

QCD at 5 Loops & the R^* operation



Franz Herzog

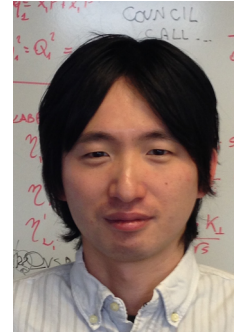
The HEPGAME Team



Giulio Falcioni



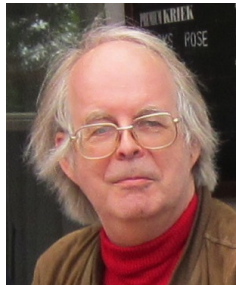
Ben Rujil



Takahiro Ueda



Franz Herzog

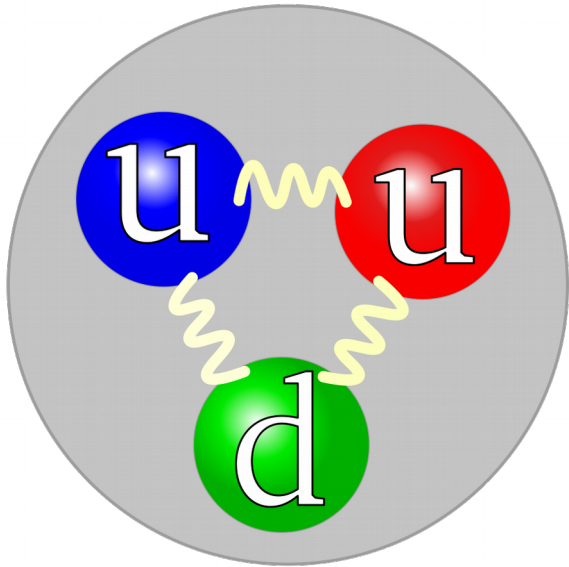


Jos Vermaseren

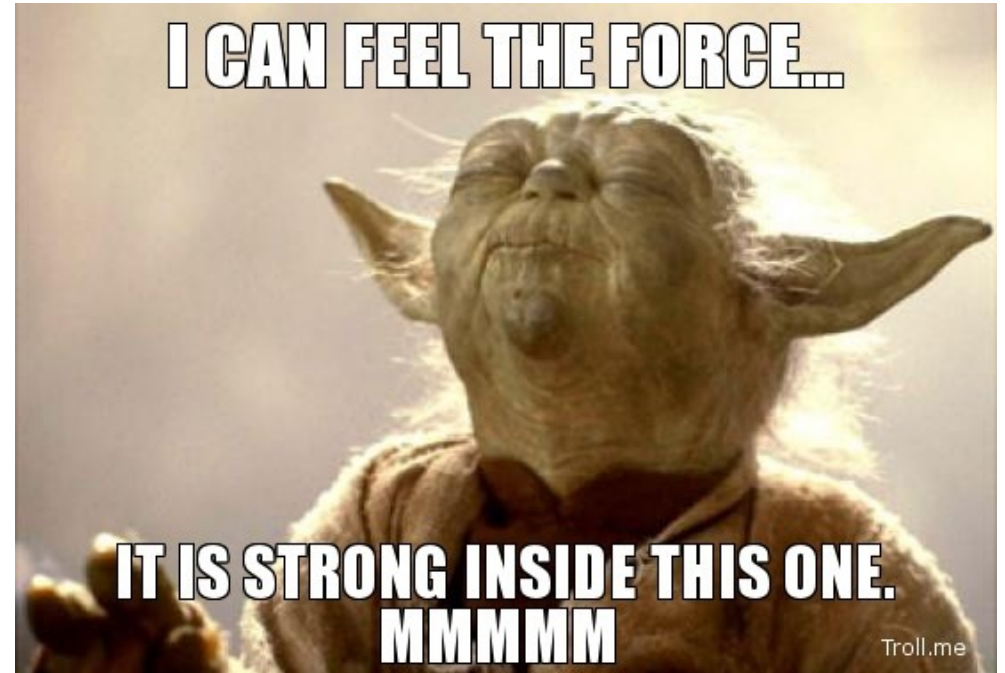


Andreas Vogt

QCD: The strongest force in the universe

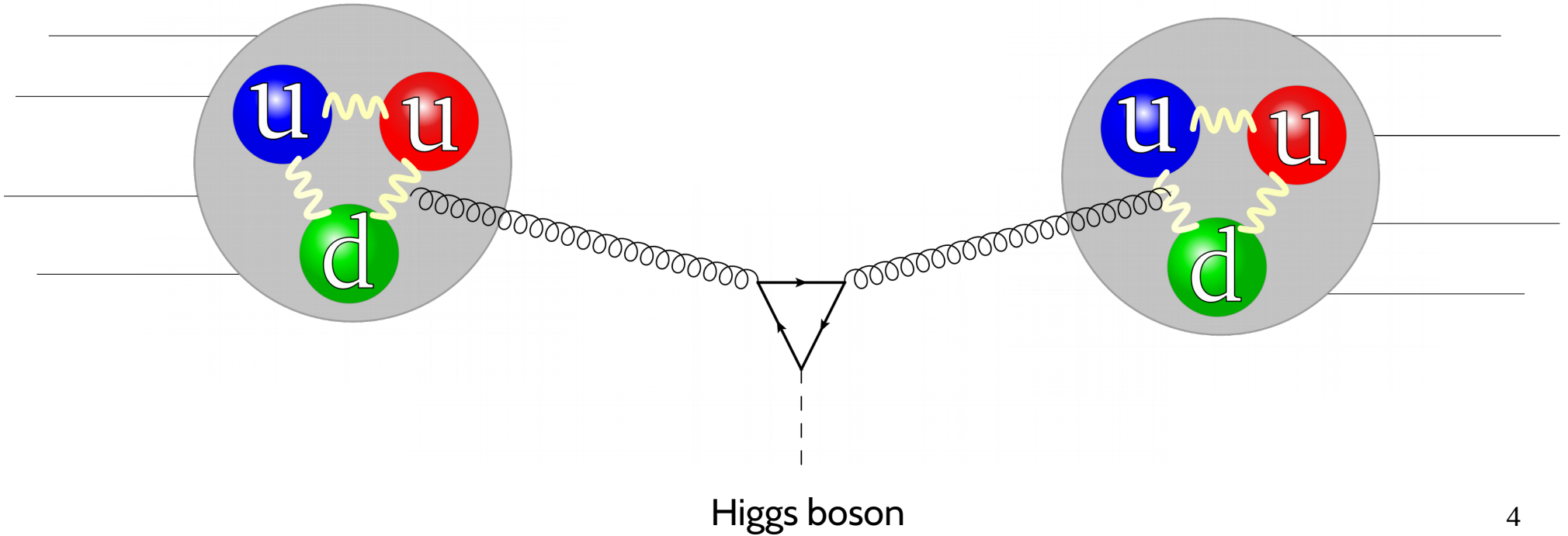


Gluons bind quarks into protons



Proton collisions

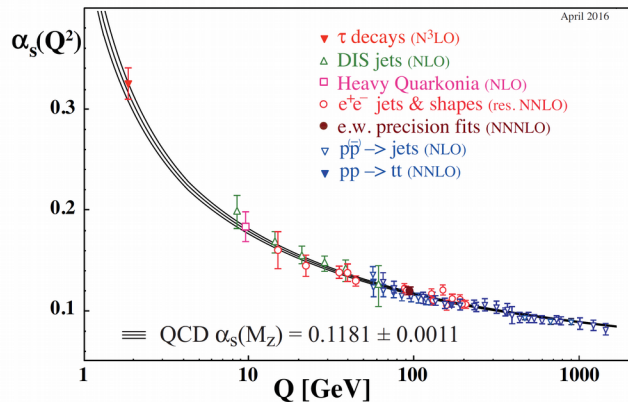
QCD is **CRUCIAL** for collider physics



QCD is asymptotically **free**

- The strong force becomes **weak** at high energy

