

Heavy-ion physics, part 2

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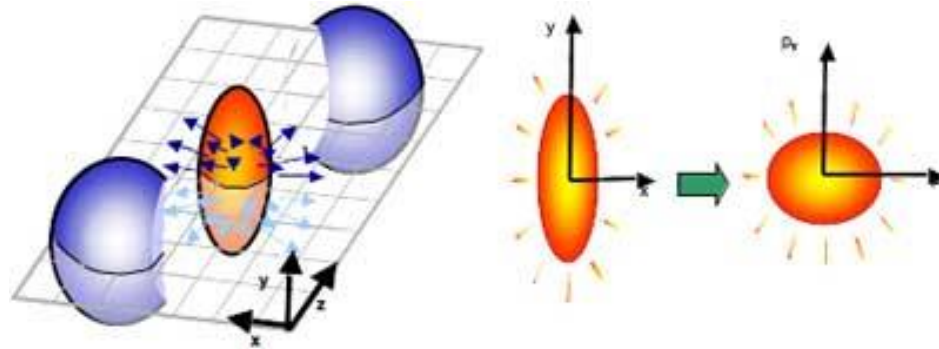
Part 2

Probes of the quark-gluon plasma

- Collective flow
- Jet quenching
- Quarkonia melting

Elliptic flow

In non-central collisions the overlap zone takes an almond shape



Given rapid thermalization, larger pressure along short axis than long one
This *spatial asymmetry* gives rise to a *momentum space anisotropy*

Analyzed by a Fourier decomposition of the particle distribution vs. azimuth

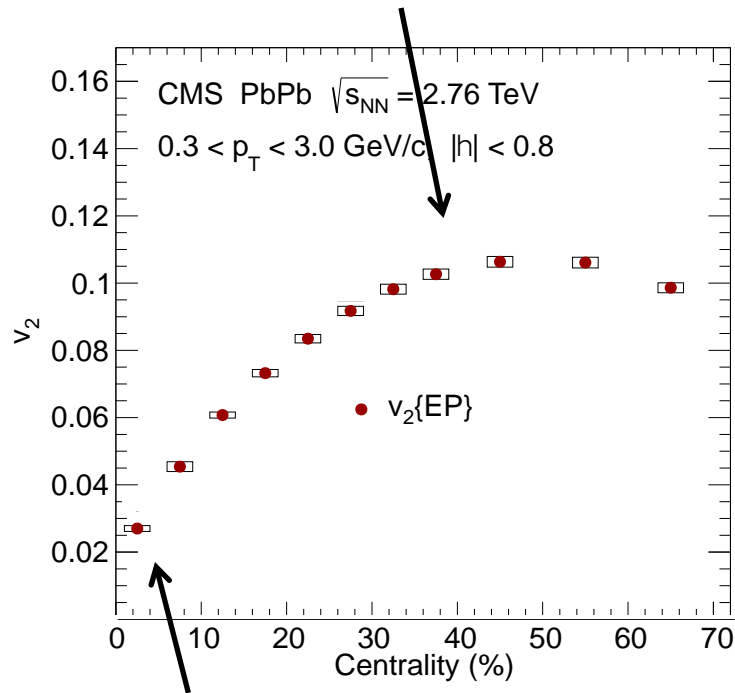
$$E \frac{d^3 \sigma}{d\mathbf{p}^3} = \frac{1}{2\pi} \frac{d^2 \sigma}{p_T dp_T dy} \left(1 + 2 \sum_n v_n \cos[n(\phi - \psi_R)] \right)$$

Given the overlap shape, 2nd harmonic v_2 dominates: “elliptic flow”

Naively odd terms not expected to contribute due to symmetry

Elliptic flow vs. centrality

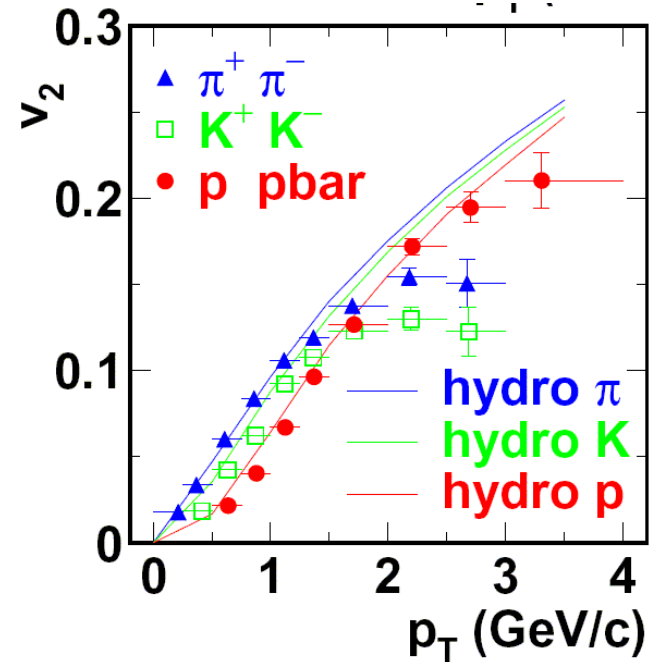
v_2 data conforms to basic expectations:
reaches peak in semi-central collisions



Small in central events
due to symmetric overlap

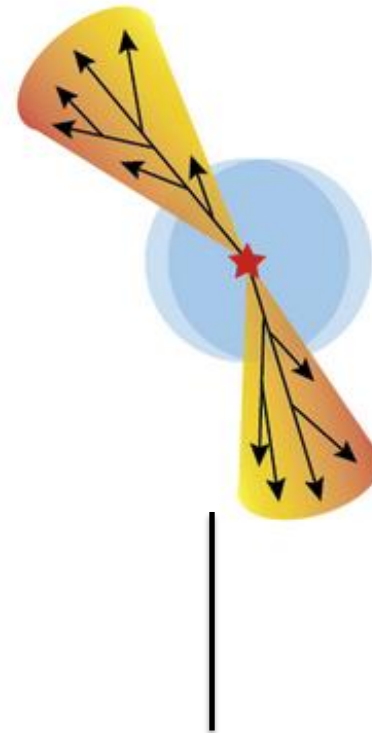
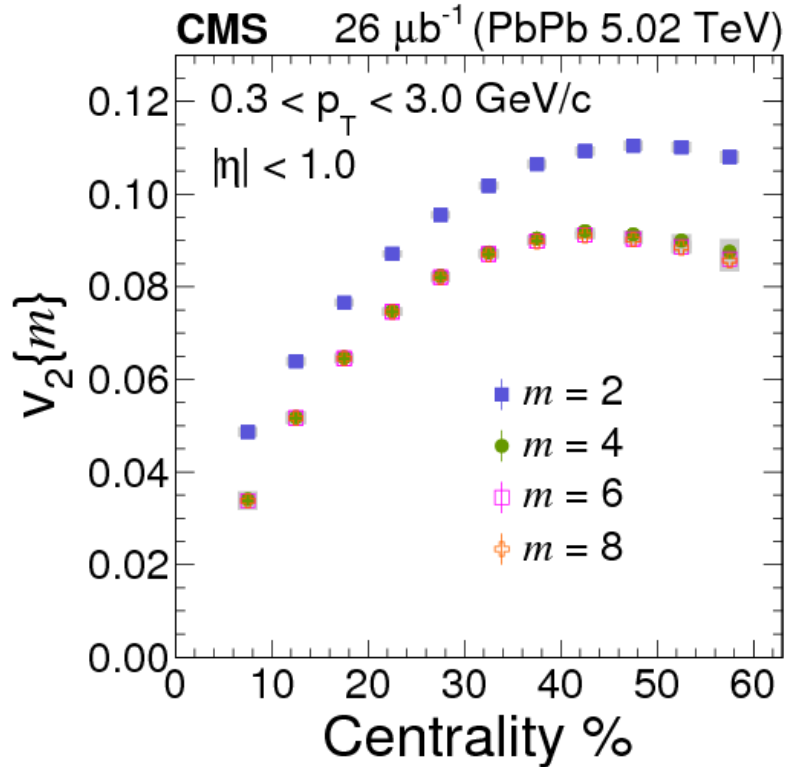
Naively, expect elliptic flow to turn off in peripheral events
where system should be too small to thermalize

Lighter particles more
easily swept up in flow



Flow effects dominant
until about 2 GeV

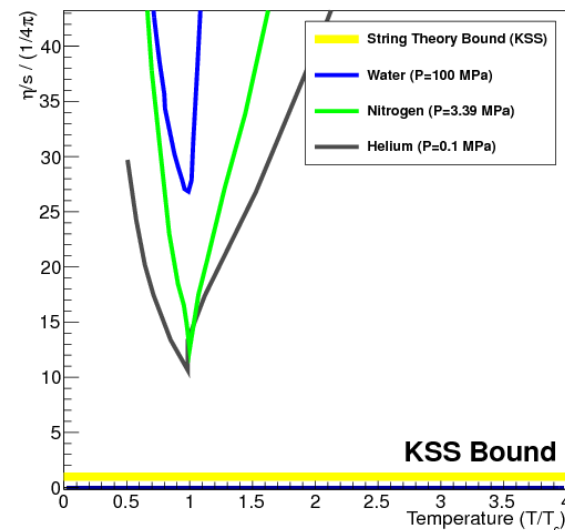
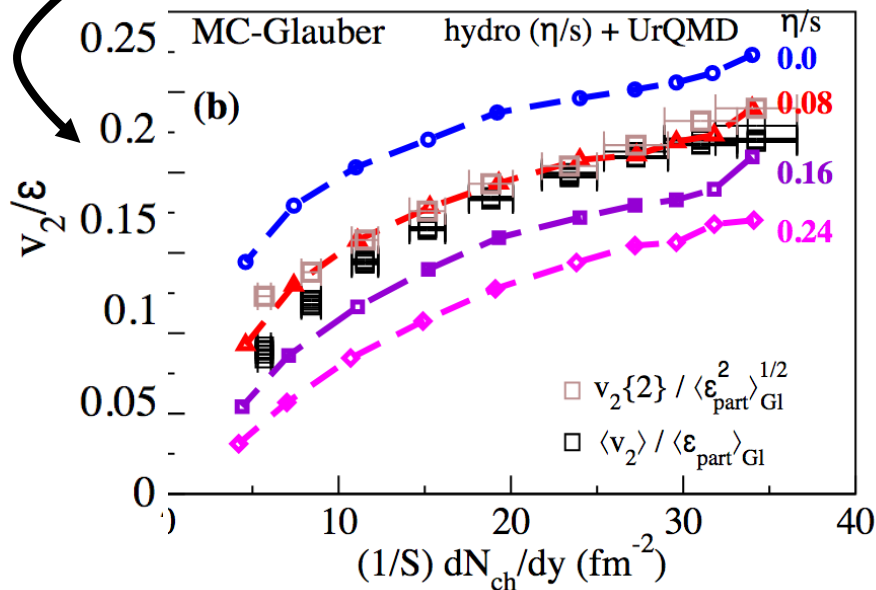
Multiparticle correlations



For particle pairs, strong contribution from “non-flow”, e.g. jets
Minimized by studying multiparticle correlations via “cumulants”

Hydrodynamics

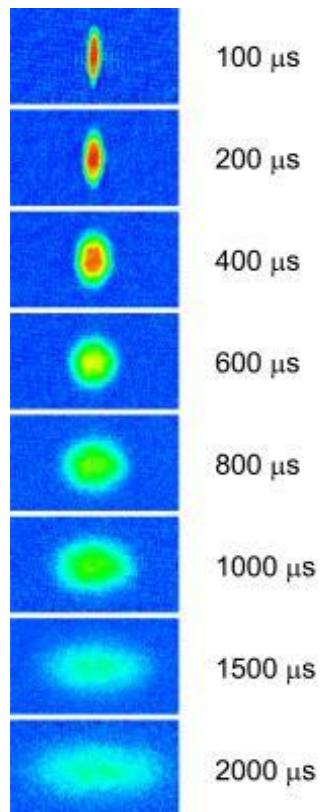
- Collective flow is modeled by relativistic hydrodynamics
- Viscous terms are corrections to ideal hydrodynamics
- Dissipation from shear viscosity would damp elliptic flow
- Sophisticated models couple hydro w/ hadronic rescattering phase



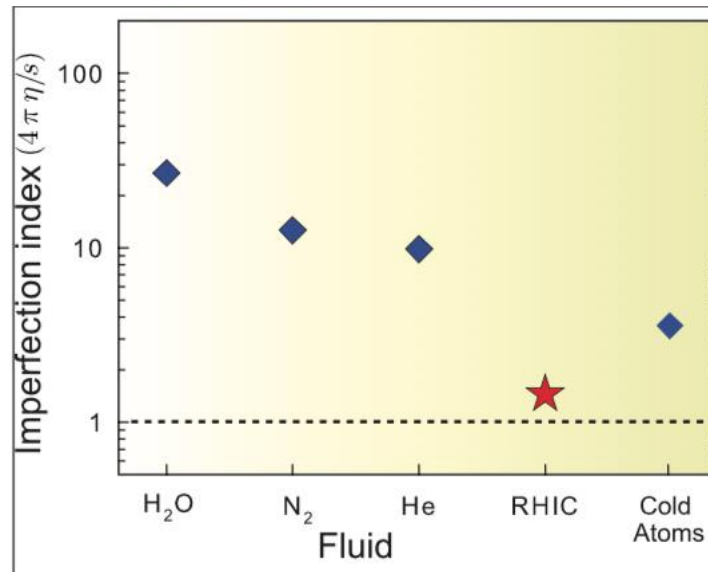
- The ratio η/s is known to have a minimum near phase transition
- Data indicates that viscous corrections are very small $\rightarrow \eta/s \sim 0.08$

