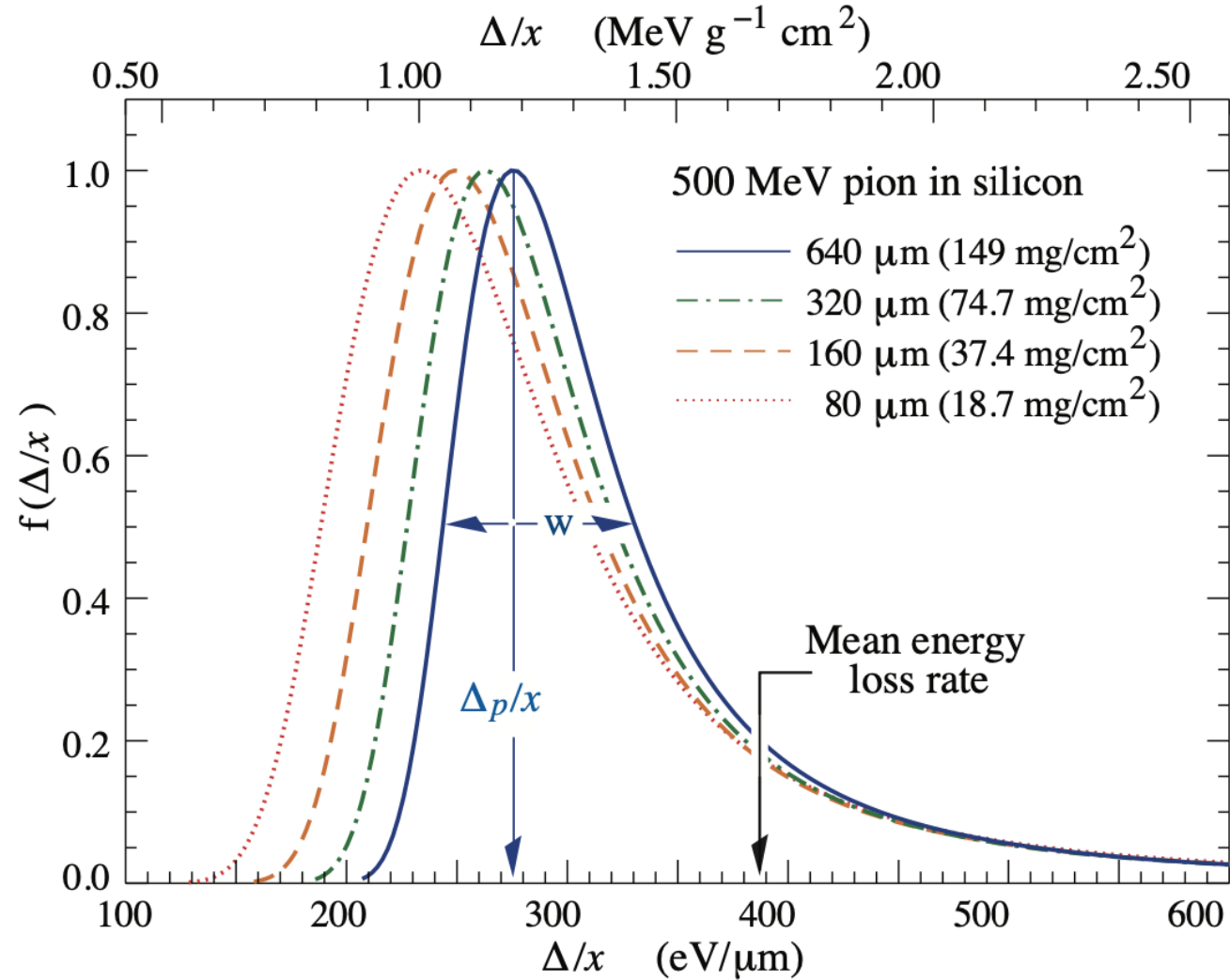
Fluctuations in Energy loss

- Bethe-Bloch describes the average energy loss, not the most probable!
- Landau gave a description of the fluctuations, assuming only the EM-interactions with the electrons in the medium:

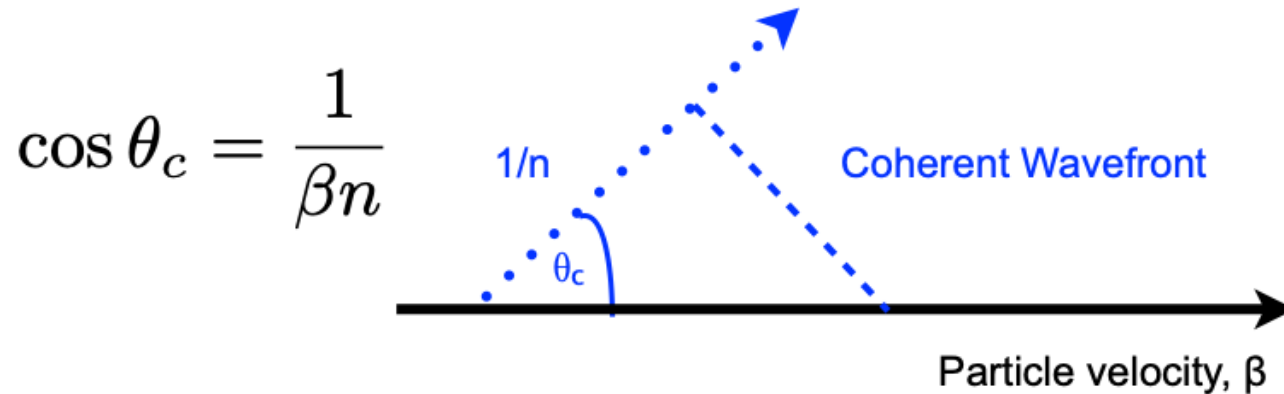
$$p((\Delta E)_{\text{mp}}; \beta\gamma, \Delta x) = \frac{1}{\pi \Delta x} \int_{t=0}^{\infty} e^{-t} \cos \left[t \frac{(\Delta E)_{\text{mp}} - \beta\gamma}{\Delta x} + \frac{2t}{\pi} \ln \left(\frac{t}{\Delta x} \right) \right] dt$$

$$\begin{array}{ll} \Delta x & \text{Thickness of medium} \\ r & \text{density of medium} \end{array} \quad (\Delta E)_{\text{mp}} = \xi \left[\ln \left(\frac{2m_e c^2 \beta^2 \gamma^2}{I} \right) + \ln \left(\frac{\xi}{I} \right) + j - \beta^2 - \delta(\beta\gamma) \right] \quad \begin{array}{l} \xi = \rho \frac{K Z z^2}{2 A \beta^2} \Delta x \\ j = 0.200 \end{array}$$

