

1 Towards a Pixel TPC part I: construction and test of a
2 32-chip GridPix detector

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9 **Abstract**

10 A Time Projection Chamber (TPC) module with 32 GridPix chips was con-
11 structed and the performance was measured using data taken in a testbeam
12 at DESY in 2021. The GridPix chips each consist of a Timepix3 ASIC
13 (TPX3) with an integrated amplification grid and have a high efficiency of
14 about 85% to detect single ionisation electrons. In the testbeam setup, the
15 module was placed in between two sets of Mimosas26 silicon detector planes
16 that provided external high precision tracking and the whole detector setup
17 was slid into the PCMAG magnet at DESY. The TPC could be operated
18 reliably and used a 93.6/5.0/1.4 gas mixture (by volume) of Ar/iC₄H₁₀/CO₂
19 with a small amount of oxygen and water vapour. The analysed data were
20 taken at electron beam momenta of 5 and 6 GeV/c and at magnetic fields of
21 0 and 1 T.

22 The result for the transverse diffusion coefficient D_T is $(287 \pm 0.5) \mu\text{m}/\sqrt{\text{cm}}$
23 at $B = 0$ T and D_T $(121 \pm 0.5) \mu\text{m}/\sqrt{\text{cm}}$ at $B = 1$ T. The longitudinal dif-
24 fusion coefficient D_L is measured to be $(251 \pm 14) \mu\text{m}/\sqrt{\text{cm}}$ at $B = 0$ T

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25 and $(224 \pm 14) \mu\text{m}/\sqrt{\text{cm}}$ at $B = 1 \text{ T}$. Results for the tracking systematical
26 uncertainties in xy (pixel plane) were measured to be smaller than $13 \mu\text{m}$
27 with and without magnetic field. The tracking systematical uncertainties in
28 z (drift direction) were smaller than $15 \mu\text{m}$ ($B = 0 \text{ T}$) and $20 \mu\text{m}$ ($B = 1$
29 T).

30 *Keywords:*

31 Micromegas, gaseous pixel detector, micro-pattern gaseous detector,
32 Timepix, GridPix, pixel time projection chamber

33 1. Introduction

34 Earlier publications on a single chip [1] and four chip (quad) GridPix de-
35 tectors [2] showed the potential of the GridPix technology and the large range
36 of applications for these devices [3]. In particular, it was demonstrated that
37 single ionisation electrons can be detected with high efficiency and accuracy,
38 allowing excellent 3D track position measurements and particle identification
39 based on the number of electrons and clusters.

40 As a next step towards a Pixel Time Projection Chamber for a future
41 collider experiment [4], [5], a module consisting of 32 GridPix chips based on
42 the TPX3 chip was constructed.

43 A GridPix detector consists of a CMOS pixel TPX3 chip [6] with inte-
44 grated amplification grid added by photo-lithographic - Micro-electromechanical
45 Systems (MEMS) - post-processing techniques. The TPX3 chip can be op-
46 erated with a low threshold of $515 e^-$, and has a low equivalent noise charge
47 of about $70 e^-$. The GridPix single chip and quad detectors have a very fine
48 granularity of $55 \times 55 \mu\text{m}^2$ and a high efficiency of about 85% - demonstrated

49 in this paper - to detect single ionisation electrons.

50 Based on the experience gained with these detectors a 32 GridPix detector
51 module - consisting of 8 quad detectors - was built. A drift box defining the
52 electric field and gas envelop was constructed. A readout system for up to
53 128 chips with 4 multiplexers readout by one Speedy Pixel Detector Readout
54 (SPIDR) board [7] [8] was designed. After a series of tests using the laser
55 setup [9] and cosmics in the laboratory at Nikhef, the detector was taken to
56 DESY for a two week testbeam campaign.

57 At DESY, the 32-chip detector was placed in between two sets of Mi-
58 mosa26 silicon detector planes and mounted on a movable stage. The whole
59 detector setup was slid into the centre of the PCMAG magnet at DESY. A
60 beam trigger was provided by scintillator counters. The data reported here
61 were taken at different stage positions and electron beam momenta of 5 and
62 6 GeV/c and at magnetic fields of 0 and 1 T. The performance of the 32
63 GridPix detector module was measured using these data sets.

64 In this paper, part I of the results will be presented with the main focus
65 on the detector spatial resolution and tracking performance. A second follow
66 up paper will discuss the dE/dx (or dN/dx) and other results.