Perfect fluidity

Strongly-interacting Fermi gasses display similar anisotropic expansion



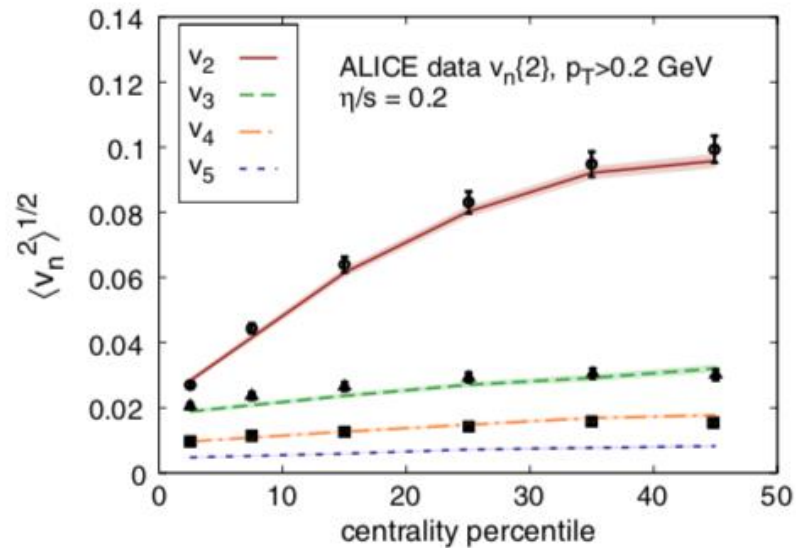
However, viscous corrections are smaller for the QGP than other known systems



Conformal field theory (via *AdS/CFT duality*) gives η/s of $1/4\pi \sim 0.08$

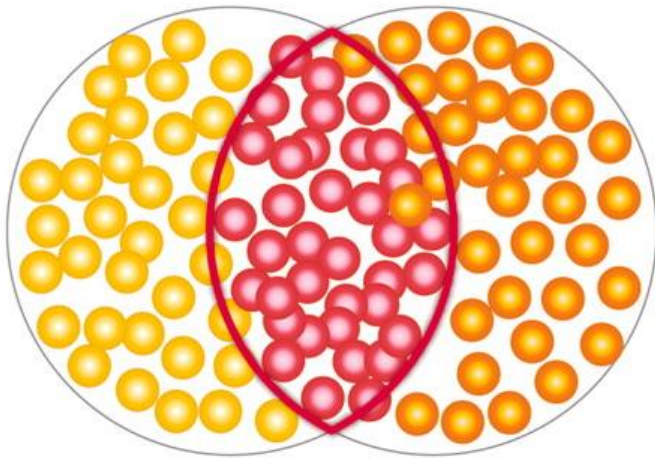
This is conjectured to be a universal lower bound 7

Higher harmonics

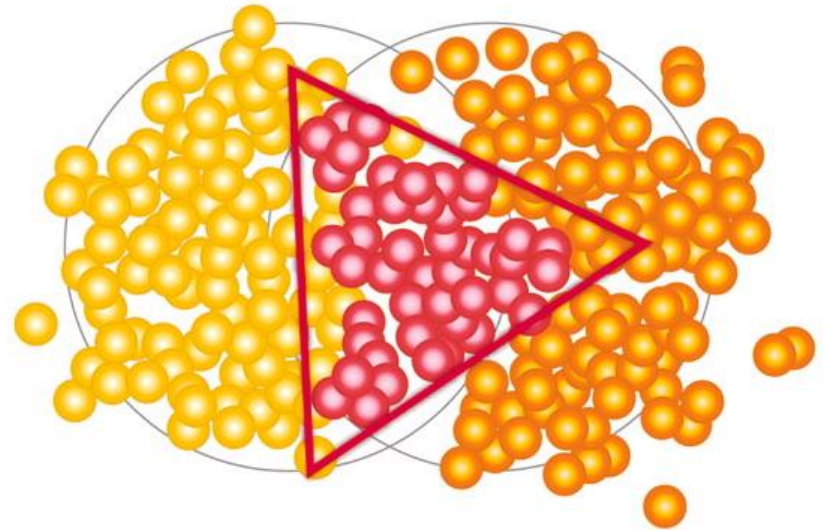


More recently higher harmonics found to be surprisingly large (v_3, v_4)
What is physical picture behind these components of flow?

Geometry fluctuations



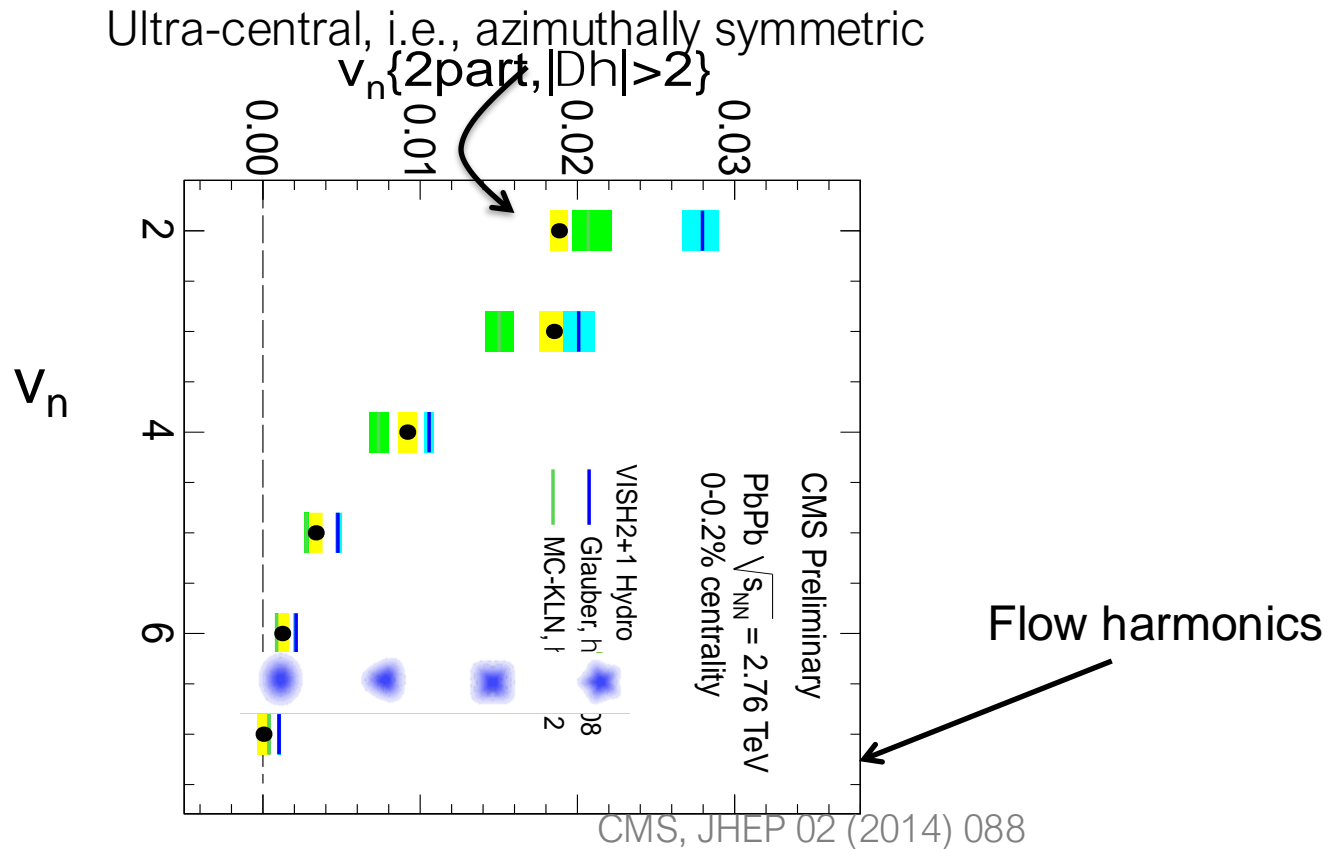
Elliptic flow



Triangular flow

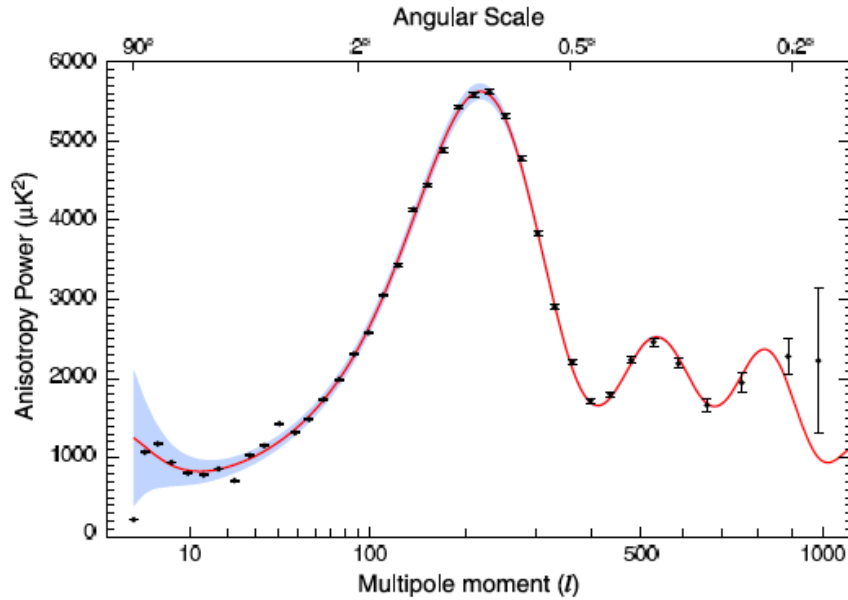
Higher harmonics explained by event-by-event fluctuations of overlap shape
These fluctuations arise naturally in MC implementation of Glauber model
Persist even when collision is azimuthally symmetric

Mapping initial state fluctuations



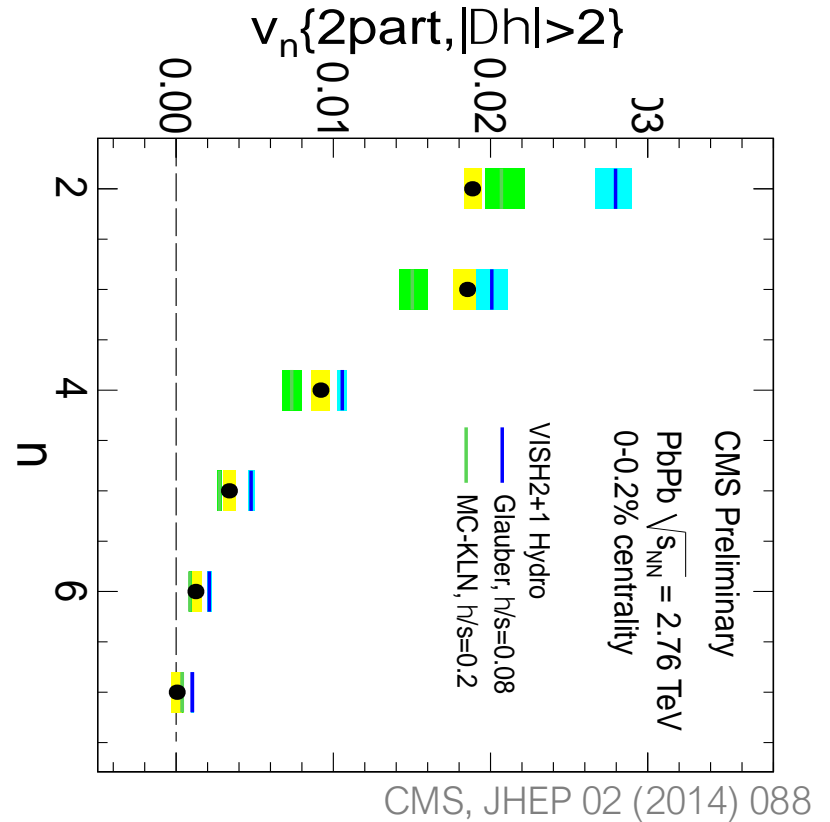
In addition to shear viscosity, higher harmonics sensitive to initial conditions
 Standard Glauber conditions give very low η/s , close to conjectured bound
 But “lumpier” initial conditions in MC-KLN correspond to larger η/s