Landau Fluctuations



Cherenkov Radiation



Discovered in 1934 by
Cherenkov and Vavilov

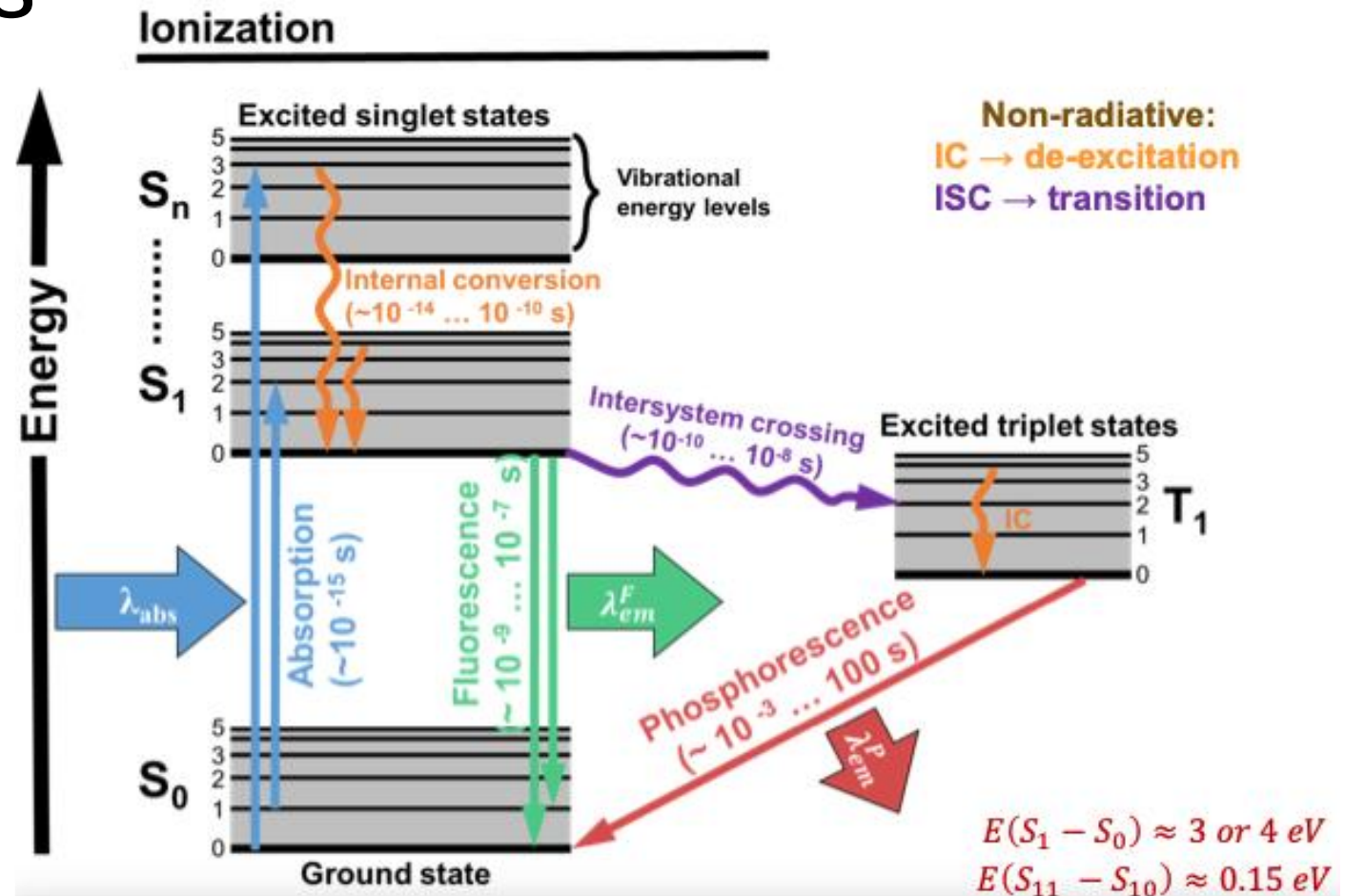
Described in 1937 by Frank and
Tamm

$$\frac{d^2 N}{dx d\lambda} = \frac{2\pi\alpha z^2}{\lambda^2} \left(1 - \frac{1}{\beta^2 n^2(\lambda)} \right) \quad N_{\text{p.e.}} = L \frac{\alpha^2 z^2}{r_e m_e c^2} \int \epsilon(E) \sin^2 \theta_c(E) dE$$

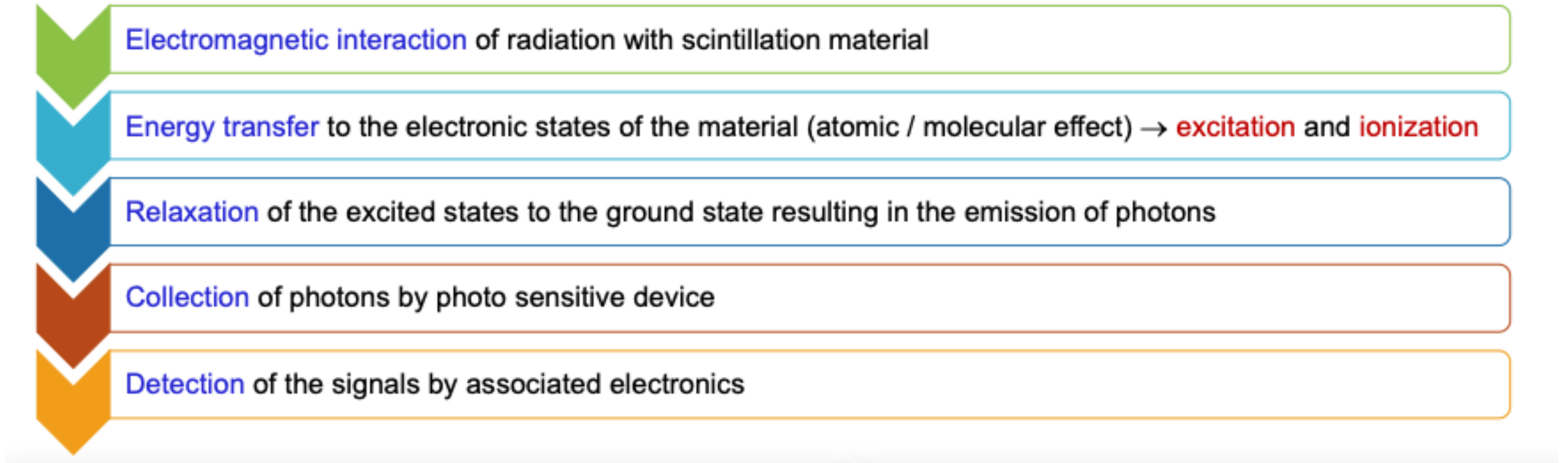
- Charged particles cause radiation if their velocity is greater than the local phase velocity of light
- This radiation is coherent at a particular angle relative to the direction of the particle
- In water energy loss is about 10^4 less than that of ionization!

Organic scintillators

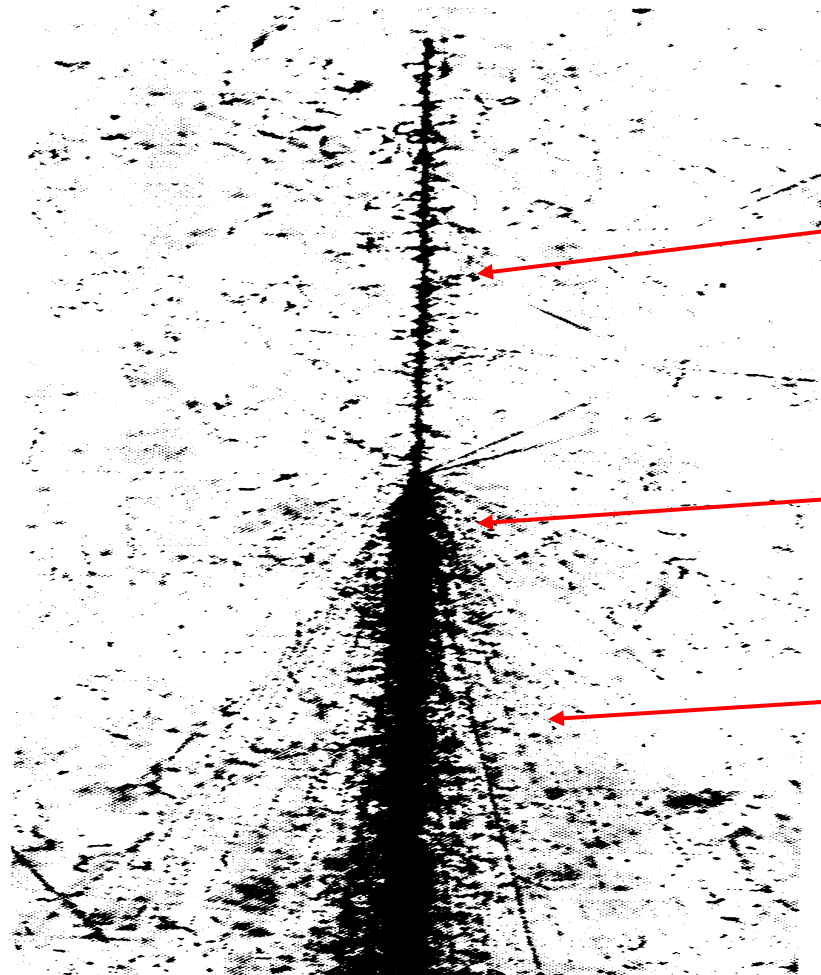
- Low-Z materials
- Fast response time
- Best for electrons and muons