$$\alpha_s(\mu) \sim \text{[Feynman diagram: two quark lines meeting at a vertex with a gluon line]} < 1$$



The Nobel Prize in Physics 2004

"for the discovery of asymptotic freedom in the theory of the strong interaction"



David J. Gross



H. David Politzer



Frank Wilczek

QCD is asymptotically **free**

➤ weak coupling → **Feynman Diagrams**

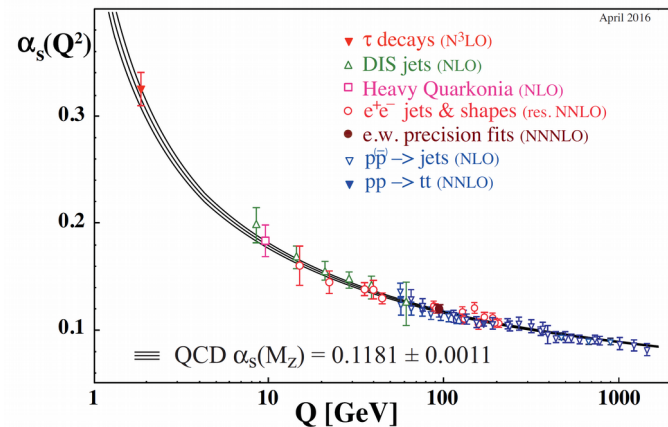
$$\begin{aligned} \mathcal{A}_{gg \rightarrow H} = & \text{tree} + \alpha_s \left(\text{1-loop} + \dots \right) \\ & + \alpha_s^2 \left(\text{2-loop} + \dots \right) \\ & + \alpha_s^3 \left(\text{3-loop} + \dots \right) \\ & + \mathcal{O}(\alpha_s^4) \end{aligned}$$



Q: How did they prove asymptotic freedom?