67 **2. The 32-GridPix detector module**

68 A 32 GridPix detector module was built using the quad detector module
69 [2] as a basic building block. The quad module consists of four GridPix
70 chips and is optimised for a high fraction of sensitive area of 68.9%. The
71 external dimensions are 39.60 mm \times 28.38 mm. The four chips which are
72 mounted on a cooled base plate (COCA), are connected with wire bonds

73 to a common central 6 mm wide PCB. A 10 mm wide guard electrode is
74 placed over the wire bonds 1.1 mm above the aluminium grids, in order to
75 prevent field distortions of the electric drift field. The guard electrode is the
76 main inactive area, and its dimensions are set by the space required for the
77 wire bonds. On the back side of the quad module, the PCB is connected to
78 a low voltage regulator. The aluminium grids of the GridPix detectors are
79 connected by 80 μm insulated copper wires to a high voltage (HV) filtering
80 board. The quad module consumes about 8W of power of which 2W is used
81 in the LV regulator.

82 Eight quad modules were embedded in a box, resulting in a GridPix
83 detector module with a total of 32 chips. A schematic 3-dimensional drawing
84 of the detector is shown in Figure 1. A schematic drawing of the quad
85 detectors in the module is shown in Figure 2, where also the beam direction
86 is indicated.

87 The internal dimensions of the box are 79 mm along the x -axis, 192 mm
88 along the y -axis, and 53 mm along the z -axis (drift direction), and it has a
89 maximum drift length (distance between cathode and readout anode) of 40
90 mm. The drift field is shaped by a series of parallel CuBe field wires of 75
91 μm diameter with a wire pitch of 2 mm. Guard strips are located on all of
92 the four sides of the active area. In addition, six guard wires - shown with
93 dashed lines (one colored red) in Figure 2 - are suspended over the boundaries
94 of the chips, to minimise distortions of the electric drift field. The wires are
95 located at a distance of 1.15 mm from the grid planes, and their potential is
96 set to the drift potential at this drift distance. The box has two 50 μm thick
97 Kapton windows to allow the beam to pass with minimal multiple scattering.

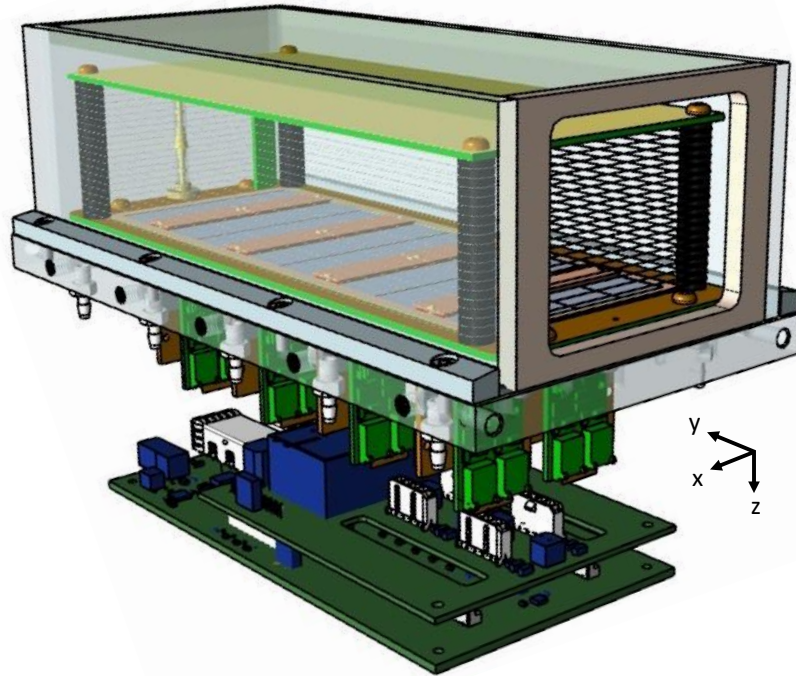


Figure 1: Schematic 3-dimensional rendering of the 32-GridPix module detector for illustration purposes.

98 The gas volume of 780 ml is continuously flushed at a rate of ~ 50 ml/min
 99 (about 4 volumes/hour) with premixed T2K TPC gas. This gas is a mixture
 100 consisting of 95% Ar, 3% CF_4 , and 2% iC_4H_{10} suitable for large TPCs because
 101 of the low transverse diffusion in a magnetic field and the high drift velocity.