Schematic of scintillator useage



Air Showers



Cosmic Ray

Collision in the Emulsion

Particle Cascade

- A cosmic ray induced shower in an emulsion



Energy Transfer

Hadronic
energy

$$\frac{2}{3}E_0$$

$$\frac{2}{3} \left(\frac{2}{3}E_0 \right)$$

⋮

$$E_{\text{had}} = \left(\frac{2}{3} \right)^n E_0$$

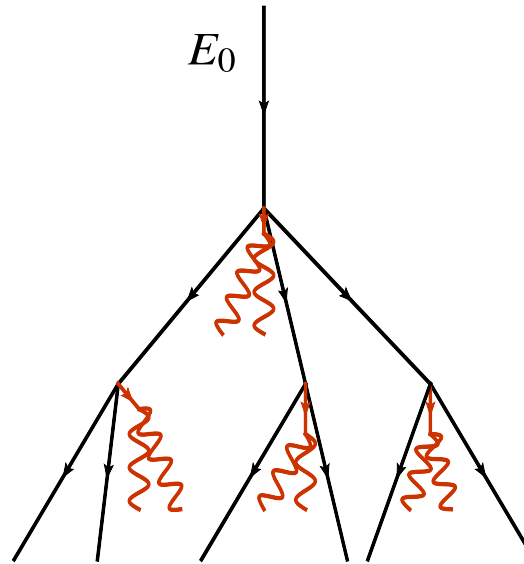
Electromagnetic
energy

$$\frac{1}{3}E_0$$

$$\frac{1}{3}E_0 + \frac{1}{3} \left(\frac{2}{3}E_0 \right)$$

⋮

$$E_{\text{em}} = \left[1 - \left(\frac{2}{3} \right)^n \right] E_0$$

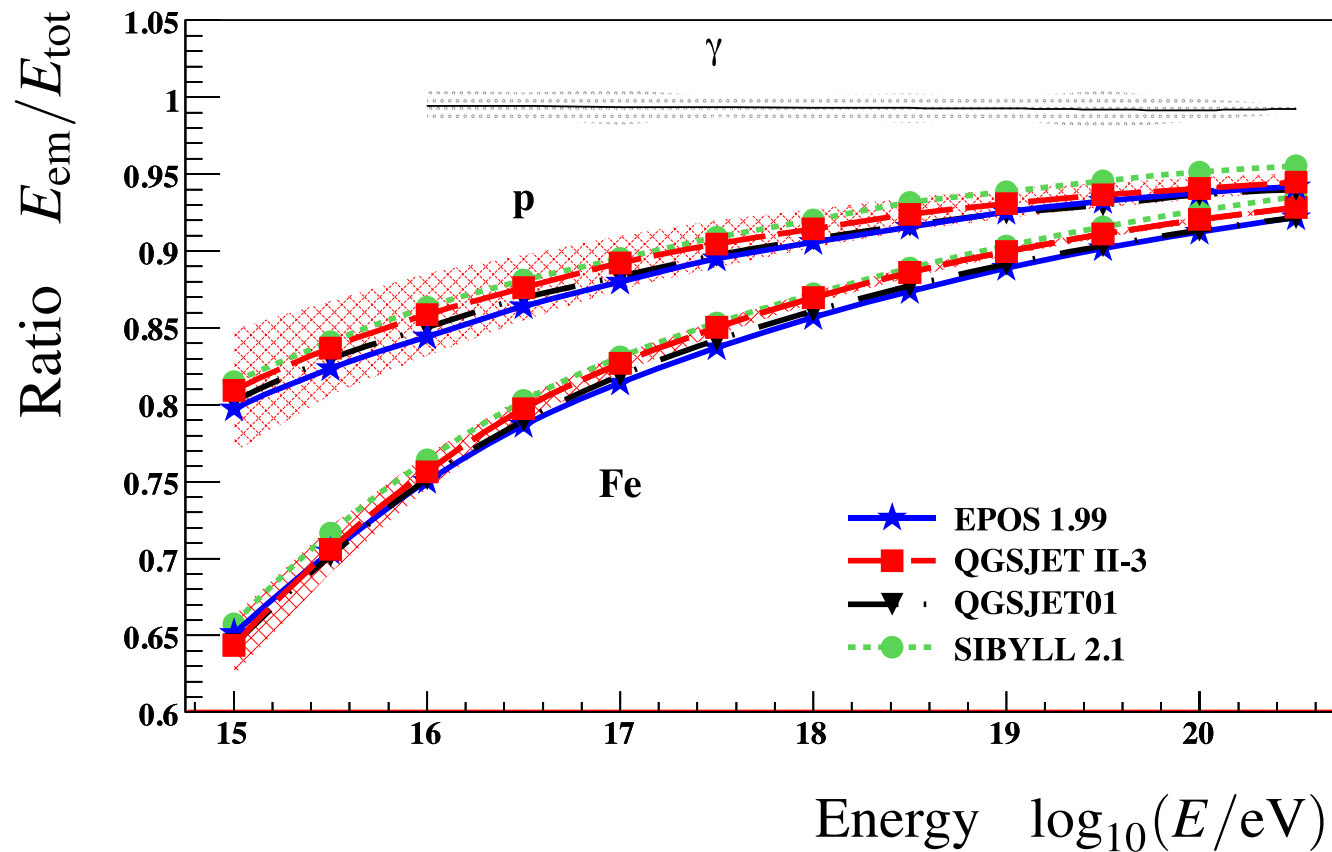


*After n
generations ...*

$$\begin{aligned} n = 5, & \ E_{\text{had}} \sim 12\% \\ n = 6, & \ E_{\text{had}} \sim 8\% \end{aligned}$$

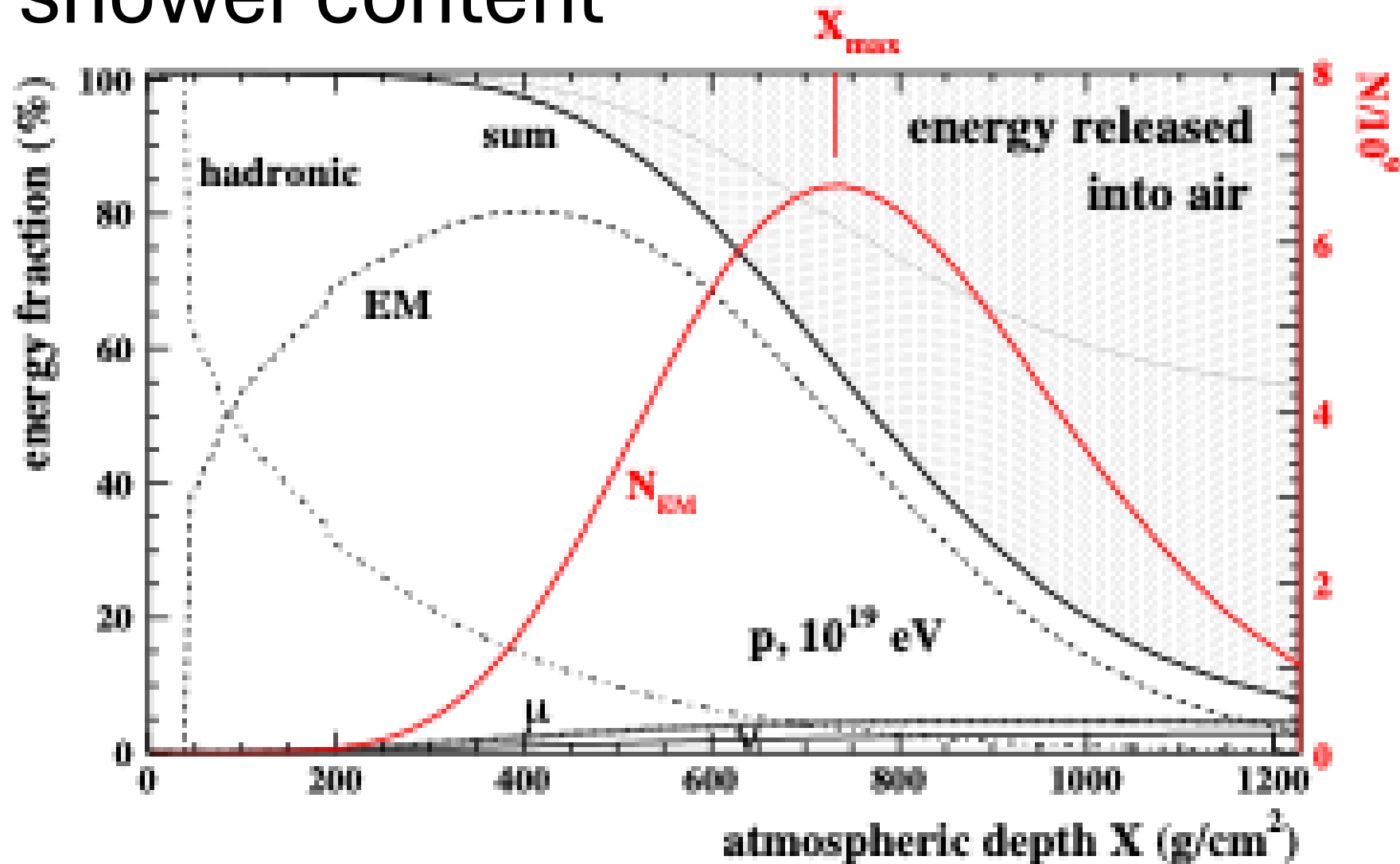
Energy Transfer

(RE, Pierog, Heck, ARNPS 2011)