A: They computed a sign!

$$\beta(Q) = Q^2 \frac{d\alpha_s(Q)}{dQ^2} = -\beta_0 \alpha_s^2 - \beta_1 \alpha_s^3 - \beta_2 \alpha_s^4 - \beta_3 \alpha_s^5 - \beta_4 \alpha_s^6 - \dots$$



Complexity at higher loops

# Loops	1	2	3	4	5
# Diagrams	7	66	1 219	29 310	831 942
# terms of worst diagrams	9	81	729	6 561	59 049
Dimension of integrals	2	5	8	11	14

History of the QCD beta function

➤ β_0 :

- 1965 [Vanyashin, Terentev]; 1970 [Khriplovich]; 1972 [T'Hooft];
- 1973 [Gross, Wilczek; Politzer]

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- β_3 :
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➤ β_4 :

- 2016 [Baikov, Chetyrkin, Kühn]
- 2017 [FH, Ruijl, Ueda, Vermaseren, Vogt]

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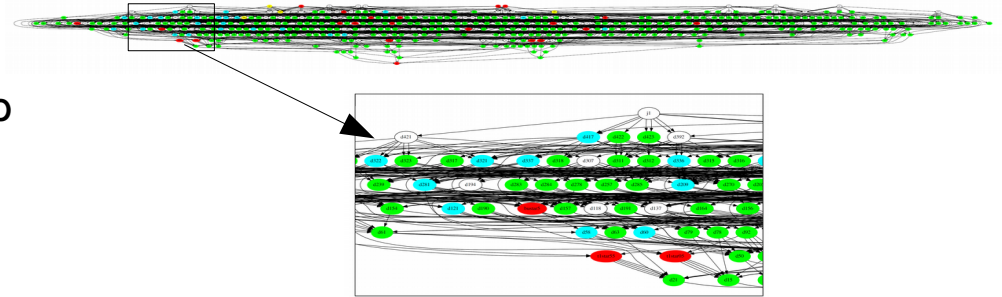
**1 year
4 days**

Computing time

Calculation

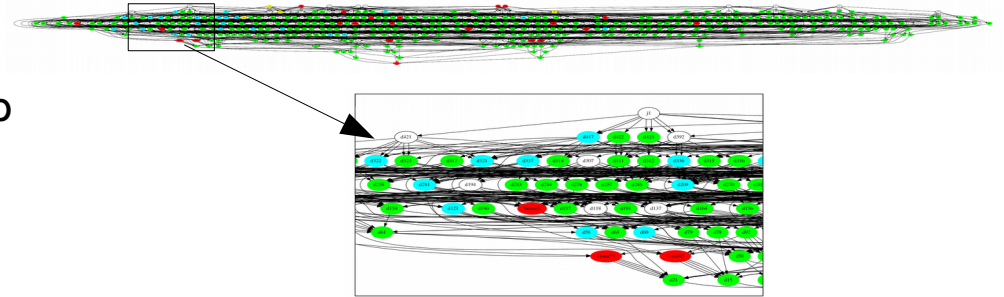
Calculation

- Forcer [arXiv:1704.06650](https://arxiv.org/abs/1704.06650)
 - Parametric solution of IBPs for up to 4-loop massless self energy graphs

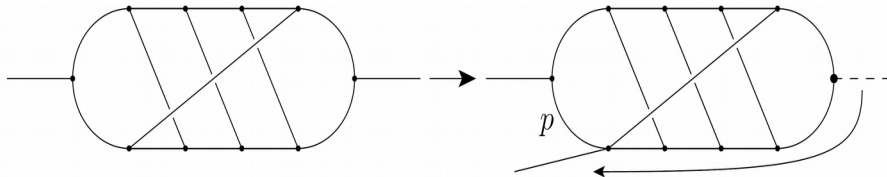


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 - fully automated in Form (& Maple for tests)
 - Allows extraction of poles of L loop euclidean integrals from L-1 loops via infrared rearrangement:

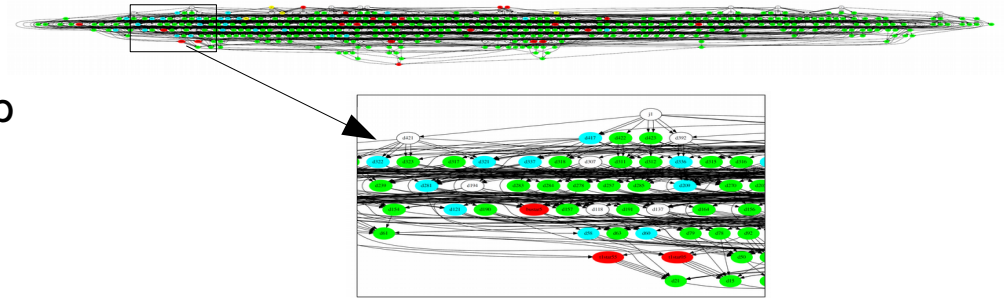


$$R^* \left(\text{Diagram 1} \right) = \text{Diagram 1} + \tilde{\Delta} \left(\text{Diagram 2} \right) \text{Diagram 3} + \tilde{\Delta} \left(\text{Diagram 2} \right) \Delta \left(\text{Diagram 4} \right) + \Delta \left(\text{Diagram 4} \right)$$



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$$R^* \text{ (circle with two dots) } = \text{ (circle with two dots) } + \tilde{\Delta} \left(\text{ (vertical line with two dots) } \right) \text{ (circle with one dot) } + \tilde{\Delta} \left(\text{ (vertical line with two dots) } \right) \Delta \left(\text{ (circle) } \right) + \Delta \left(\text{ (circle with two dots) } \right)$$

- Background field gauge:
 - Extract beta from background field self energy

$$\Pi_B(Q) \sim B \text{ (wavy line) } \text{ (loop diagram with vertices } \hat{A} \text{ and } \hat{A} \text{)} \text{ (wavy line) } B + \dots \quad Z_g \sqrt{Z_B} = 1$$

General gauge group

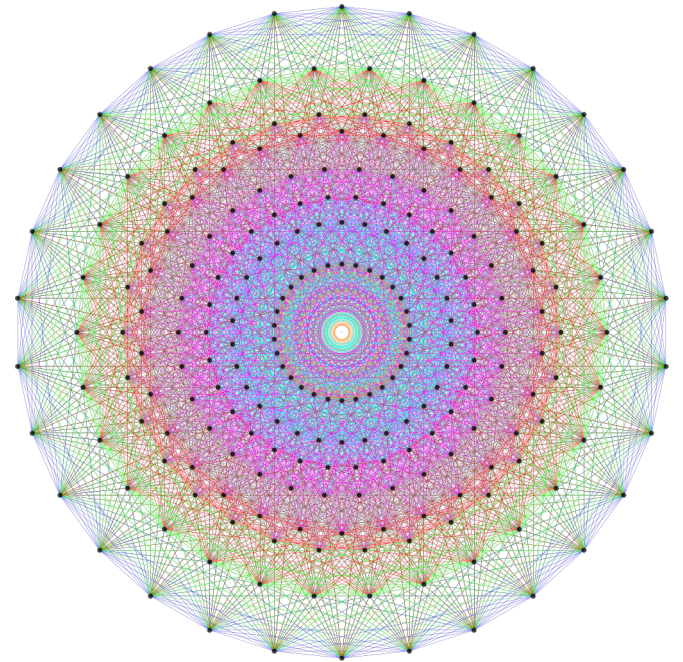
Besides C_A, C_F we only need the symmetric group invariant tensors

T^a are the generators of the fermionic representation.

C^a are the generators of the adjoint representation.

$$d_F^{abcd} = \frac{1}{6} \text{Tr}(T^a T^b T^c T^d + \text{five } bcd \text{ permutations})$$

$$d_A^{abcd} = \frac{1}{6} \text{Tr}(C^a C^b C^c C^d + \text{five } bcd \text{ permutations})$$



Group theory factors for Feynman diagrams

T. van Ritbergen^a, A.N. Schellekens^b, J.A.M. Vermaseren^b

UM-TH-98-01
NIKHEF-98-004

A projection of the 8 dimensional 4_{21} polytope
invariant under the the E_8 Lie group