Analogy to CMB



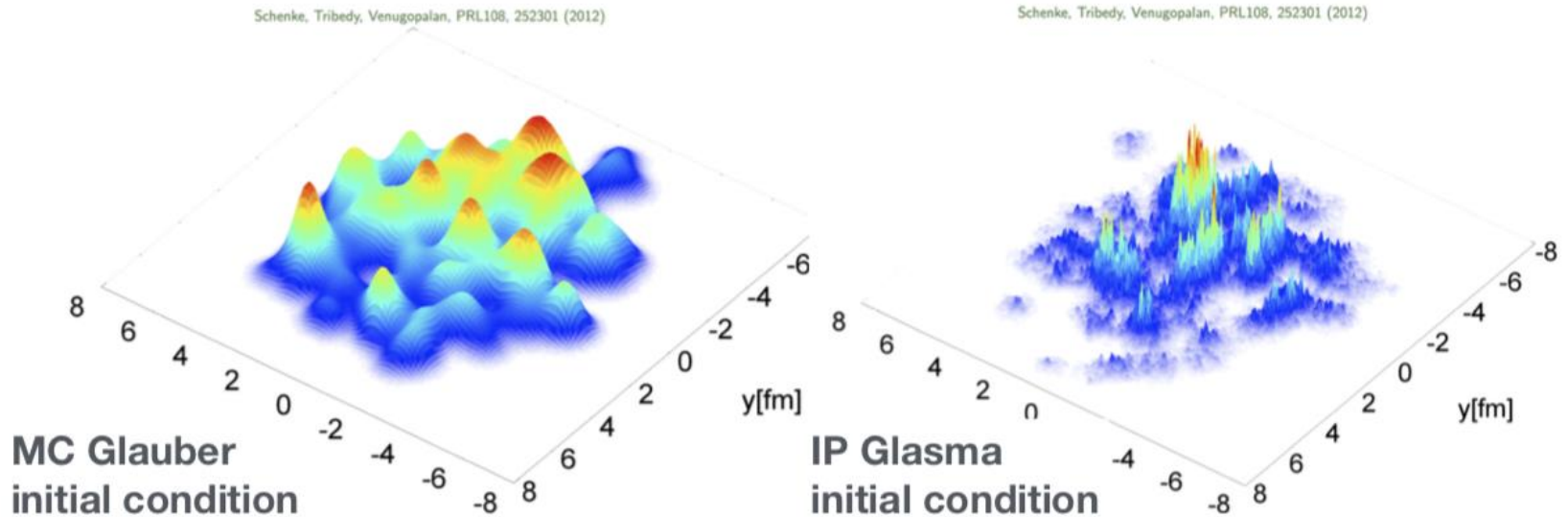
NASA/WMAP Science team

Power spectrum of the cosmic microwave background reflect fluctuations in the density of matter in the early universe before inflation



Fourier harmonics from long-range correlations reflect **fluctuations of the initial state** of dense QCD matter before hydrodynamic expansion

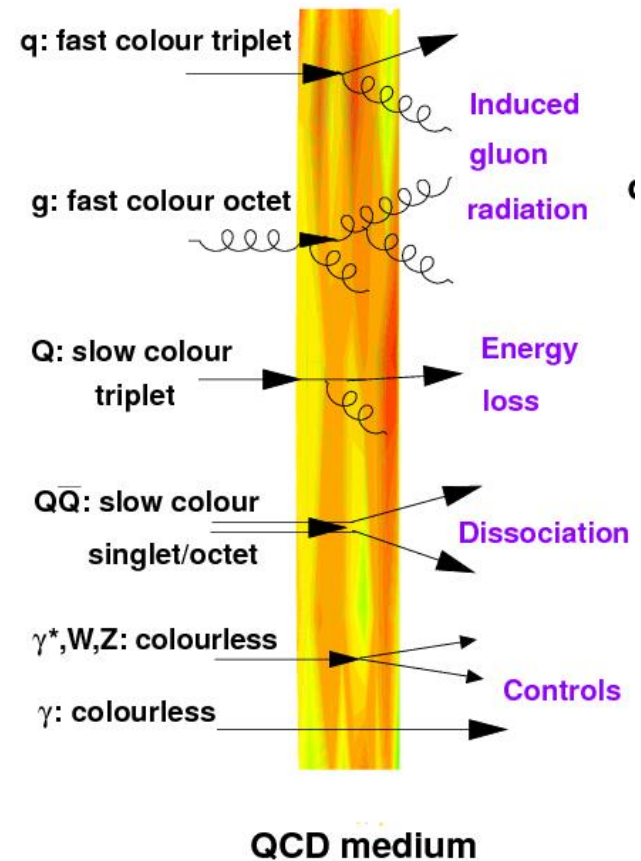
Initial conditions



- Glauber treatment only considers fluctuations at the scale of the nucleon
- Sub-nuclear fluctuations at the level of the gluon field may be important; such effects should be present in small systems (pA or even pp)

“Hard Probes” of the QGP

- Discussed “bulk observables” that affect the majority of particles that are produced in heavy-ion collisions
- Would be nice to scatter particles off the QGP, à la Rutherford
- Next best thing to a particle gun: hard processes
 - High momentum partons \rightarrow jets
 - Heavy quarks, not produced thermally
 - Bound states of heavy quarks, whose binding may be screened
 - Colorless probes

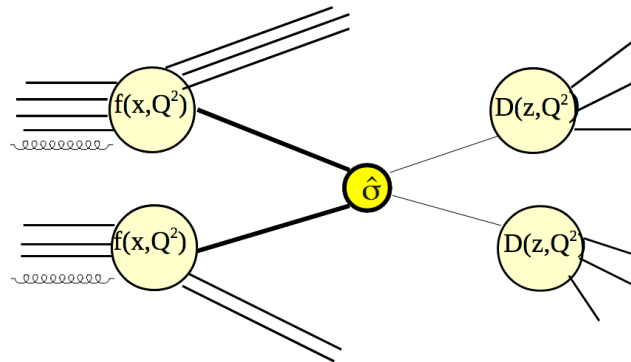


Jets in elementary collisions

High p_T particle production calculable via collinear factorization

Parton scattering cross-section

Parton distribution functions



Fragmentation functions

$$d\sigma = f(x_1, Q^2) \otimes f(x_2, Q^2) \otimes d\hat{\sigma} \otimes D(z, Q^2)$$

PDF and FF are not calculable, but are universal

A jet the output of a clustering algorithm, removes sensitivity to hadronization

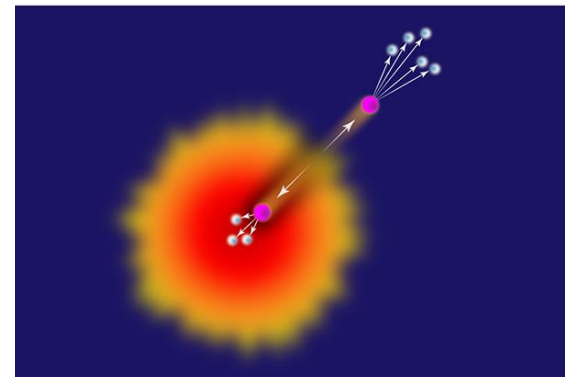
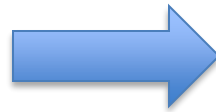
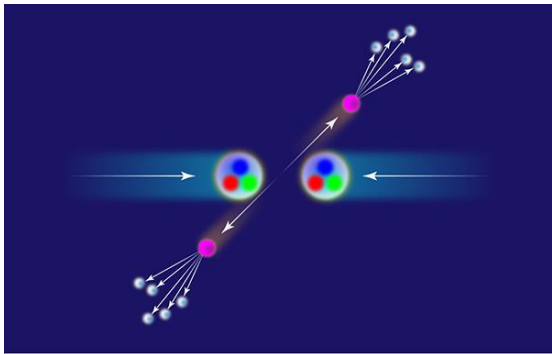
Jet quenching

$$d\sigma = f(x_1, Q^2) \otimes f(x_2, Q^2) \otimes d\hat{\sigma} \otimes P(\Delta E) \otimes D(z', Q^2)$$

PDFs are modified in nuclei,
but relatively small effect

Outgoing parton should lose
energy as they cross the QGP

Assuming partons fragment outside QGP,
fragmentation unmodified (but shifted)

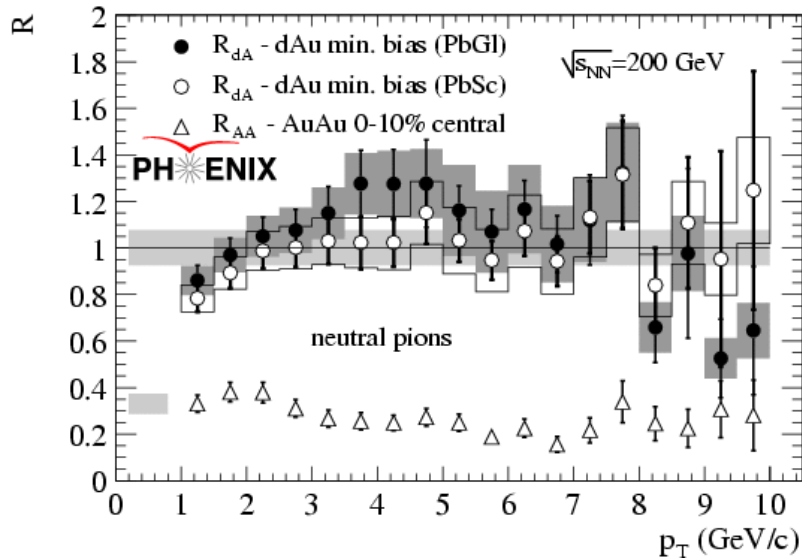


High p_T hadron suppression

Nuclear Modification Factor:

$$R_{AA} \equiv \frac{dN^{AA}/dp_T}{\langle N_{\text{coll}} \rangle dN^{pp}/dp_T} = \frac{dN^{AA}/dp_T}{\langle T_{AA} \rangle d\sigma^{pp}/dp_T}$$

PHENIX measured $\pi^0 \rightarrow 2\gamma$
99% branching ratio



$R_{AA} = 1$ indicates “ N_{coll} scaling”,
i.e., no nuclear effects

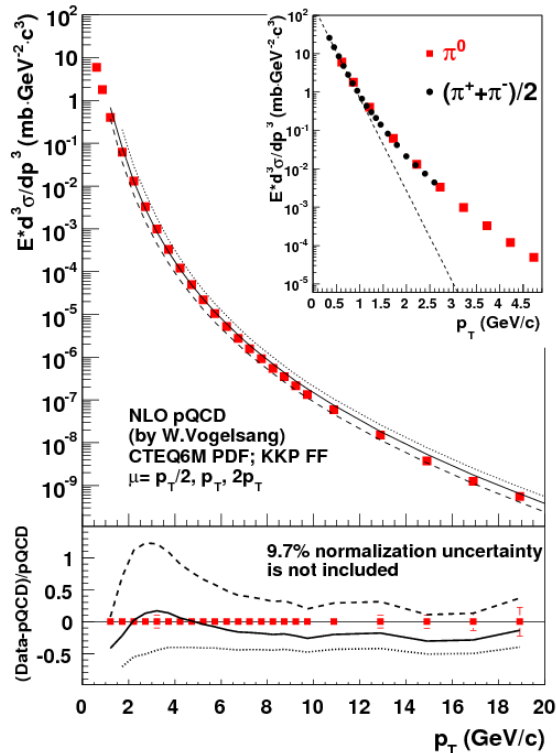
Deuteron-Gold collisions
 → No large parton energy loss
 → No large initial state effects,
 e.g., modification of parton
 distribution inside nuclei

Au-Au yield suppressed by factor of 5!
 Roughly independent of p_T

Fractional energy loss

What would a flat R_{AA} imply for e-loss?

High p_T spectra typically well-described by a power-law



$$\frac{d\sigma}{dp_T} \propto \frac{1}{p_T^n}$$

RHIC: $n \approx 7$
 LHC: $n \approx 5$

Assume *fractional* e-loss: $\Delta p_T = C \cdot p_T$

$$R_{AA} = C^{n-1}$$

e.g., $R_{AA} = 0.2$ @ RHIC
 $\rightarrow C = 0.8$ (20% e-loss)

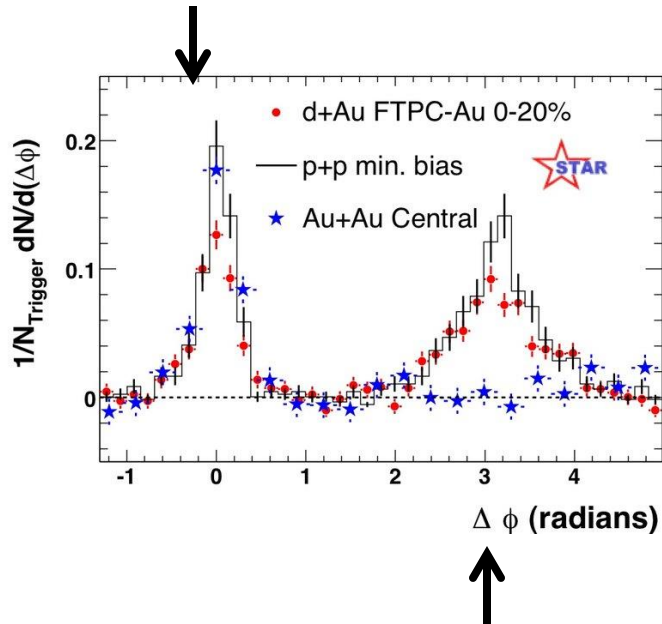
Flat R_{AA} consistent with fractional e-loss

R_{AA} depends on both e-loss and spectral shape

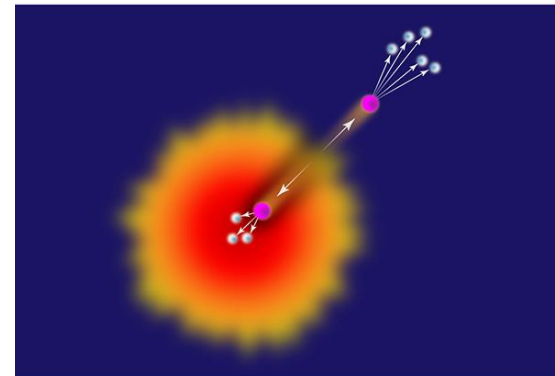
Dihadron azimuthal correlations

Small angle correlation from hadrons within a single jet

Select *trigger* hadron,
 $4 < p_T^{\text{trig}} < 6 \text{ GeV}$
 $\otimes \sqrt{}$ to partner hadron,
 $2 < p_T < p_T^{\text{trig}}$



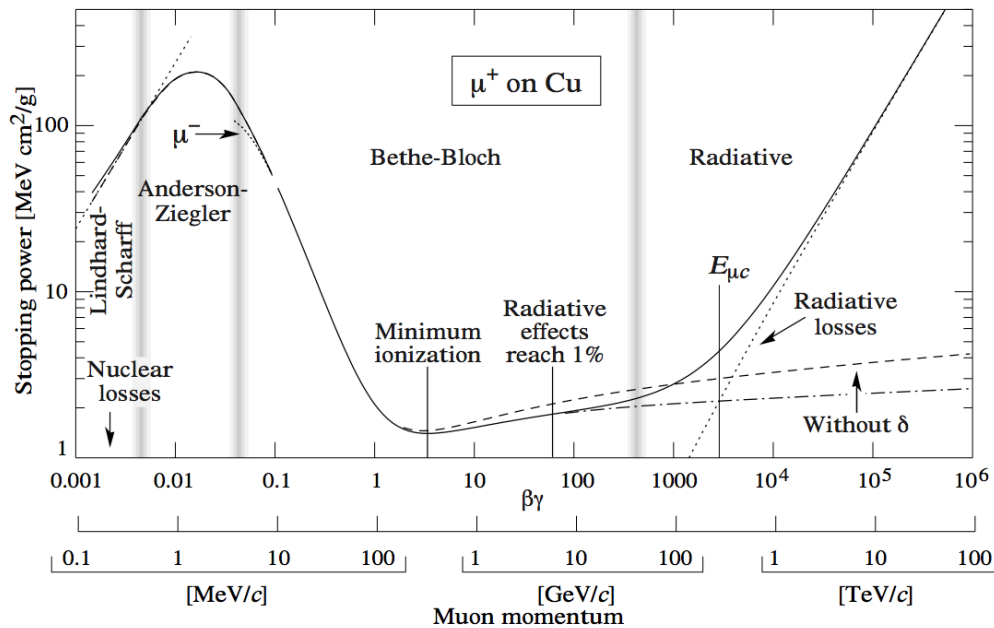
Back-to-back from recoiling jets
 Disappears in central Au-Au!