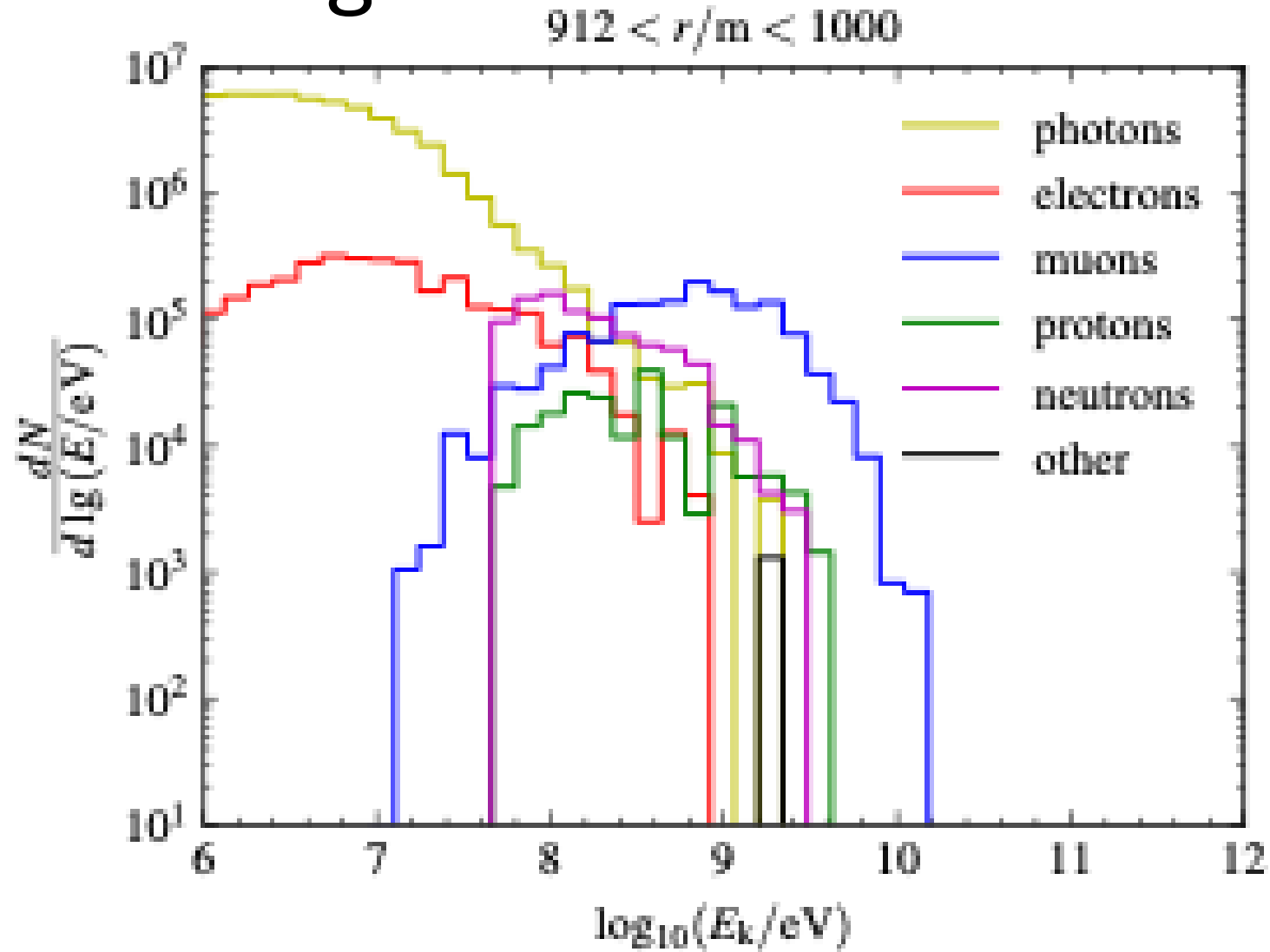


At low energy muons carry away a significant fraction of the energy

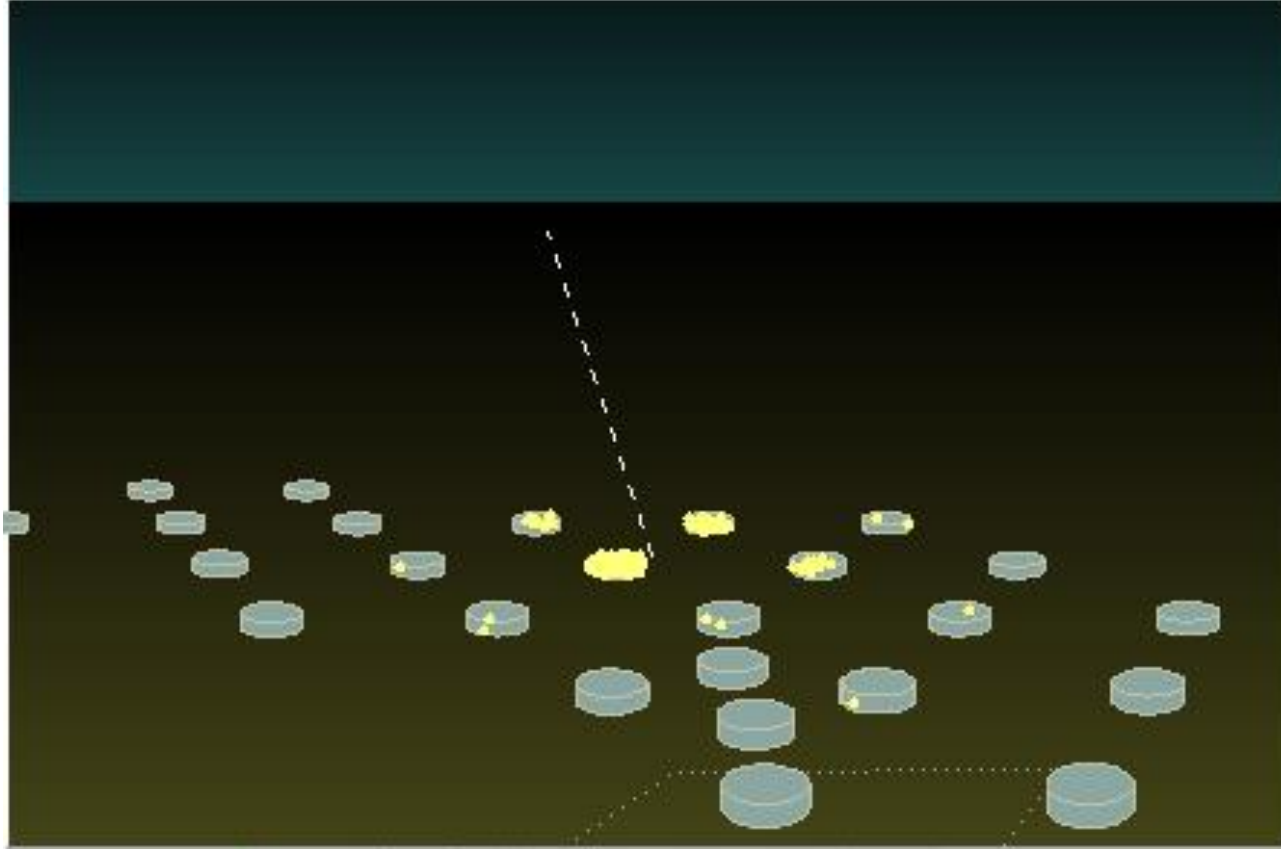
Air shower content



Particle Energies



Rough description of An Air shower



An electromagnetic current of changing size moving through the atmosphere at the speed of light