Old Results

$$\beta_0 = \frac{11}{3} C_A - \frac{4}{3} T_F n_f ,$$

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$$\beta_2 = \frac{2857}{54} C_A^3 - \frac{1415}{27} C_A^2 T_F n_f - \frac{205}{9} C_F C_A T_F n_f + 2 C_F^2 T_F n_f \\ + \frac{44}{9} C_F T_F^2 n_f^2 + \frac{158}{27} C_A T_F^2 n_f^2 ,$$

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$$\beta_3 = C_A^4 \left(\frac{150653}{486} - \frac{44}{9} \zeta_3 \right) + \frac{d^{abcd} d_A^{abcd}}{N_A} \left(-\frac{80}{9} + \frac{704}{3} \zeta_3 \right) \\ + C_A^3 T_F n_f \left(-\frac{39143}{81} + \frac{136}{3} \zeta_3 \right) + C_A^2 C_F T_F n_f \left(\frac{7073}{243} - \frac{656}{9} \zeta_3 \right) \\ + C_A C_F^2 T_F n_f \left(-\frac{4204}{27} + \frac{352}{9} \zeta_3 \right) + \frac{d_F^{abcd} d_A^{abcd}}{N_A} n_f \left(\frac{512}{9} - \frac{1664}{3} \zeta_3 \right) \\ + 46 C_F^3 T_F n_f + C_A^2 T_F^2 n_f^2 \left(\frac{7930}{81} + \frac{224}{9} \zeta_3 \right) + C_F^2 T_F^2 n_f^2 \left(\frac{1352}{27} - \frac{704}{9} \zeta_3 \right) \\ + C_A C_F T_F^2 n_f^2 \left(\frac{17152}{243} + \frac{448}{9} \zeta_3 \right) + \frac{d_F^{abcd} d_F^{abcd}}{N_A} n_f^2 \left(-\frac{704}{9} + \frac{512}{3} \zeta_3 \right) \\ + \frac{424}{243} C_A T_F^3 n_f^3 + \frac{1232}{243} C_F T_F^3 n_f^3 ,$$

New Result

$$\begin{aligned}
\beta_4 = & C_A^5 \left(\frac{8296235}{3888} - \frac{1630}{81} \zeta_3 + \frac{121}{6} \zeta_4 - \frac{1045}{9} \zeta_5 \right) \\
& + \frac{d_A^{abcd} d_A^{abcd}}{N_A} C_A \left(-\frac{514}{3} + \frac{18716}{3} \zeta_3 - 968 \zeta_4 - \frac{15400}{3} \zeta_5 \right) \\
& + C_A^4 T_F n_f \left(-\frac{5048959}{972} + \frac{10505}{81} \zeta_3 - \frac{583}{3} \zeta_4 + 1230 \zeta_5 \right) \\
& + C_A^3 C_F T_F n_f \left(\frac{8141995}{1944} + 146 \zeta_3 + \frac{902}{3} \zeta_4 - \frac{8720}{3} \zeta_5 \right) \\
& + C_A^2 C_F^2 T_F n_f \left(-\frac{548732}{81} - \frac{50581}{27} \zeta_3 - \frac{484}{3} \zeta_4 + \frac{12820}{3} \zeta_5 \right) \\
& + C_A C_F^3 T_F n_f \left(3717 + \frac{5696}{3} \zeta_3 - \frac{7480}{3} \zeta_5 \right) - C_F^4 T_F n_f \left(\frac{4157}{6} + 128 \zeta_3 \right) \\
& + \frac{d_A^{abcd} d_A^{abcd}}{N_A} T_F n_f \left(\frac{904}{9} - \frac{20752}{9} \zeta_3 + 352 \zeta_4 + \frac{4000}{9} \zeta_5 \right) \\
& + \frac{d_F^{abcd} d_A^{abcd}}{N_A} C_A n_f \left(\frac{11312}{9} - \frac{127736}{9} \zeta_3 + 2288 \zeta_4 + \frac{67520}{9} \zeta_5 \right) \\
& + \frac{d_F^{abcd} d_A^{abcd}}{N_A} C_F n_f \left(-320 + \frac{1280}{3} \zeta_3 + \frac{6400}{3} \zeta_5 \right) \\
& + C_A^3 T_F^2 n_f^2 \left(\frac{843067}{486} + \frac{18446}{27} \zeta_3 - \frac{104}{3} \zeta_4 - \frac{2200}{3} \zeta_5 \right) \\
& + C_A^2 C_F T_F^2 n_f^2 \left(\frac{5701}{162} + \frac{26452}{27} \zeta_3 - \frac{944}{3} \zeta_4 + \frac{1600}{3} \zeta_5 \right) \\
& + C_F^2 C_A T_F^2 n_f^2 \left(\frac{31583}{18} - \frac{28628}{27} \zeta_3 + \frac{1144}{3} \zeta_4 - \frac{4400}{3} \zeta_5 \right) \\
& + C_F^3 T_F^2 n_f^2 \left(-\frac{5018}{9} - \frac{2144}{3} \zeta_3 + \frac{4640}{3} \zeta_5 \right) \\
& + \frac{d_F^{abcd} d_A^{abcd}}{N_A} T_F n_f^2 \left(-\frac{3680}{9} + \frac{40160}{9} \zeta_3 - 832 \zeta_4 - \frac{1280}{9} \zeta_5 \right) \\
& + \frac{d_F^{abcd} d_F^{abcd}}{N_A} C_A n_f^2 \left(-\frac{7184}{3} + \frac{40336}{9} \zeta_3 - 704 \zeta_4 + \frac{2240}{9} \zeta_5 \right) \\
& + \frac{d_F^{abcd} d_F^{abcd}}{N_A} C_F n_f^2 \left(\frac{4160}{3} + \frac{5120}{3} \zeta_3 - \frac{12800}{3} \zeta_5 \right) \\
& + C_A^2 T_F^3 n_f^3 \left(-\frac{2077}{27} - \frac{9736}{81} \zeta_3 + \frac{112}{3} \zeta_4 + \frac{320}{9} \zeta_5 \right) \\
& + C_A C_F T_F^3 n_f^3 \left(-\frac{736}{81} - \frac{5680}{27} \zeta_3 + \frac{224}{3} \zeta_4 \right) \\
& + C_F^2 T_F^3 n_f^3 \left(-\frac{9922}{81} + \frac{7616}{27} \zeta_3 - \frac{352}{3} \zeta_4 \right) \\
& + \frac{d_F^{abcd} d_F^{abcd}}{N_A} T_F n_f^3 \left(\frac{3520}{9} - \frac{2624}{3} \zeta_3 + 256 \zeta_4 + \frac{1280}{3} \zeta_5 \right) \\
& + C_A T_F^4 n_f^4 \left(\frac{916}{243} - \frac{640}{81} \zeta_3 \right) - C_F T_F^4 n_f^4 \left(\frac{856}{243} + \frac{128}{27} \zeta_3 \right)
\end{aligned}$$