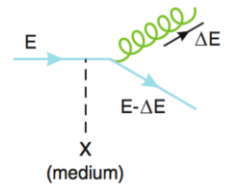
Interpretation: Trigger hadron preferentially produced near surface, while recoil jet traverses the QGP

Energy loss in QCD matter

Passage of particles thru (QED) matter



At large energy radiative e-loss should dominate
 → gluon bremsstrahlung



For mean free path $\lambda \ll$ medium size L
 we are in the LPM regime

$$\Delta E \propto \alpha_s C_F \hat{q} L^2$$

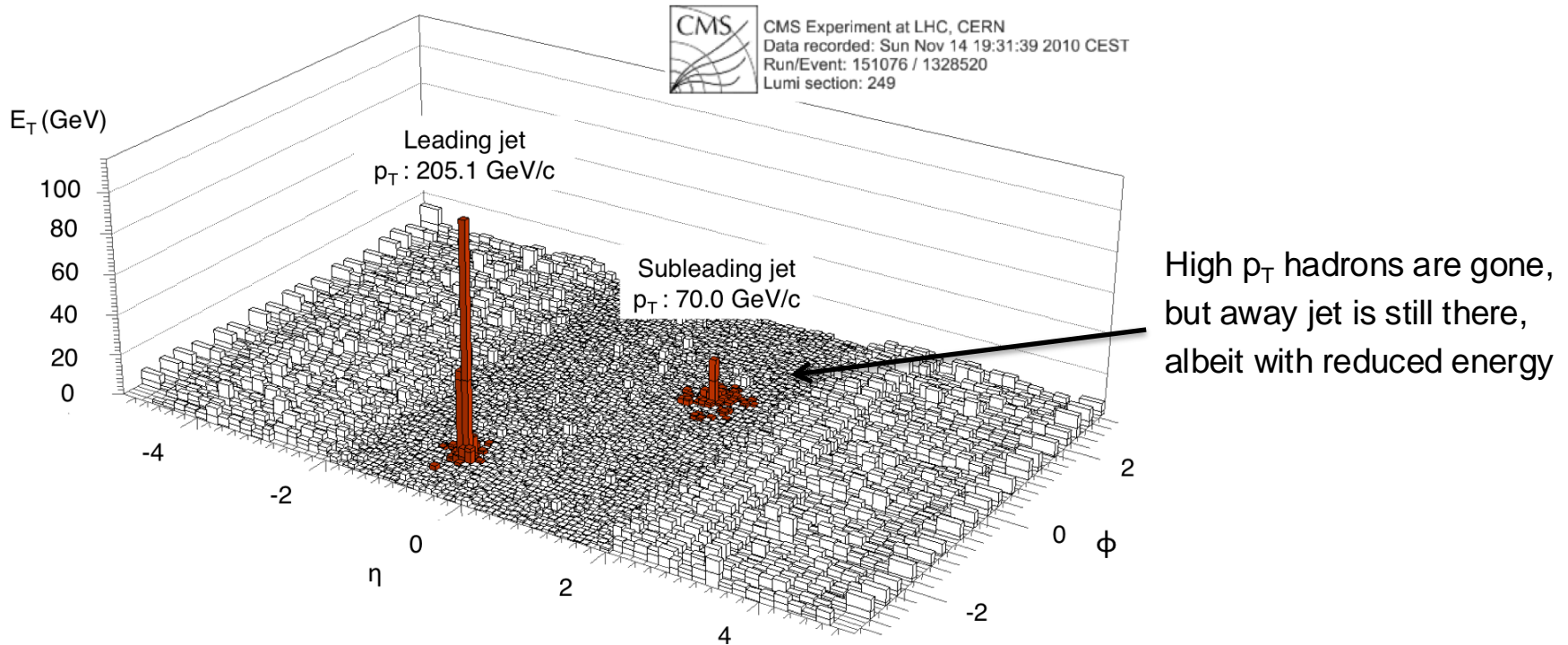
L^2 instead of L , due to destructive interference between scattering centers (LPM effect)

Quenching depends on color factor

$$C_F = \begin{cases} 3 & \text{gluon jets} \\ 4/3 & \text{quark jets} \end{cases}$$

Stopping power q depends on medium density
 → jets are a tomographic probe of the QGP

Fully reconstructed jets

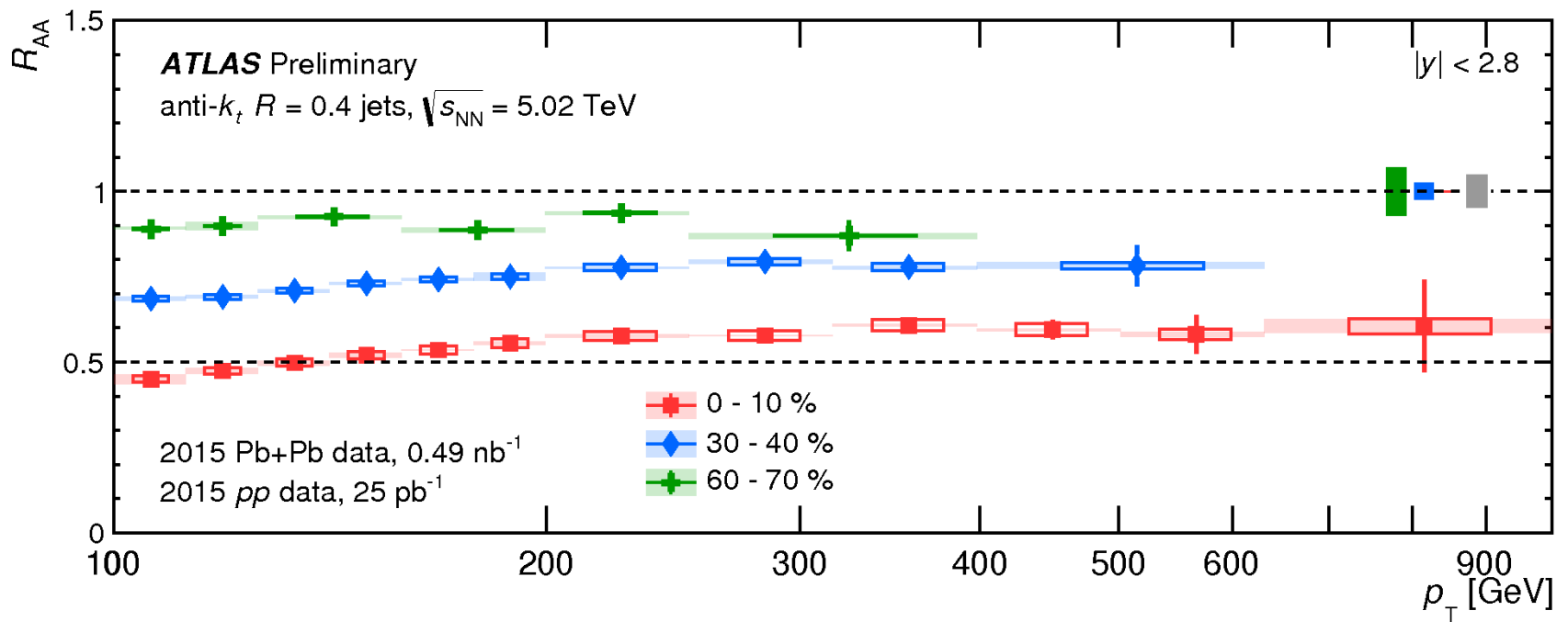


Jet reconstruction become possible at the LHC for the first time in heavy ions

- Hard scattering more abundant at larger collision energy
- Availability of large acceptance, hermetic calorimeters in CMS & ATLAS

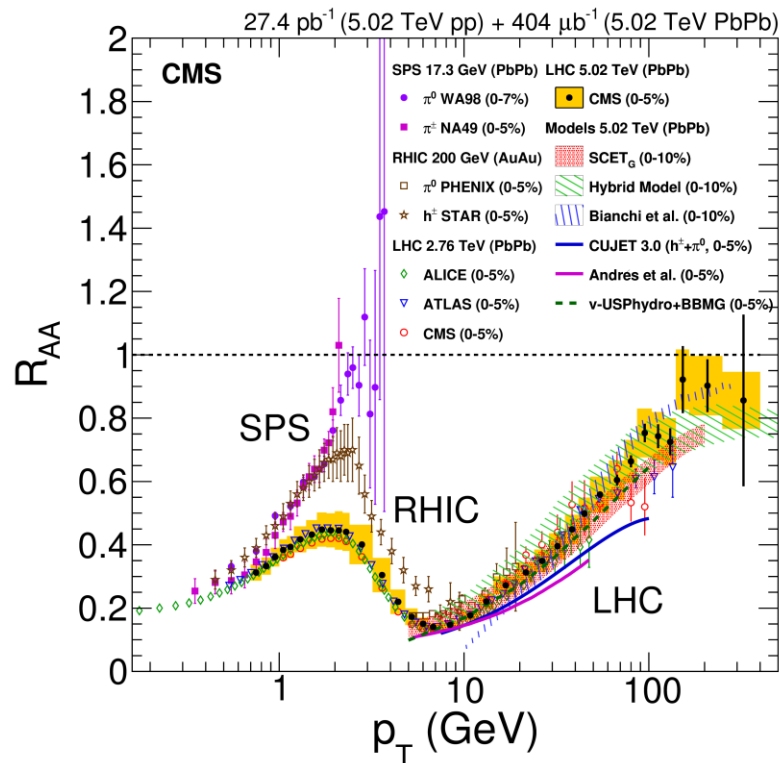
Jet R_{AA}

Wide range of jet momenta probed by PbPb data at the LHC



At large p_T , remarkably flat R_{AA} in central collisions,
reminiscent of the hadron suppression at RHIC

Hadrons at the LHC



Instead a slow but steady rise, showing little suppression at 100 GeV!

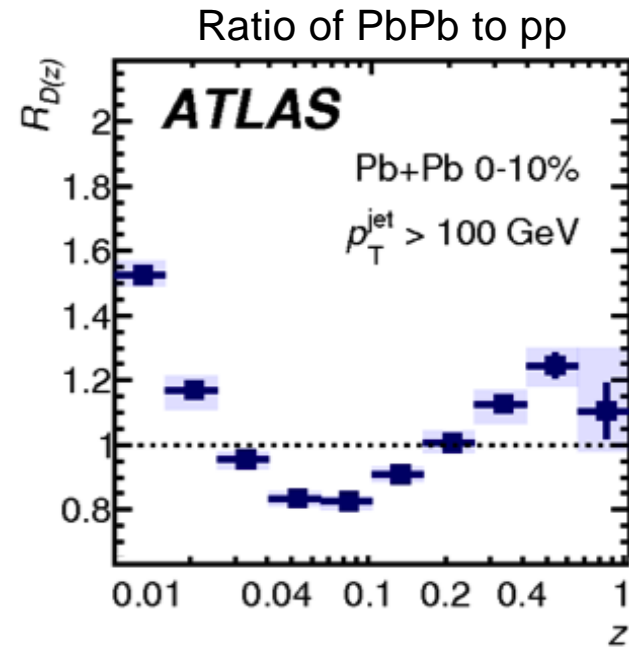
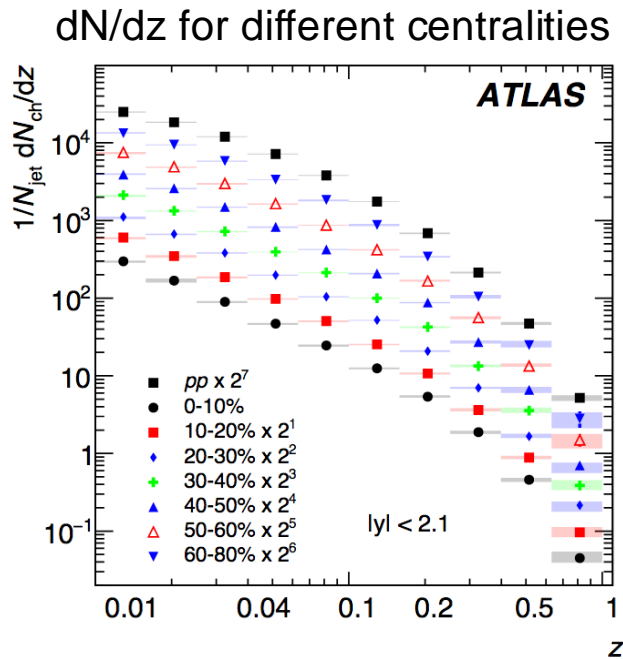
How to reconcile charged hadron and jet suppression?