Creation of a radio signal

Macroscopic: Compare the air shower to a wire in which a changing current is induced that is moving perpendicular to its length with the speed of light

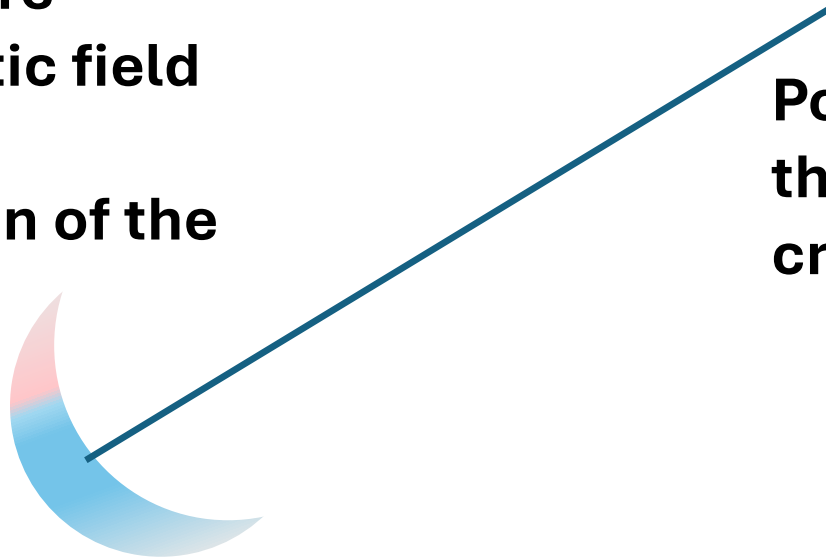
Microscopic: Trajectories of individual electrons and positrons are influenced by magnetic fields causing synchrotron radiation

Macroscopic picture of the Radio Signal

Electrons and positrons are deflected in Earth magnetic field

Net current in the direction of the Lorentz Force

Positive Ions travel slower than the air shower particles, but are created at the speed of the shower



Simulation Codes:

EVA

MGMR3D

(Krijn de Vries, Olaf Scholten)

Shower front has a net-negative charge due to:

Compton Scattering

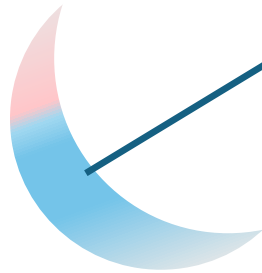
Knock out electrons

Positron annihilation

Microscopic picture of the Radio Signal

Follow individual particles

Emission due to acceleration of charged particles



Is added to a code for simulation of the creation of particle simulation of Air showers code

Simulation Codes:

COREAS

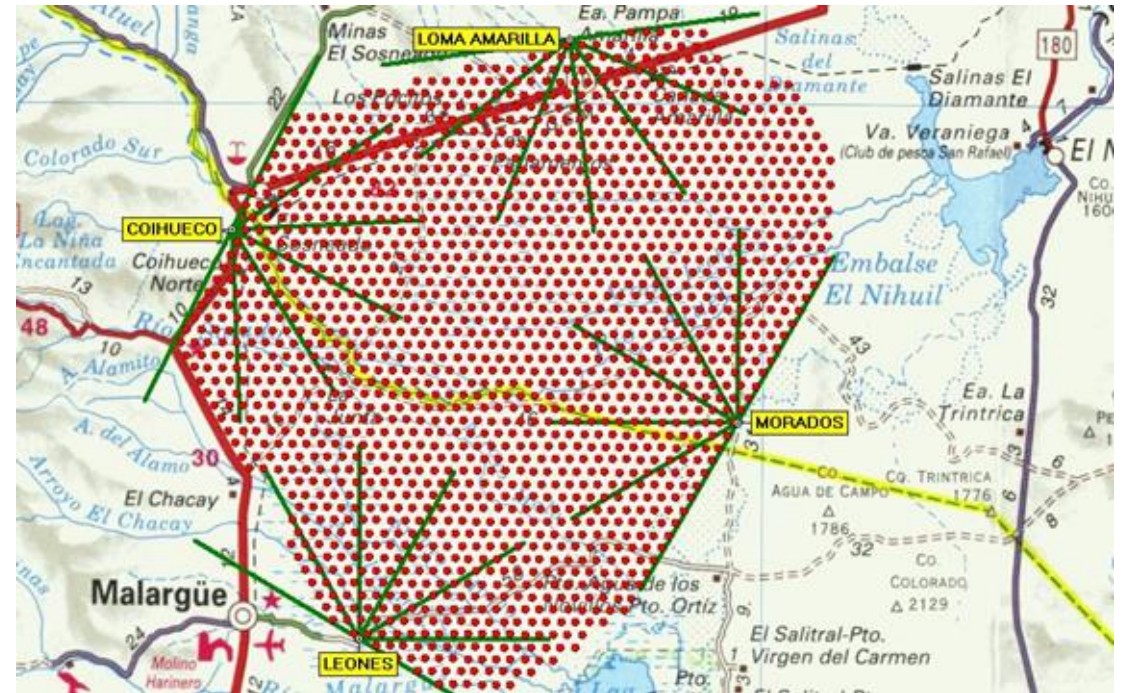
ZHAIRES

(Huege,

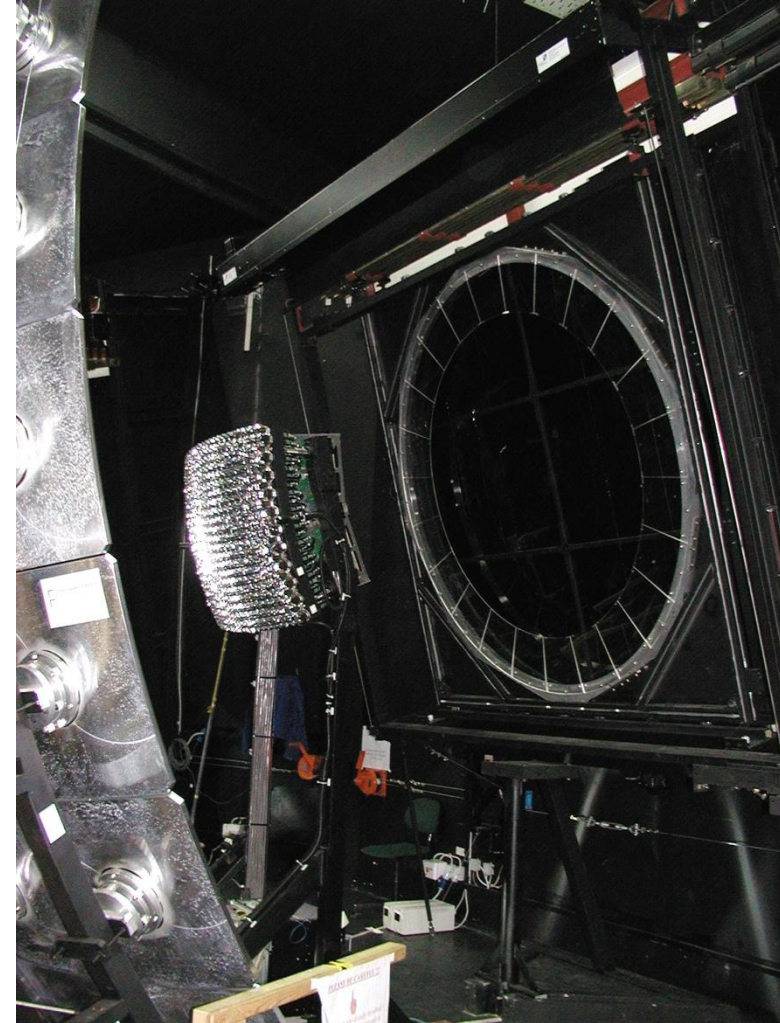
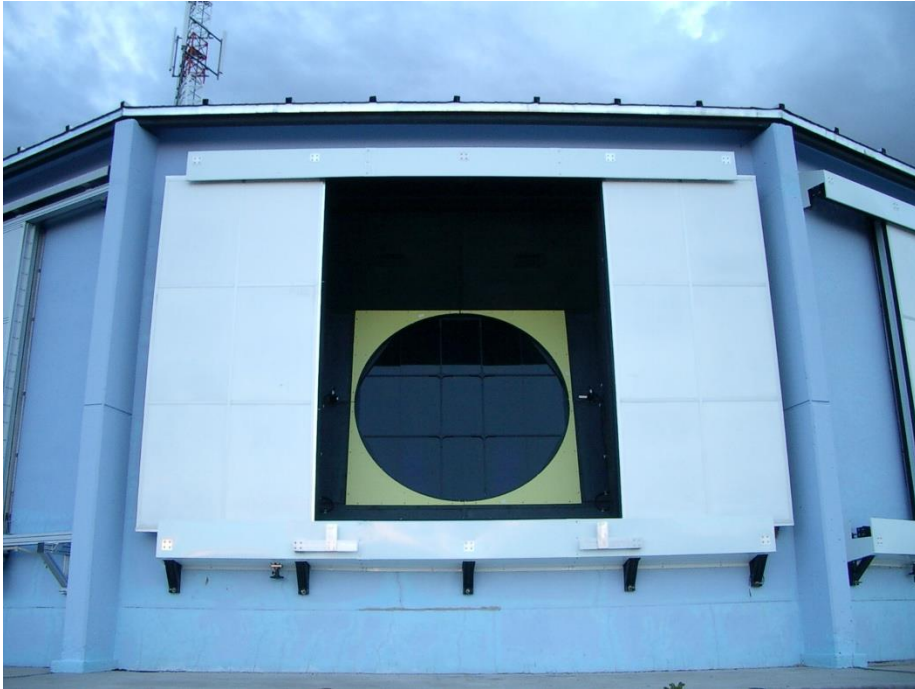
Zas, Halsen, Stanev)