QCD

[agrees with result of Baikov, Chetyrkin and Kühn]

$$\beta_0 = 11 - \frac{2}{3}n_f, \quad \beta_1 = 102 - \frac{38}{3}n_f,$$

$$\beta_2 = \frac{2857}{2} - \frac{5033}{18}n_f + \frac{325}{54}n_f^2,$$

$$\beta_3 = \frac{149753}{6} + 3564\zeta_3 + n_f \left(-\frac{1078361}{162} - \frac{6508}{27}\zeta_3 \right) \\ + n_f^2 \left(\frac{50065}{162} + \frac{6472}{81}\zeta_3 \right) + \frac{1093}{729}n_f^3$$

$$\beta_4 = \frac{8157455}{16} + \frac{621885}{2}\zeta_3 - \frac{88209}{2}\zeta_4 - 288090\zeta_5 \\ + n_f \left(-\frac{336460813}{1944} - \frac{4811164}{81}\zeta_3 + \frac{33935}{6}\zeta_4 + \frac{1358995}{27}\zeta_5 \right) \\ + n_f^2 \left(\frac{25960913}{1944} + \frac{698531}{81}\zeta_3 - \frac{10526}{9}\zeta_4 - \frac{381760}{81}\zeta_5 \right) \\ + n_f^3 \left(-\frac{630559}{5832} - \frac{48722}{243}\zeta_3 + \frac{1618}{27}\zeta_4 + \frac{460}{9}\zeta_5 \right) + n_f^4 \left(\frac{1205}{2916} - \frac{152}{81}\zeta_3 \right)$$

QED

[agrees with result of Baikov, Chetyrkin, Kühn and Rittinger]

$$\beta_0 = \frac{4}{3} n_f , \quad \beta_1 = 4 n_f , \quad \beta_2 = -2 n_f - \frac{44}{9} n_f^2 ,$$

$$\beta_3 = -46 n_f + n_f^2 \left(\frac{760}{27} - \frac{832}{9} \zeta_3 \right) - \frac{1232}{243} n_f^3$$

$$\begin{aligned} \beta_4 = & n_f \left(\frac{4157}{6} + 128 \zeta_3 \right) + n_f^2 \left(-\frac{7462}{9} - 992 \zeta_3 + 2720 \zeta_5 \right) \\ & + n_f^3 \left(-\frac{21758}{81} + \frac{16000}{27} \zeta_3 - \frac{416}{3} \zeta_4 - \frac{1280}{3} \zeta_5 \right) + n_f^4 \left(\frac{856}{243} + \frac{128}{27} \zeta_3 \right) \end{aligned}$$