→ Must be dependence of quenching on fragmentation pattern of jet

Jet fragmentation functions

Not a true FF, jet used as proxy for parton $z = p_{T,\text{hadron}} / p_{T,\text{jet}}$

Not trivial: Quenching may transfer some energy outside jet



Excess at low z : likely dominated by medium response

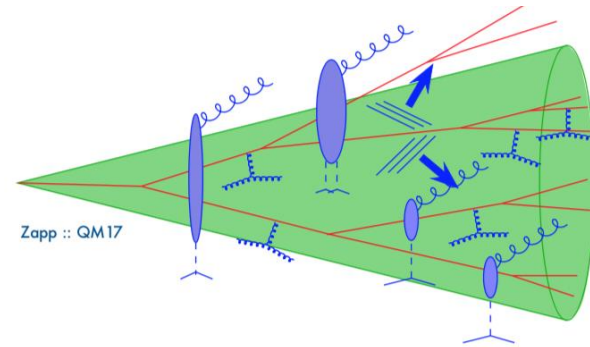
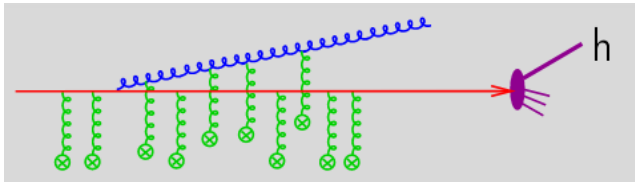
Depletion then excess at intermediate to high z :

Modification of fragmentation pattern due to quenching?

Preferential quenching of gluon jets over quark jets?

Jet quenching state-of-the-art

Single hard parton traversing QCD matter is very much a 1st approximation

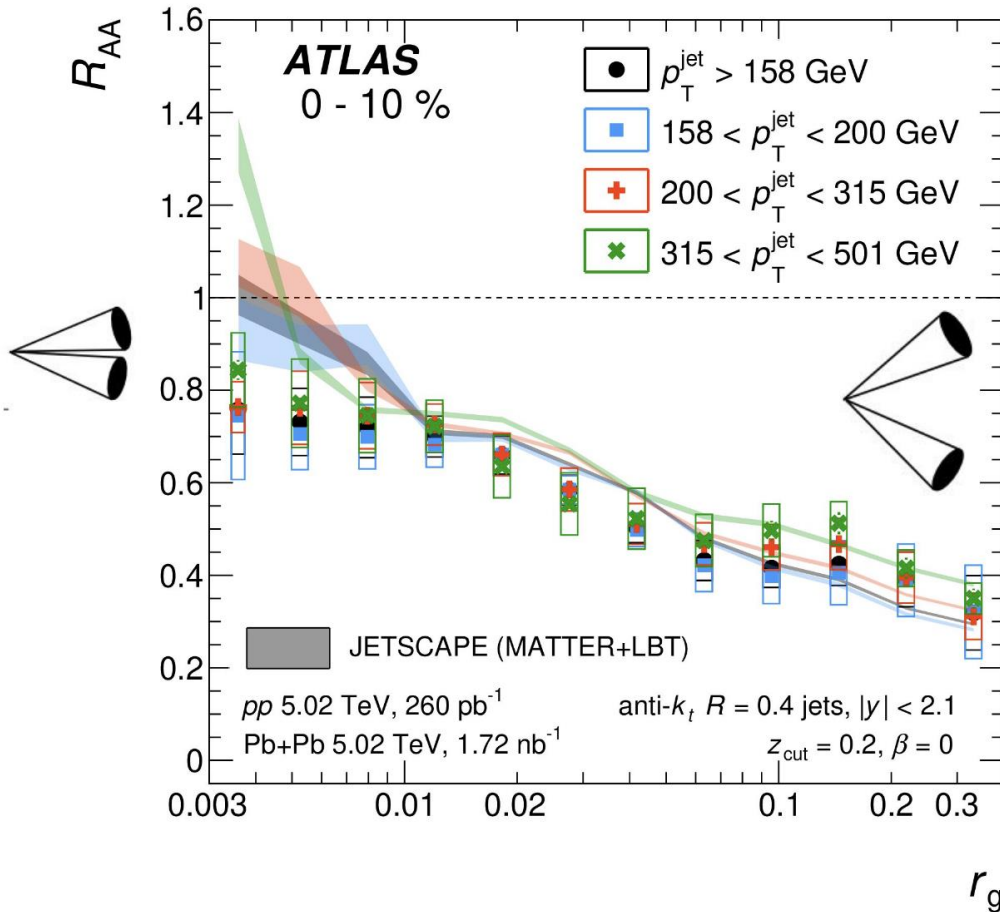


More realistic treatment requires taking into account the entire parton shower

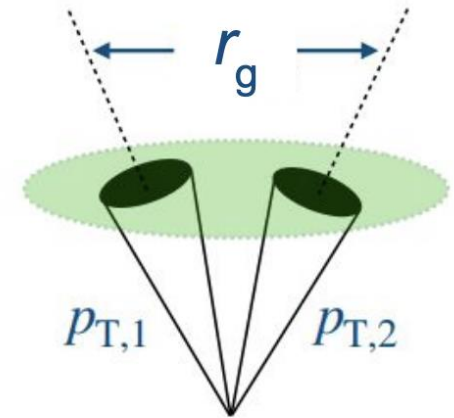
Among the effects currently being studied

- Coherence effects between nearby partons in the shower & jet substructure
- Modeling of recoil “splash”
- Development of full jet quenching MC codes

New: Jet substructure

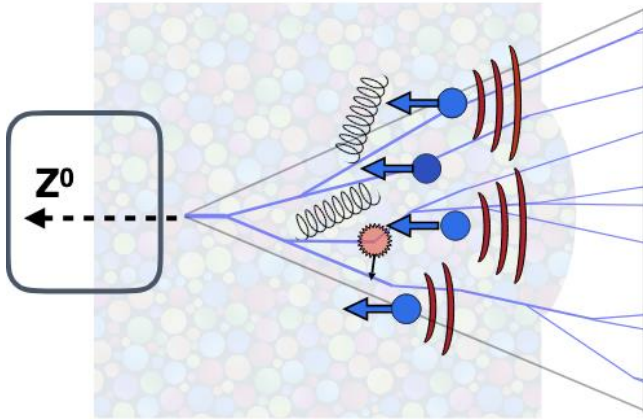
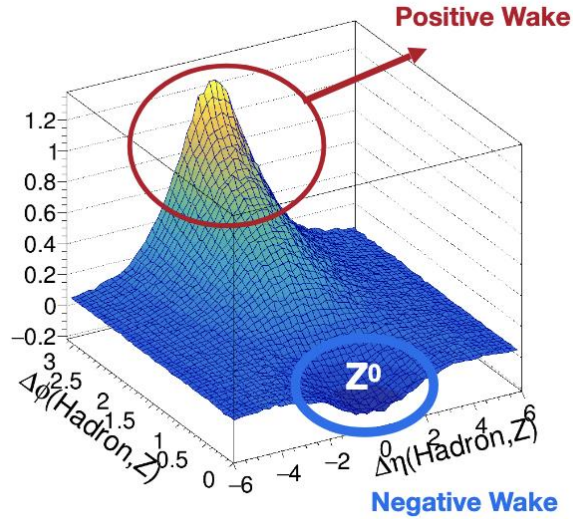


“Decluster” your jets to look at the hardest prongs

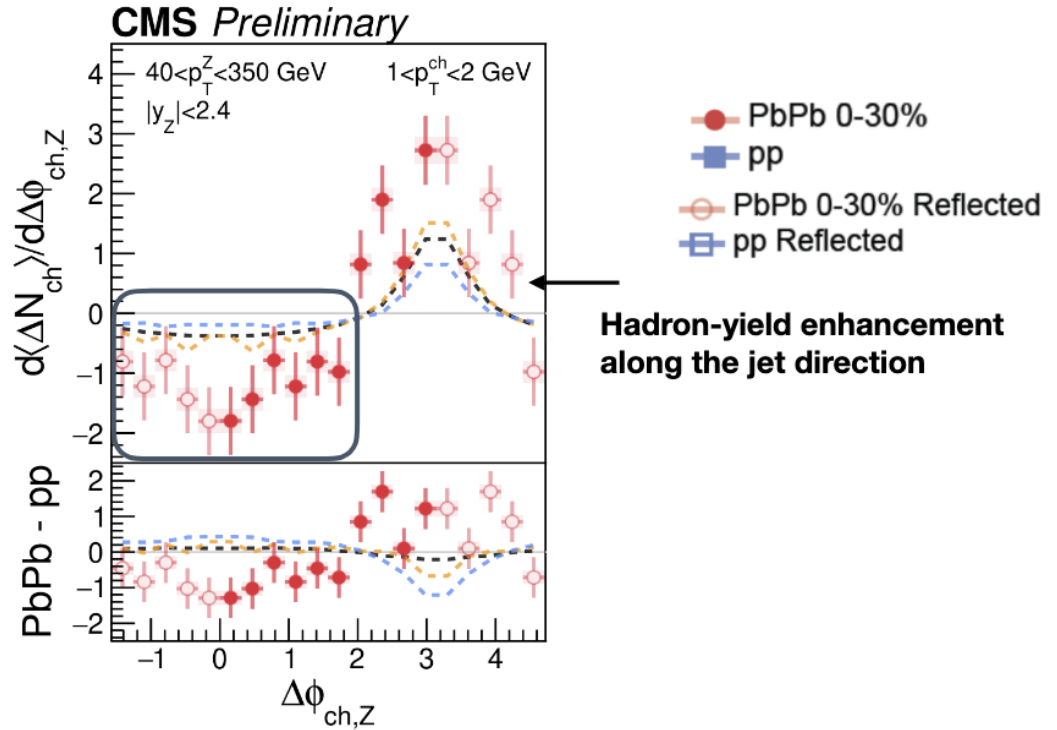


Larger suppression for more separated prongs
 → Coherence effects?

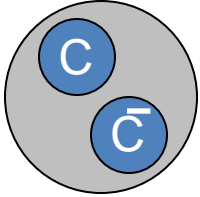
New: jet wake



Clear depletion in PbPb
on the Z side ($\Delta\phi=0$)



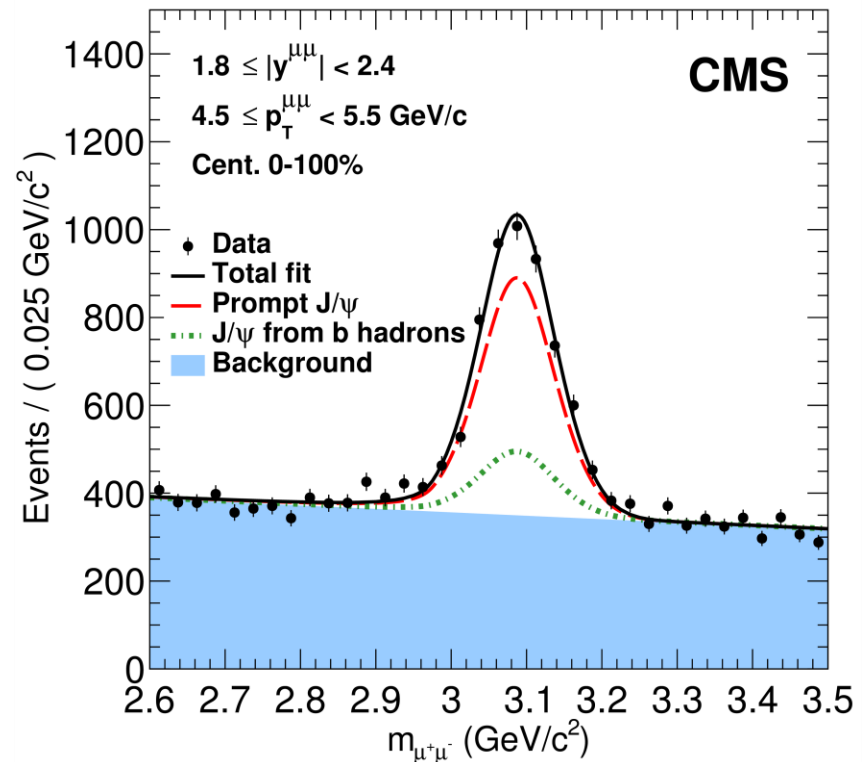
J/ψ meson factsheet



J/ψ mass peak in PbPb collisions,
via the dimuon decay channel

- Bound state of c-cbar, most abundant quarkonium state
- 1974 “November revolution”, confirmed 4th quark generation
- Large rate of decays in leptons, 6% into dimuon or dielectron
- Mass = 3.1 GeV (~ 2x charm)

PbPb 368 μb⁻¹ (5.02 TeV)



Charm quark above temperate threshold for thermal production
→ produced via hard scattering