Makes no assumption on different emission mechanisms

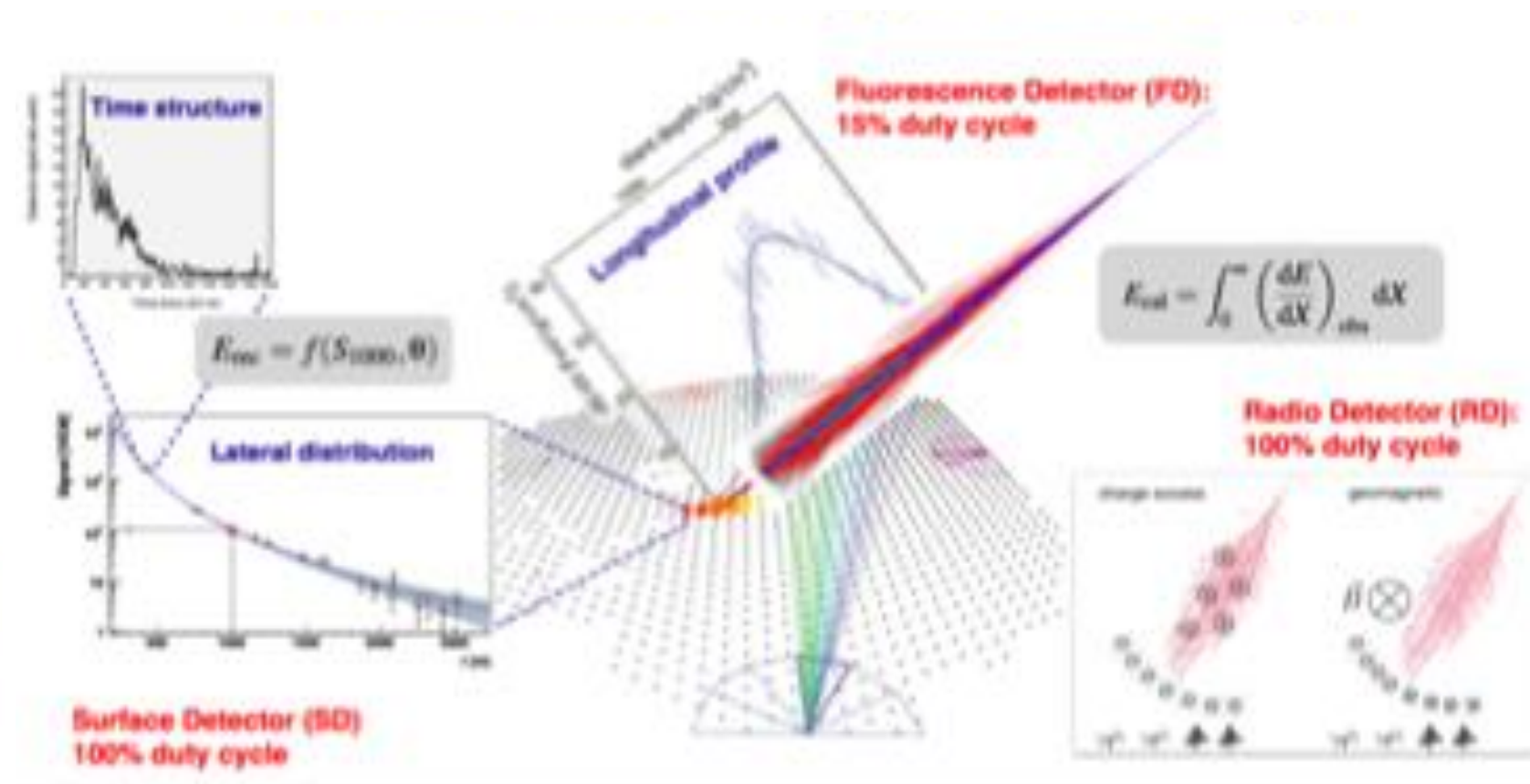
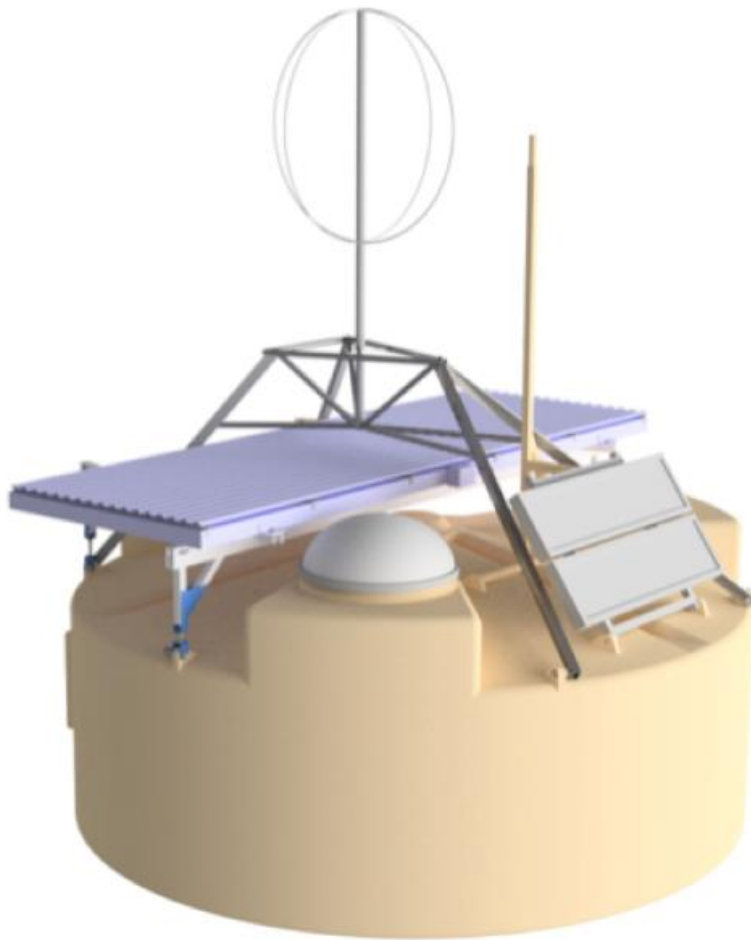
The Pierre Auger Observatory