Same method applies to decay rates

1 [hep-ph] 4 Jul 2017

On Higgs decays to hadrons and the R-ratio at N⁴LO

F. Herzog^a, B. Ruijl^{a,b}, T. Ueda^a, J.A.M. Vermaseren^a and A. Vogt^c

^a*Nikhef Theory Group
Science Park 105, 1098 XG Amsterdam, The Netherlands*

^b*Leiden Centre of Data Science, Leiden University*

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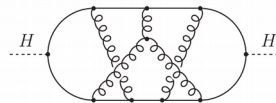
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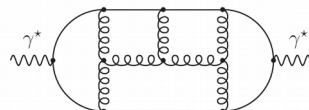
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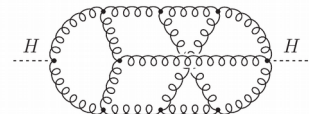
- $H \rightarrow bb$ at N4LO



- Hadronic R-ratio at N4LO



- $H \rightarrow gg$ at N4LO



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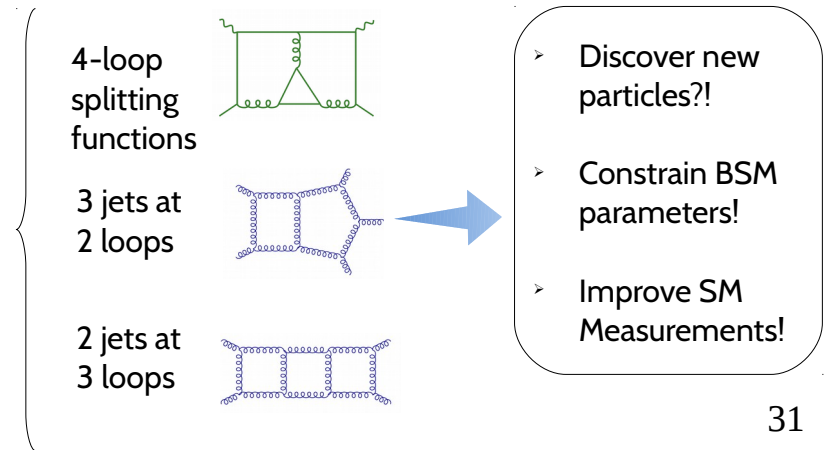
NEW!

First independent confirmation of results
by Baikov, Chetyrkin and Kuehn

Outlook

- R^* :
 - Develop general Framework for computing Renormalisation constants in **any** local QFT
 - With Forcer up to 5 loops!
 - Moments of splitting functions at 4 and 5 loops.
 - SM EFT at 2 loops, even beyond ?
 - Supersymmetry, Gravity, ..
- VID I starting April 2018:
 - Generalise R^* to massless external states:
 - Develop methods for high order LHC cross sections

$$R^*(\Gamma) = \sum_{\substack{S \subseteq \Gamma, \tilde{S} \subseteq \Gamma \\ S \cap \tilde{S} = \emptyset}} \tilde{Z}(\tilde{S}) * Z(S) * \Gamma / S \setminus \tilde{S}$$



n_f dependence

$$\tilde{\beta} \equiv -\beta(a_s)/(a_s^2 \beta_0)$$

$$\tilde{\beta}(\alpha_s, n_f=3) = 1 + 0.565884 \alpha_s + 0.453014 \alpha_s^2 + 0.676967 \alpha_s^3 + 0.580928 \alpha_s^4$$

$$\tilde{\beta}(\alpha_s, n_f=4) = 1 + 0.490197 \alpha_s + 0.308790 \alpha_s^2 + 0.485901 \alpha_s^3 + 0.280601 \alpha_s^4$$

$$\tilde{\beta}(\alpha_s, n_f=5) = 1 + 0.401347 \alpha_s + 0.149427 \alpha_s^2 + 0.317223 \alpha_s^3 + 0.080921 \alpha_s^4$$