Size of the J/ψ

Cornell potential:

$$V(r) = \underbrace{\sigma r}_{\text{Confinement}} - \underbrace{\frac{\alpha}{r}}_{\text{Coulomb-like}}$$

σ and α from spectrum of onium states

$$\sigma \approx 1 \text{ GeV/fm}, \alpha \approx \pi/12$$

$$E(r) = \underbrace{2m}_{\text{Rest}} + \underbrace{\frac{1}{2mr^2}}_{\text{kinetic}} + V(r)$$

Minimizing to get the lowest bound state

$$\frac{dE}{dr} = \frac{1}{mr^3} - \frac{\alpha}{r^2} - \sigma = 0$$

Dominant term is confining one

$$r \approx (m\sigma)^{-1/3}$$

Inserting $m = 1.5 \text{ GeV}$
 $r \approx 0.3 \text{ fm}$

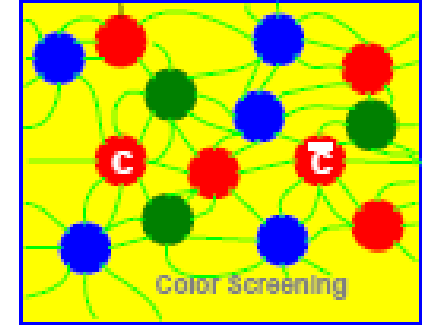
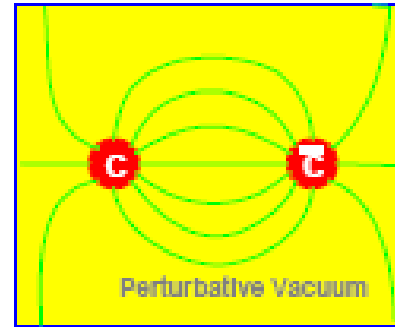
→ Several times smaller
 than typical light hadrons

Quarkonia melting

Potential is screened in the QGP

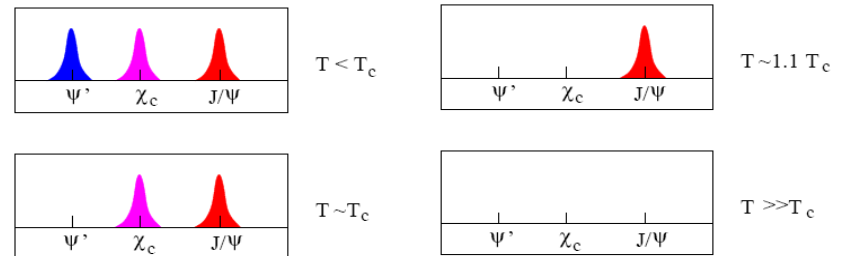
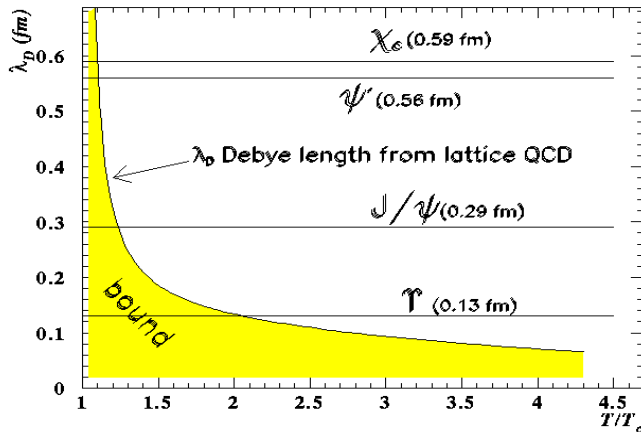
$$V(r) = -\frac{\alpha}{r} e^{-r/\lambda_D}$$

Debye screening length



Besides J/ψ , also excited states ψ' & χ_c
 Feed-down: 10% of J/ψ from ψ'
 30% of J/ψ from χ_c

λ_D provided by lattice QCD



Sequential melting of states
 according to binding energy

Charmonia @ the SPS

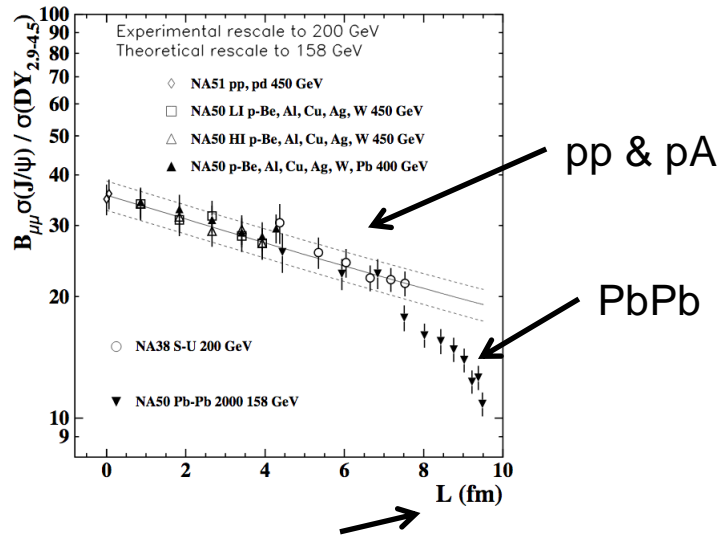
SPS collided Pb ions in fixed target mode at CERN

200 GeV Pb beam

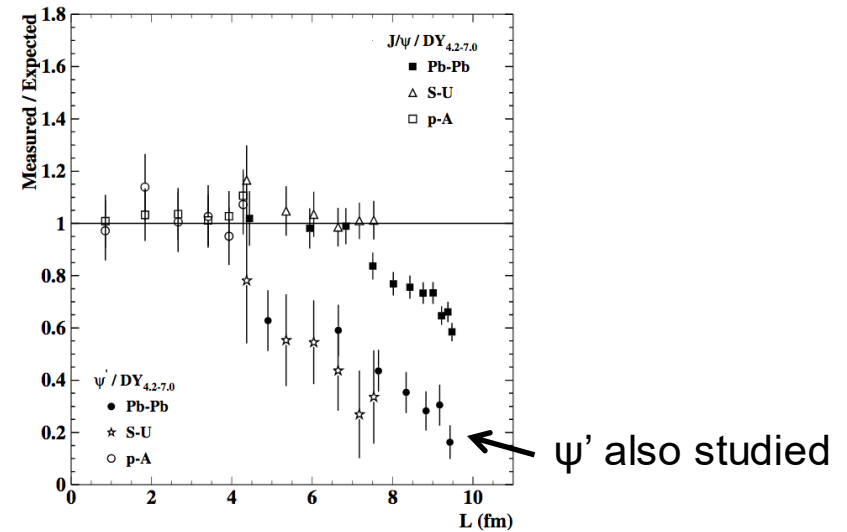
$\rightarrow \sqrt{s_{NN}} \approx 17$ GeV

Instead of R_{AA} , ratio to Drell-Yan

Ratio to nuclear absorption baseline

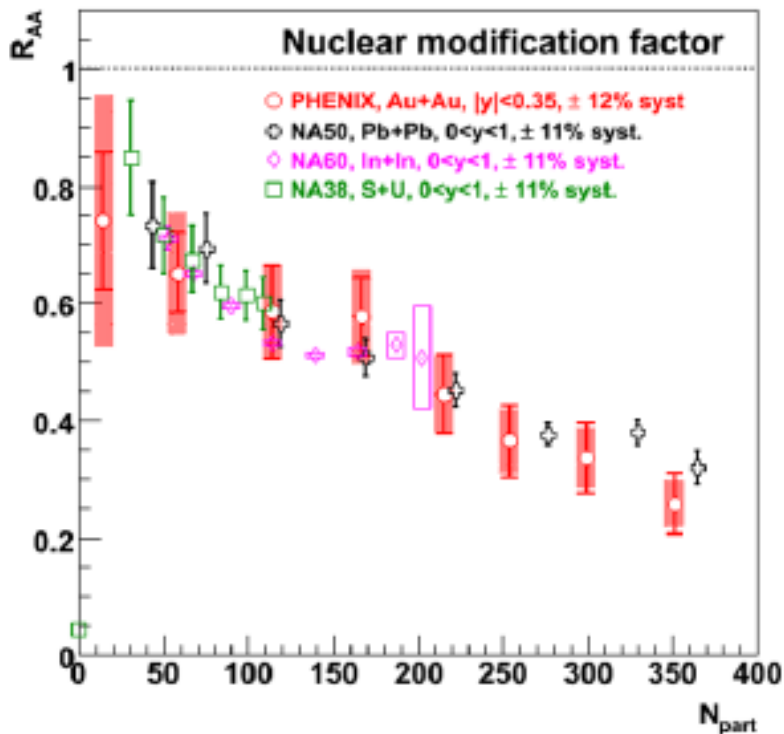


Mean nuclear length traversed (L),
from Glauber model

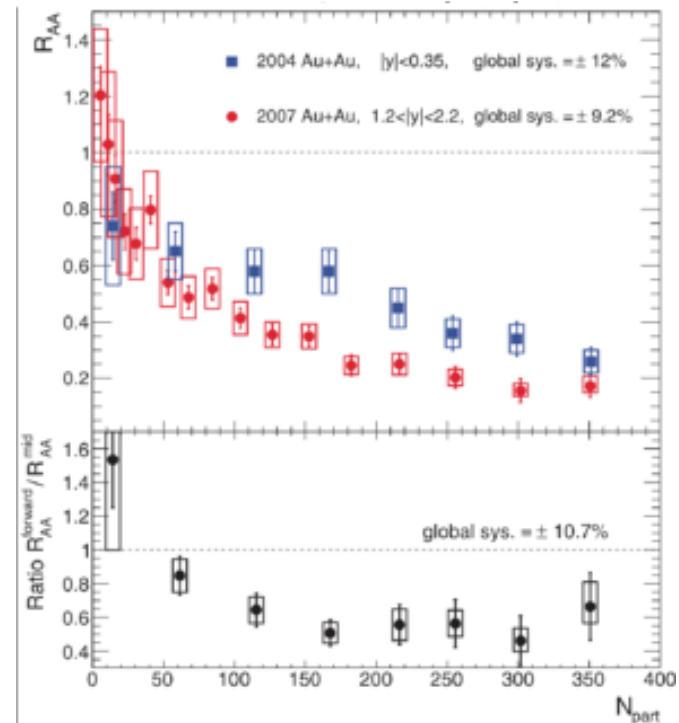


J/ψ suppressed even in light systems due to (pre-resonant) nuclear absorption
 However, in central PbPb ‘anomalous’ suppression observed
 Naïve interpretation: ψ' melts first, then χ_C , while J/ψ remains intact at this energy
 Evidence for QGP discovery, although modern interpretation more nuanced

Charmonia @ RHIC $\sqrt{s_{NN}} = 200$ GeV



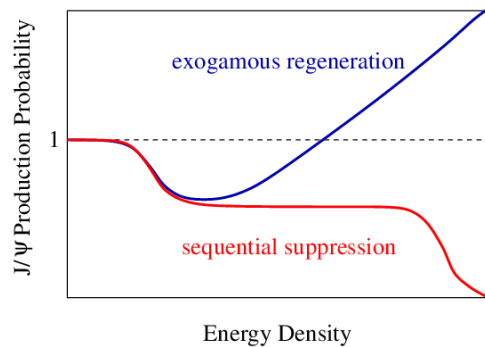
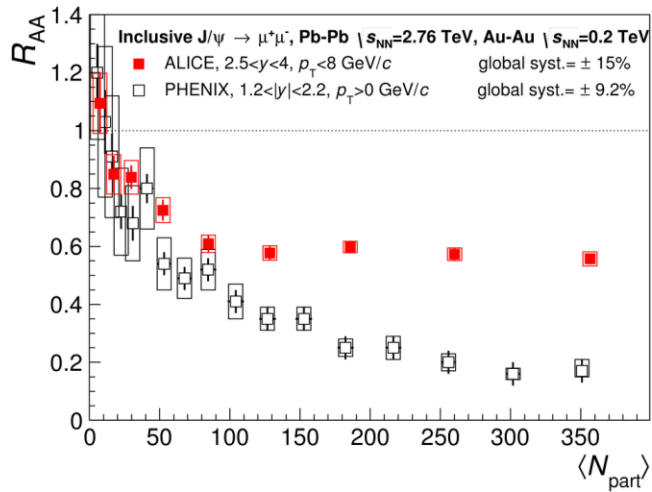
Comparable suppression to SPS
 \rightarrow Saturation of onia melting?