The Pierre Auger Observatory



The Pierre Auger Observatory

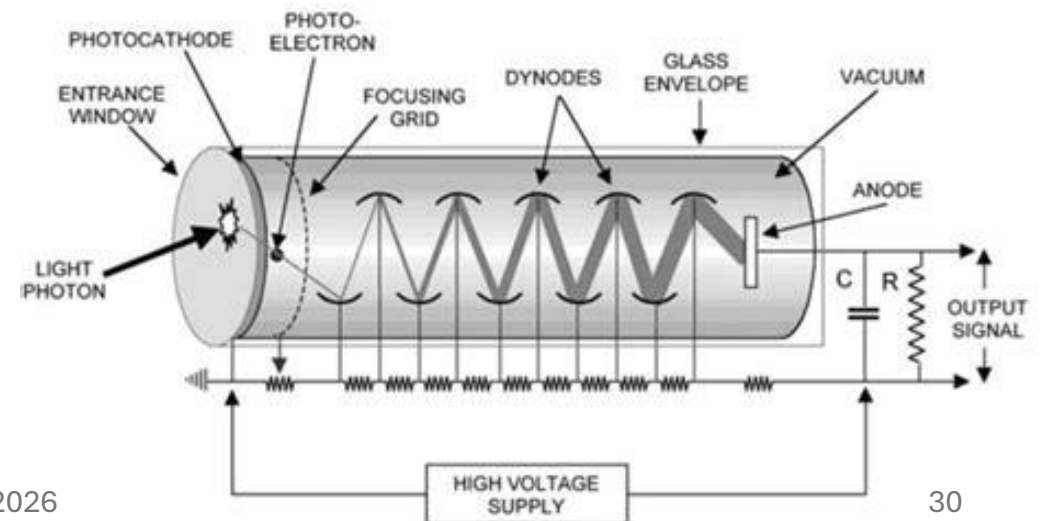
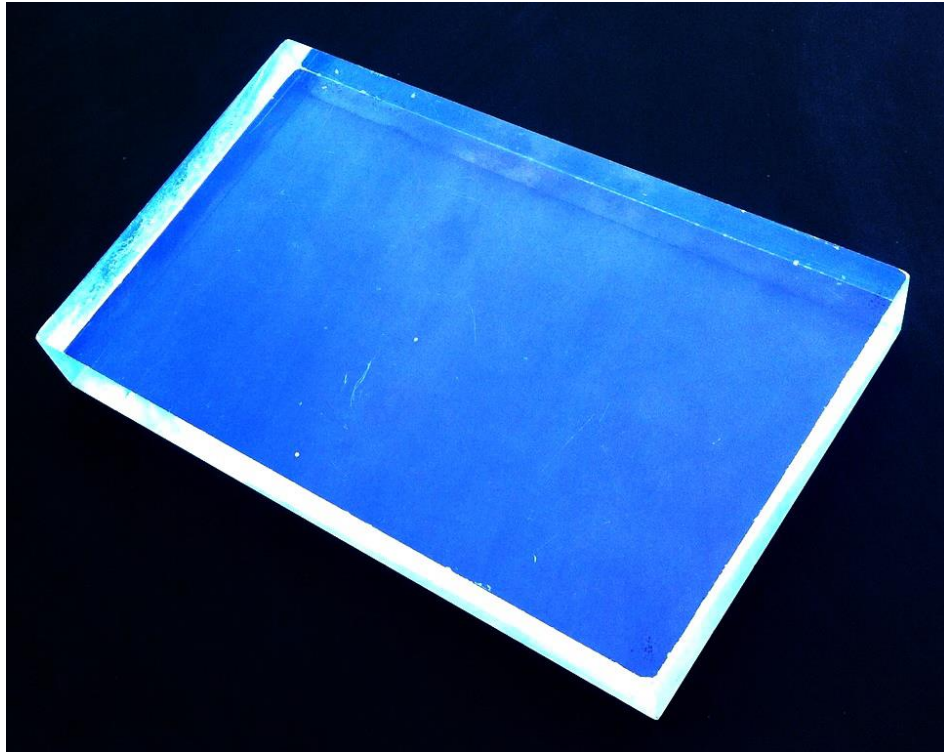


Hands-on: Determine the (vertical) muon flux

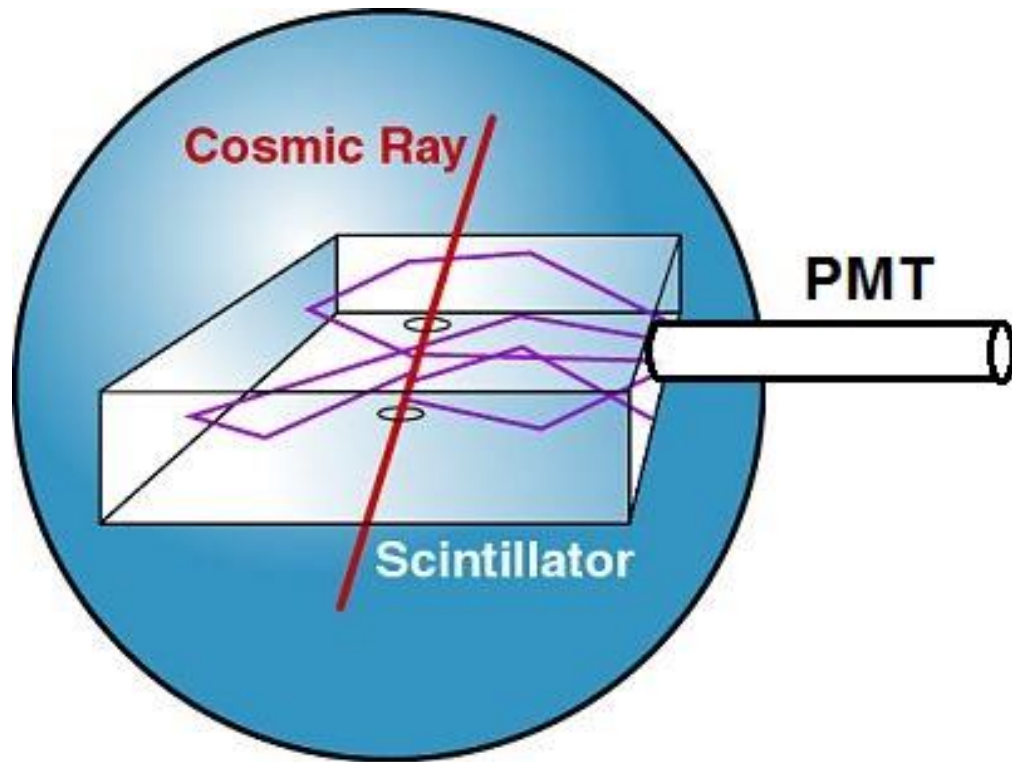
Parts:



Hands-on: Determine the (vertical) muon flux



Hands-on: Determine the (vertical) muon flux

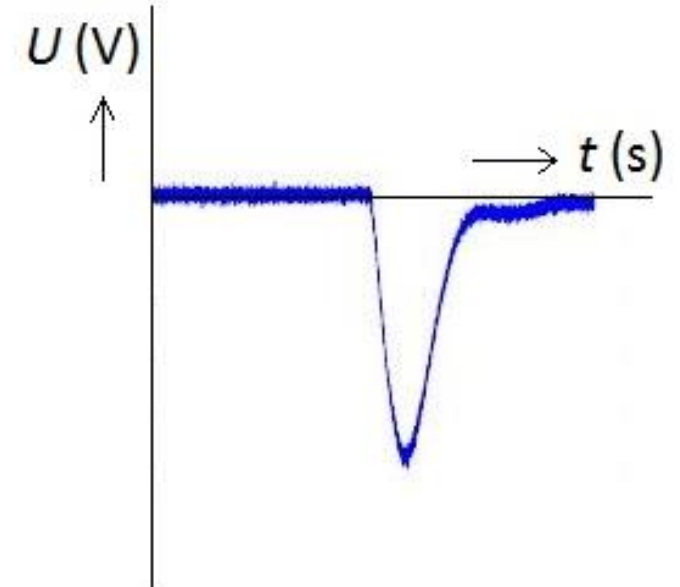


1-2% of energy loss is converted in fotons $\rightarrow 10^5$ fotons/muon

$\sim 1\%$ of fotons reach PMT

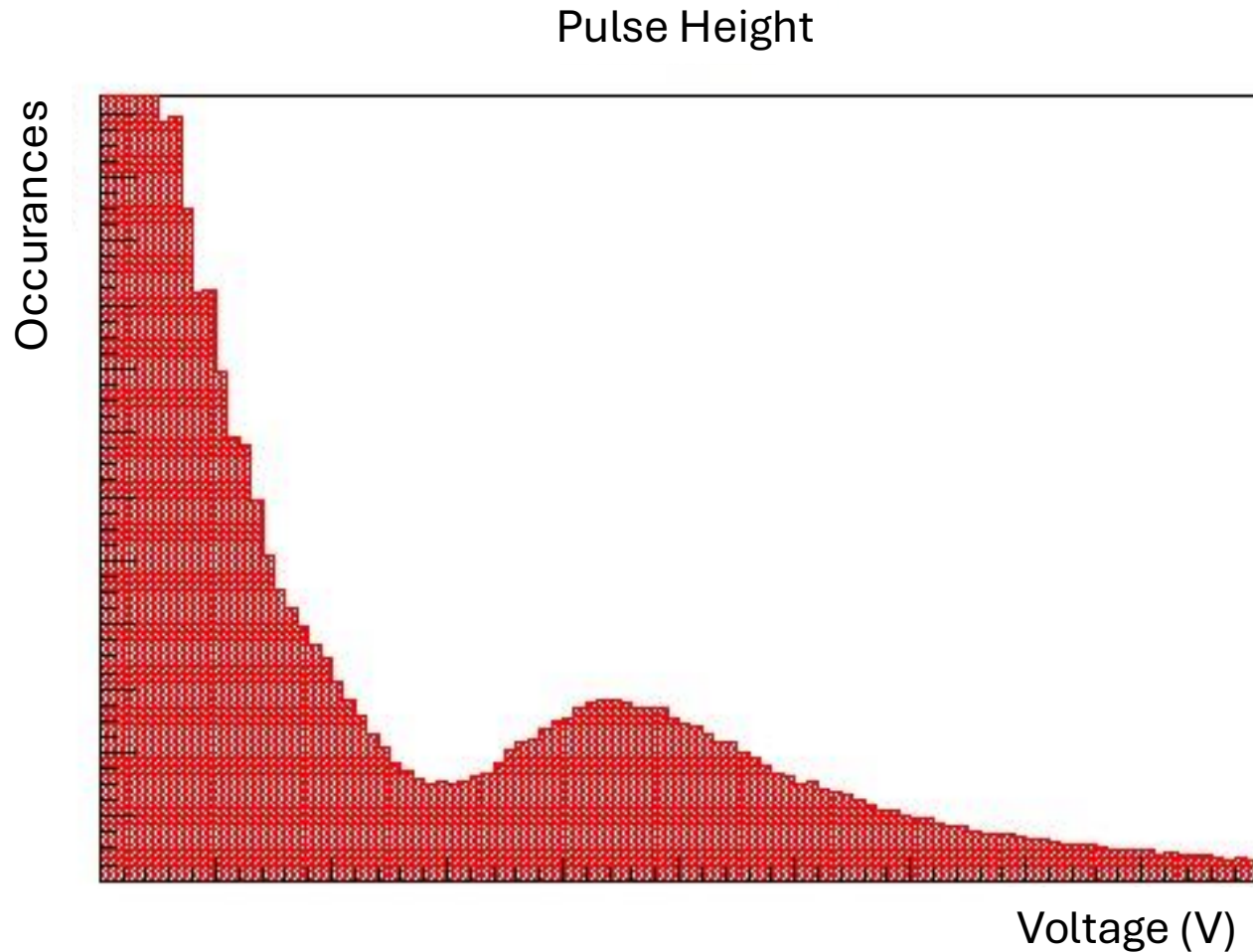
Efficiency PMT $\sim 30\%$

Amplification PMT $\sim 10^5$

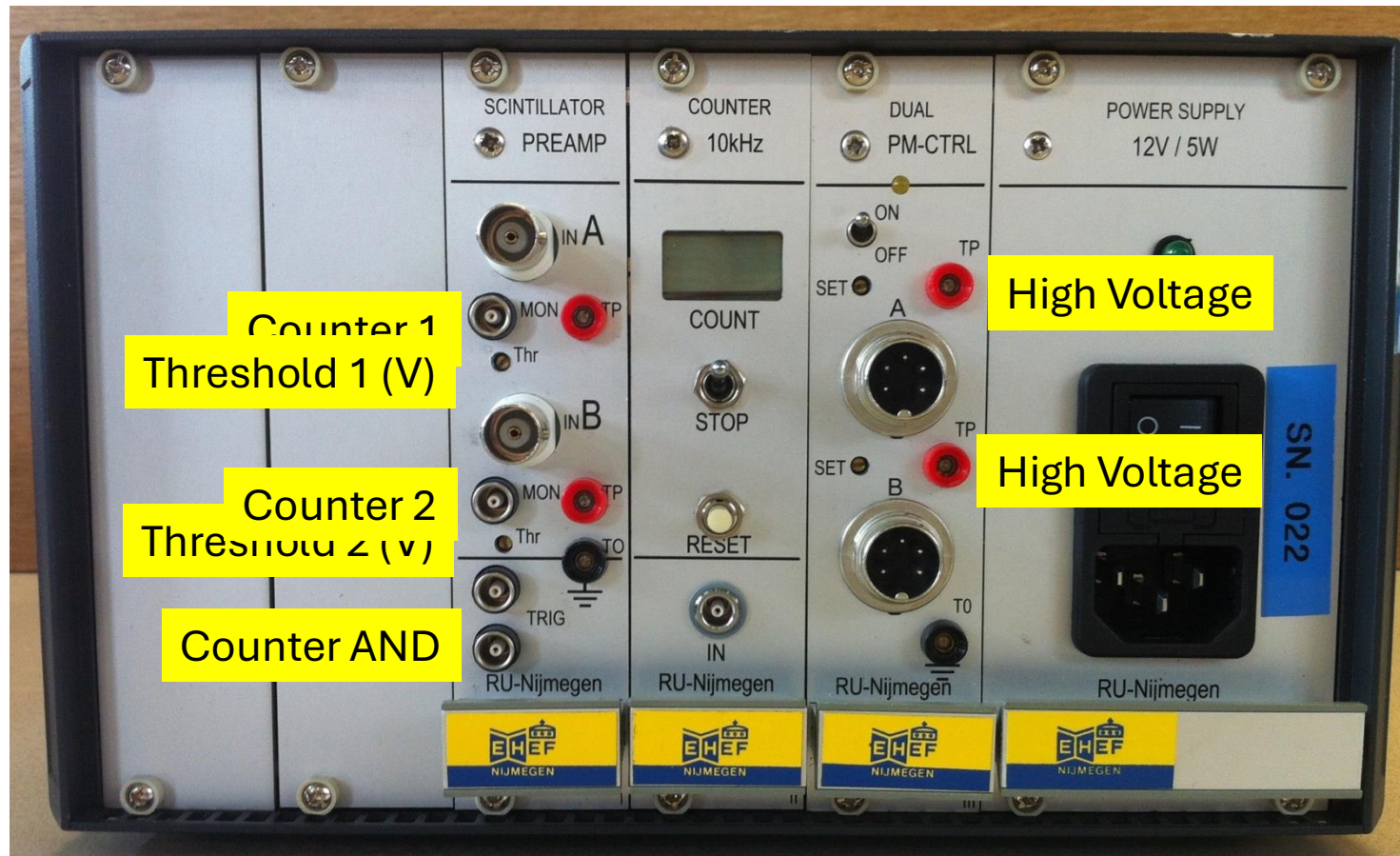


Pulseheight ~ 20 mV
Pulsewidth ~ 25 ns

Hands-on: Determine the (vertical) muon flux



Hands-on: Determine the (vertical) muon flux



Hands-on: Determine the (vertical) muon flux

- Determine good HV values and thresholds
- Determine the efficiency of your setup
- Determine the amount of noise in your setup
- Determine the vertical flux, including its uncertainty