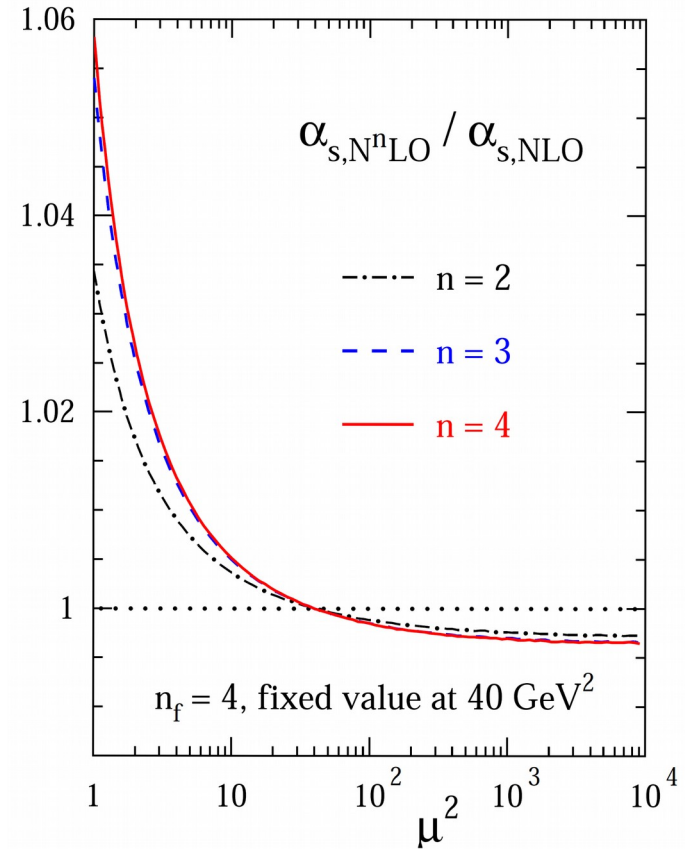
$$\tilde{\beta}(\alpha_s, n_f=6) = 1 + 0.295573 \alpha_s - 0.029401 \alpha_s^2 + 0.177980 \alpha_s^3 + 0.001555 \alpha_s^4$$

Convergence enhanced for larger n_f

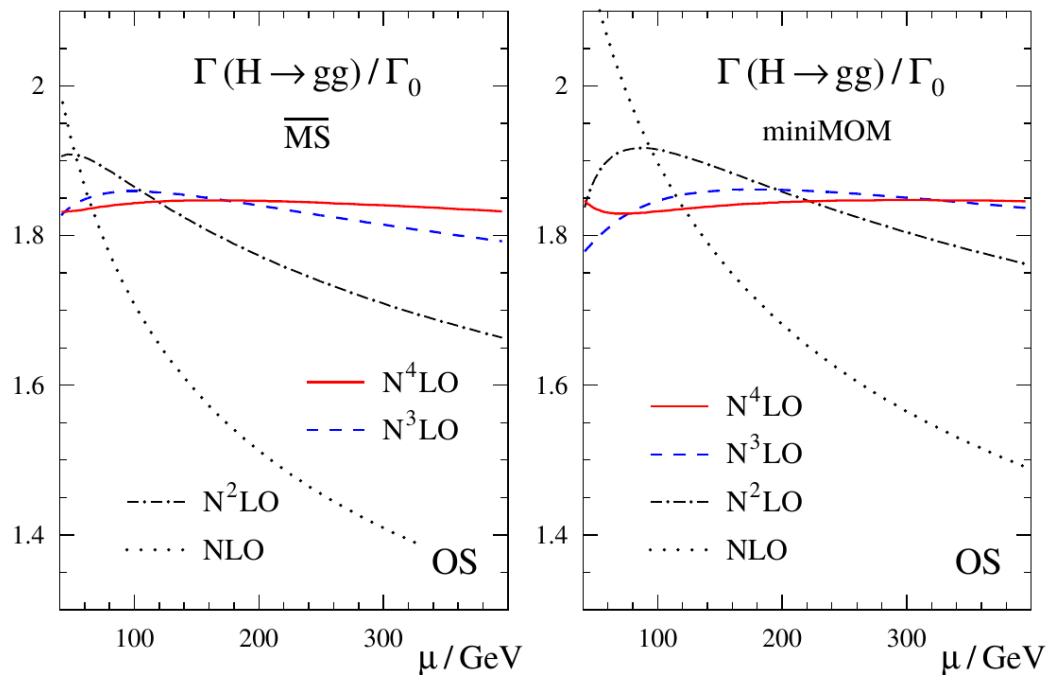
Scale Evolution at low scales

Let us hypothetically fix

$$\alpha_s(6.3\text{GeV}) = 0.2$$



Scale and scheme dependence of $\Gamma(H \rightarrow gg)$



$$\Gamma_{\text{N}^4\text{LO}}(H \rightarrow gg) = \Gamma_0 \left(1.844 \pm 0.011_{\text{series}} \pm 0.045_{\alpha_s(M_Z), 1\%} \right)$$