However, larger suppression at forward rapidity \rightarrow not expected in Debye screening picture

Melting & regeneration

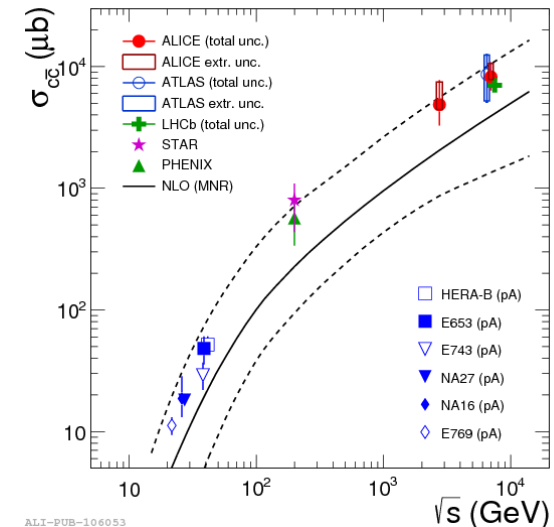
Less suppression at LHC



Interpretation: Cancellation of two effects, melting and *regeneration*

LHC: $\sqrt{s_{NN}} = 2760$ GeV

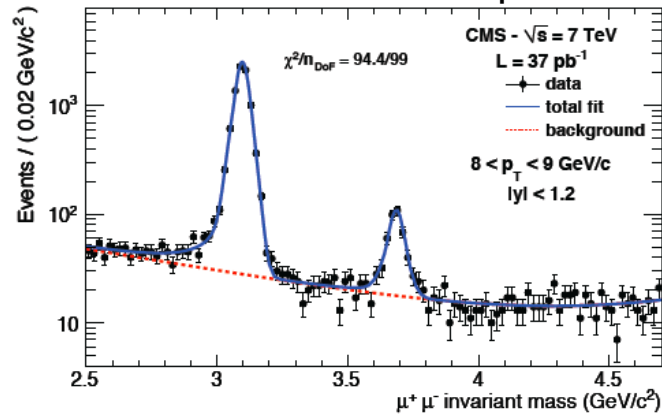
Charm copious enough at high energies to meet other charm in the QGP. Effect $\sim N_c^2$



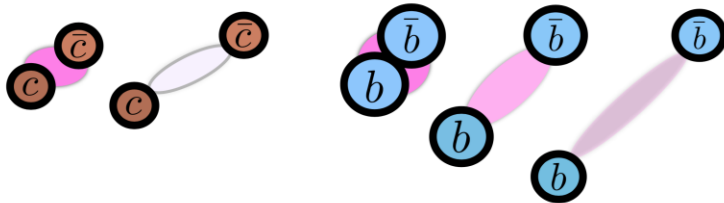
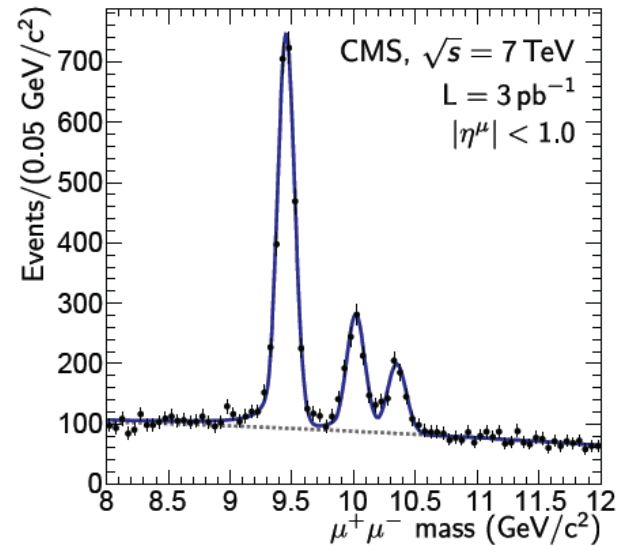
$\sigma(cc) \approx 5$ mb x 1500 coll /
65 mb ≈ 115 cc pairs in a
central PbPb collision!

Upsilon spectroscopy

Charmonium dimuon spectrum



Bottomonium dimuon spectrum

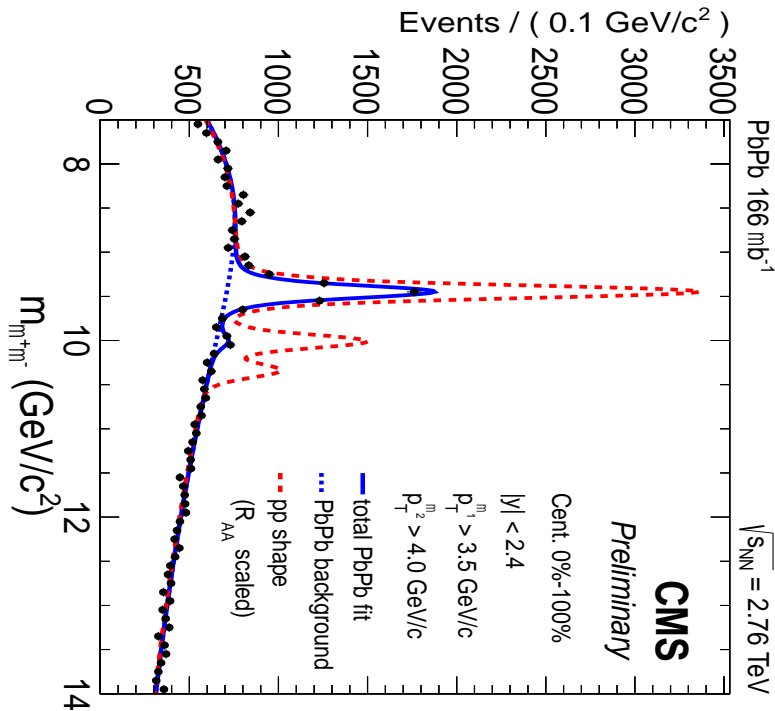


state	J/ψ	ψ'	Υ	Υ'	Υ''
mass (GeV)	3.10	3.68	9.46	10.0	10.4
radius (fm)	0.25	0.45	0.14	0.28	0.39

- 3 Υ states w/ comparable cross section, but different binding energy
- Υ(1s) more tightly bound than J/ψ

bb cross section ≈ 0.3 mb \rightarrow little regeneration expected!

Upsilon melting



$$R_{AA}(Y(1S)) = 0.453 \pm 0.014 \pm 0.046;$$

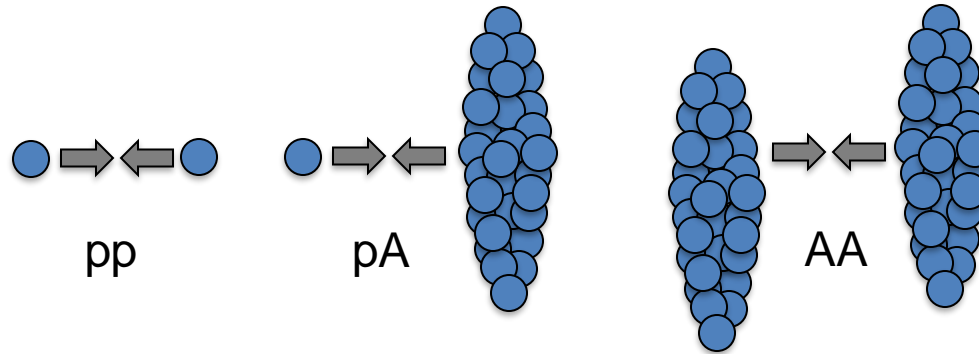
$$R_{AA}(Y(2S)) = 0.119 \pm 0.028 \pm 0.015;$$

$$R_{AA}(Y(3S)) < 0.145 \text{ at a 95\% confidence level,}$$

- Y(1s) suppressed, although $\approx 30\%$ feed-down from excited states
- Y(2s) more suppressed
- Y(3s) not even visible

- Naïve interpretation: Excited states melt, while ground state survives
- More likely different states dissociate over different volumes
- Also need to consider the role of *cold nuclear effects*

Addendum: small systems



Original paradigm

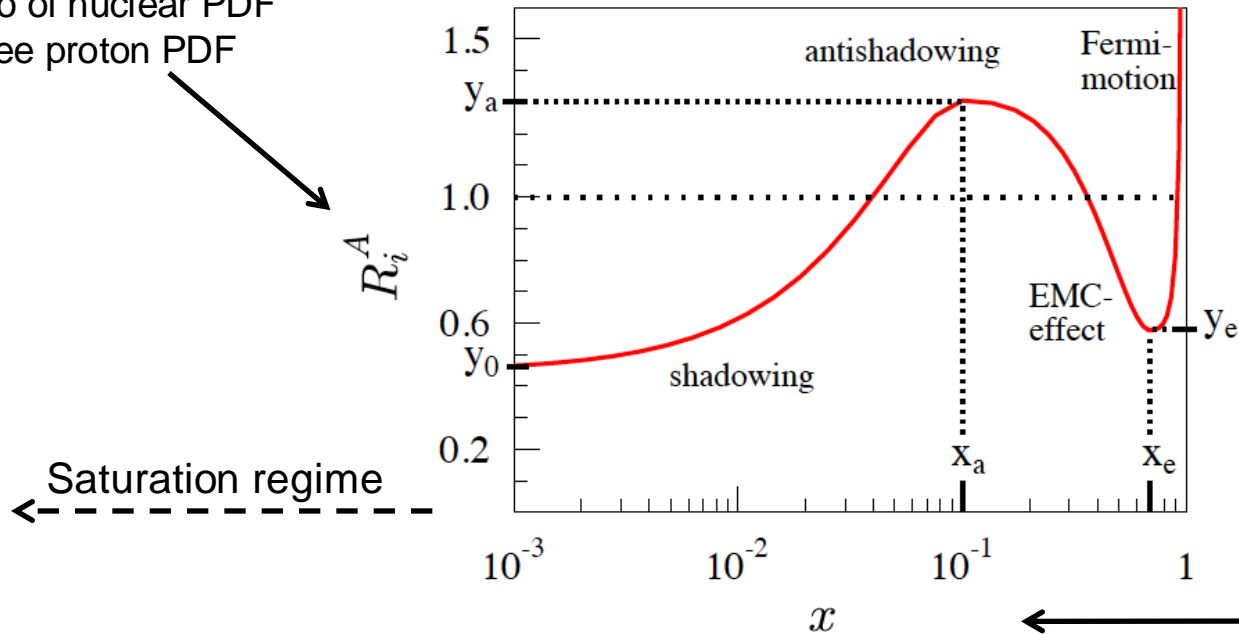
- pp collisions a baseline for absence of nuclear effects
- pA collisions a baseline for cold nuclear matter effects, i.e., nuclear effects unrelated to the quark-gluon plasma

Paradigm shift over last ~ 15 years

Certain effects thought to be tied to deconfinement show up even in smaller systems

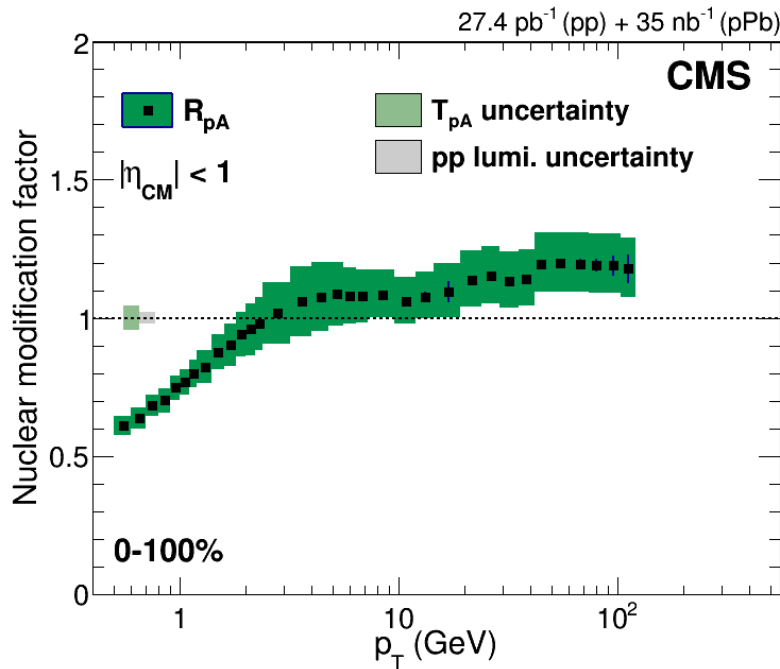
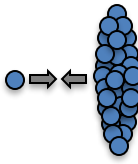
Nuclear PDFs

Ratio of nuclear PDF
to free proton PDF

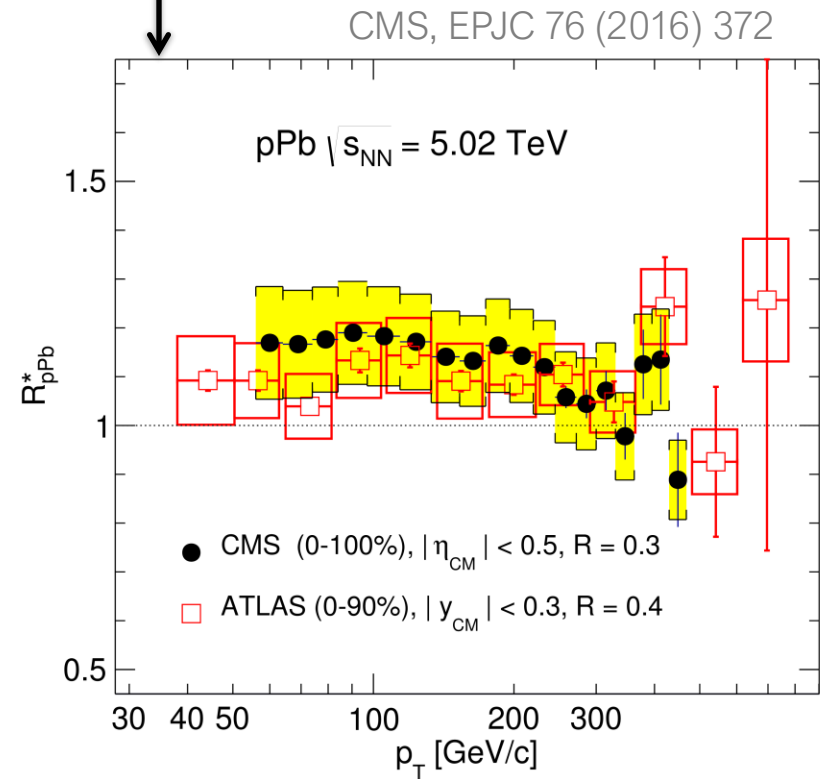


- Modest nuclear effect already expected in pA: modification of parton distributions inside nuclei
- Known from nuclear DIS measurements
- Evolution with Q^2 calculable in QCD via DGLAP equations
- Gluons must saturate at very low $x \rightarrow$ non-linear evolution

Hadron & jet R_{pA}



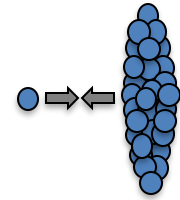
CMS, JHEP 04 (2017) 039



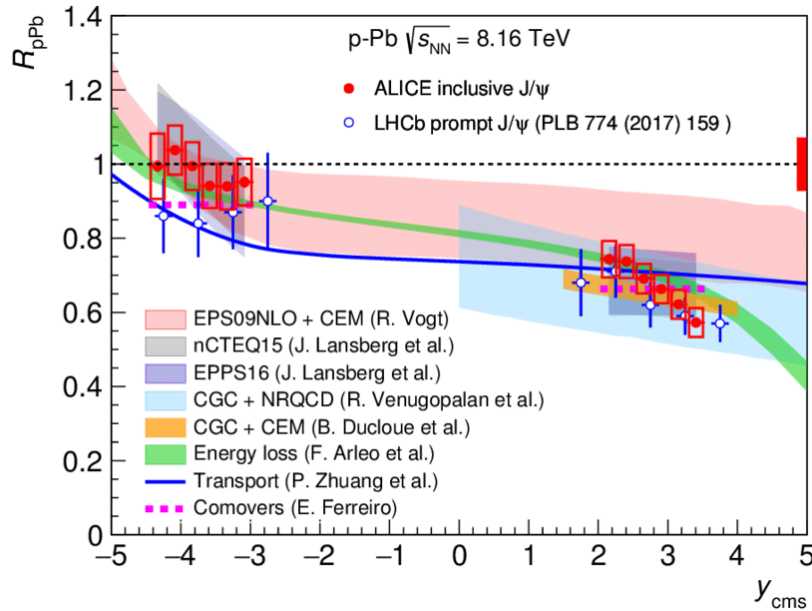
- No suppression as would be expected if jet quenching were present
- Rather, small enhancement due to nuclear parton distributions (anti-shadowing)

Quarkonia R_{pA}

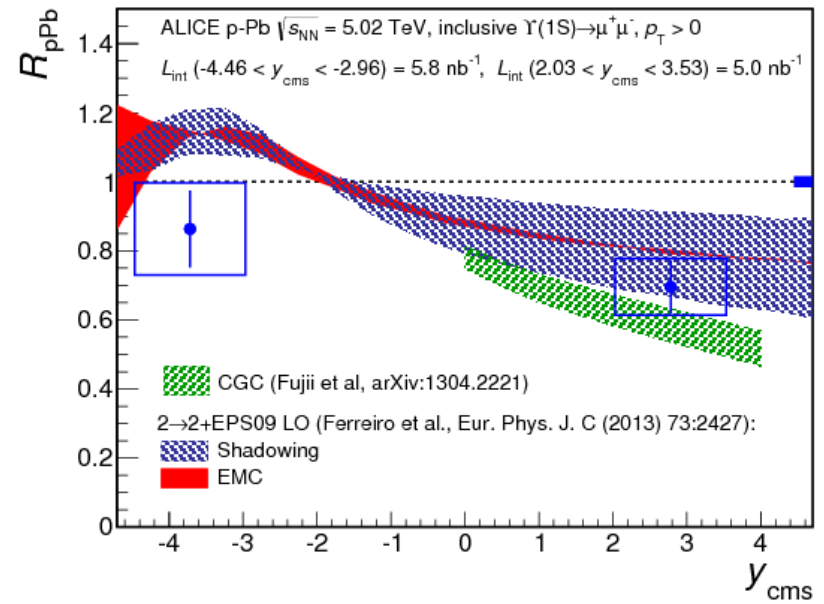
Proton-going direction \longrightarrow



J/ψ



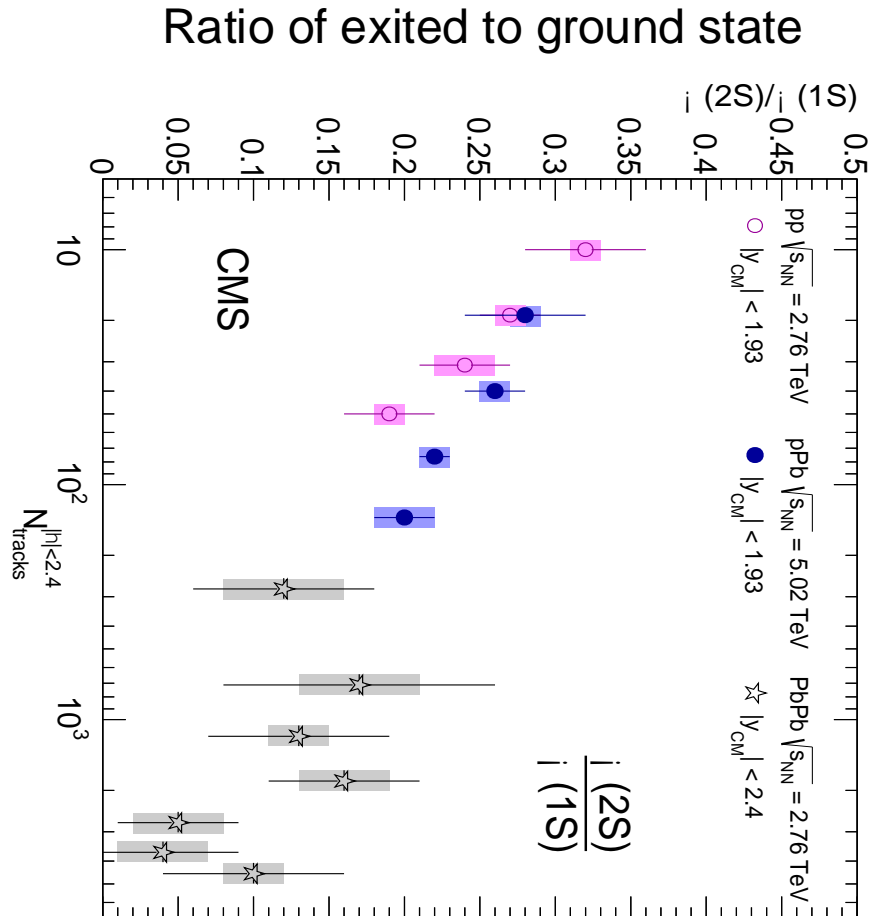
$\Upsilon(1s)$



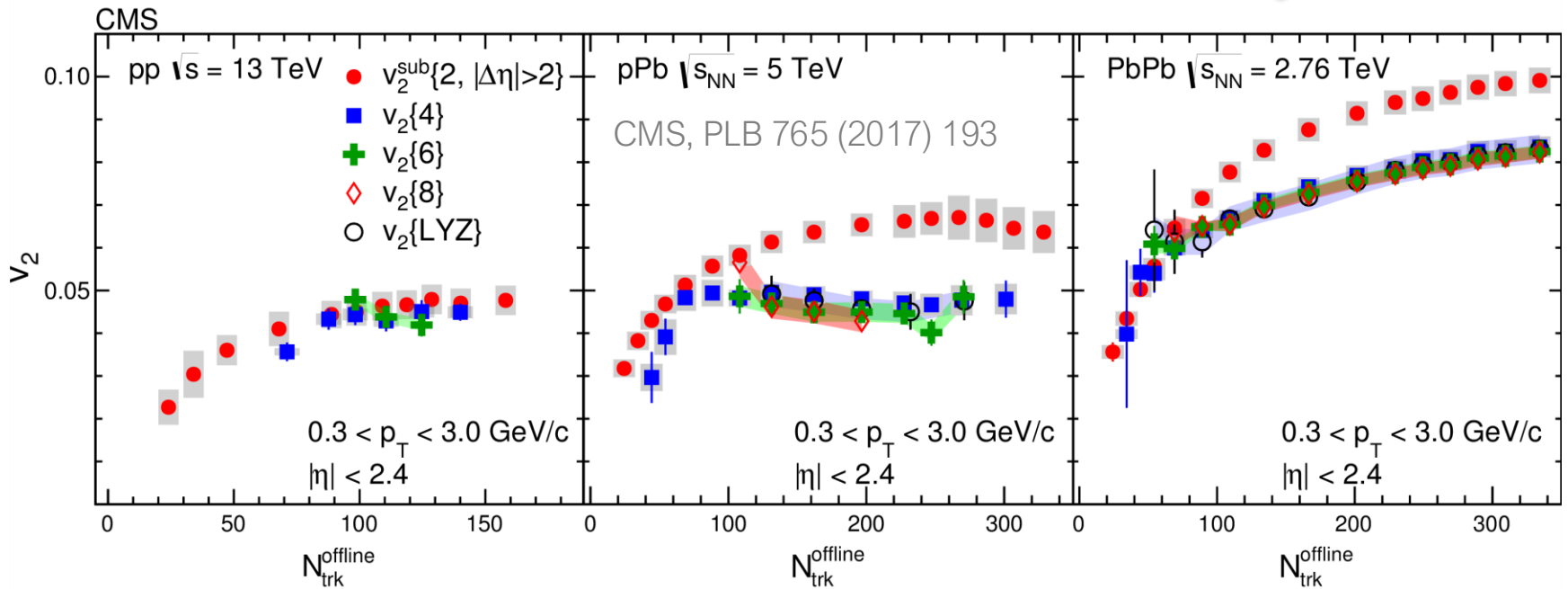
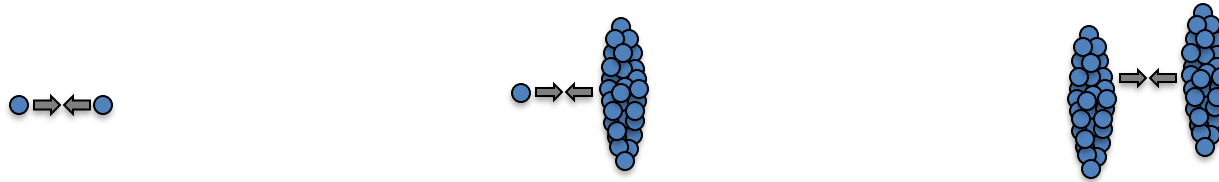
Quarkonia access lower x , where shadowing expected to dominate
 Results consistent with shadowing in proton-going direction
 Part of the Υ suppression is due to PDF effects

Y states vs. multiplicity

- Centrality difficult to define in small systems → multiplicity as a proxy
- Preferential suppression of $Y(2s)$ already seen in pp and pPb collisions
- Consistent with a smooth dependence on multiplicity



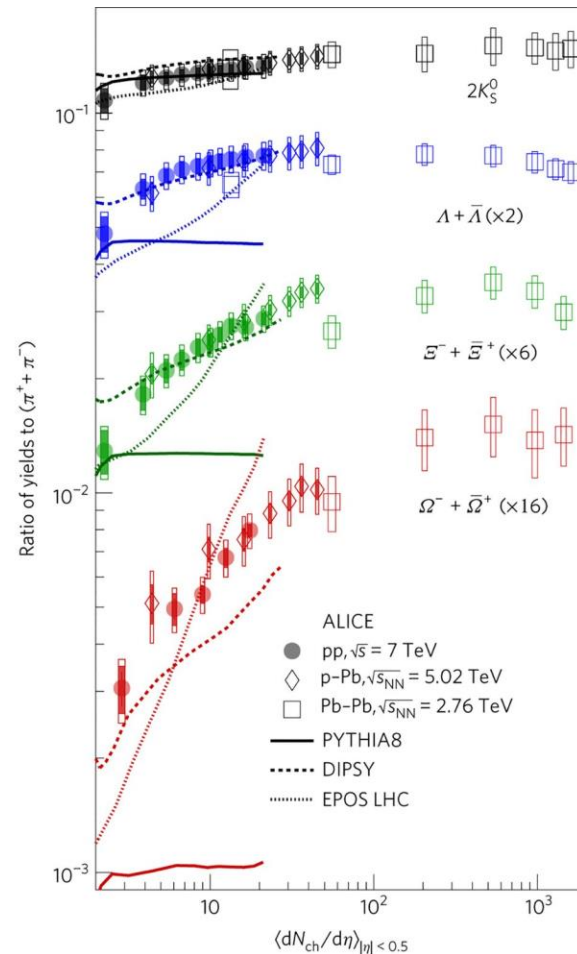
Multiparticle correlations



- No initial asymmetry in small systems, yet elliptic flow observed for mult > 100
- Source suspected to be fluctuations at the sub-nucleonic scale
- Hydrodynamic behavior at relatively small multiplicity is surprising

Strangeness in small systems

- Recall strangeness enhancement indicative of chemical equilibrium over an extended volume
- Effect also present in small systems
- Smooth dependence on multiplicity
- Increases w/ # of strange quarks



Small systems: summary

- Nuclear modification factor (R_{pA}) consistent with no QGP, rather effects from modified nuclear PDFs
- However, with increasing multiplicity, effects thought to relate to QGP formation emerge in pp and pA collisions
 - Preferential suppression of excited $Y(2s)$ over $Y(1s)$
 - Flow signals from initial state fluctuations
 - Strangeness enhancement indicating chemical equilibration
- May indicate a QGP droplet is formed in small systems, but very much an ongoing field of study

Part 2, take-home messages

- The QGP is studied with a variety of probes, a few of which have been discussed here
 - Parton energy loss in the QGP via jet quenching
 - Debye screen in the QGP via onia dissociation
 - Collectivity of the QGP via elliptic flow
- Emergence of some of these effects in small systems is an ongoing avenue of investigation