102 The data acquisition system of the quad module was adopted to allow for
 103 reading out multiple quad detectors. A multiplexer card was developed that
 104 handles four quad detectors or 16 chips and combines the TPX3 data into
 105 one data stream. For the 32 GridPix module two multiplexers are connected
 106 to a SPIDR board that controls the chips and readout process. The readout
 107 speed per chip is 160 Mbps and for the multiplexer 2.56 Gbps this corre-

108 sponds to a maximum rate of 21MHits/s. For each pixel the precise Time of
 109 Arrival (ToA) using a 640 MHz TDC and the time over threshold (ToT) are
 110 measured.

111 3. Experimental setup

112 In preparation of the two weeks DESY testbeam campaign, a support
 113 frame was designed to move the 32-chip GridPix detector module in the
 114 plane perpendicular to the beam by a remotely controlled stage such that
 115 the whole detector volume could be probed. The module was mounted upside
 116 down with respect to figure 1 to allow access to the electronics from above.
 117 The support frame also held three Mimosa26 silicon detector planes [10] -

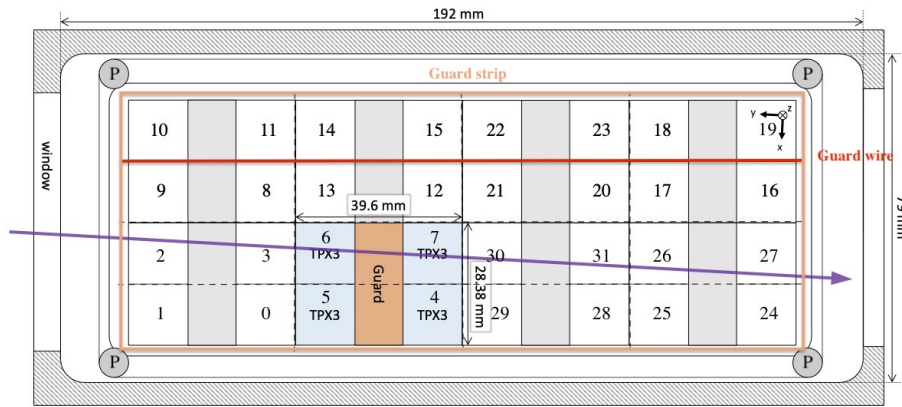


Figure 2: Schematic drawing of the 32-GridPix module detector with one example quad as viewed from the top of the quad detectors. The chips are numbered and the beam direction is shown in purple. A guard electrode of a quad detector is shown in orange. The four surrounding guard strips are shown -not to scale- in orange. Six guard wires - are shown with dashed lines (one colored red) and the pillars of the drift box are shown a circles with a P in the centre.

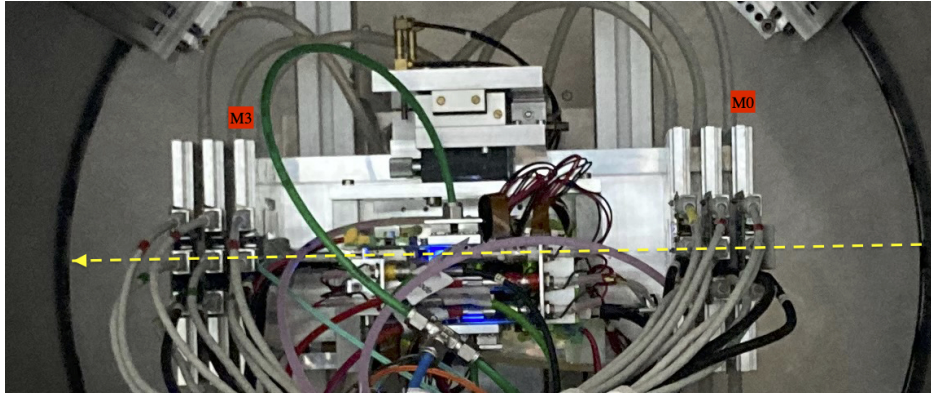


Figure 3: Photo of the detector setup - side view - at the centre of the PCMAG magnet (the circular contour). The Mimosa26 planes M0 and M3 are indicated in red as well as the beam direction (yellow).

118 with an active area of $(21.2 \text{ mm} \times 10.6 \text{ mm})$ - placed in front of the detector
119 and three Mimosa26 planes behind the detector. At DESY, the Mimosa26
120 silicon detector planes were provided by the testbeam coordinators. The
121 whole detector setup was slid towards the centre of the PCMAG magnet
122 at the DESY II testbeam facility [10]. A beam trigger was provided by a
123 double scintillator counter coincidence. The data were taken at different
124 stage positions to cover the whole sensitive TPC volume. Runs with electron
125 beam momenta of 5 and 6 GeV/c and at magnetic fields of 0 and 1 T were
126 analysed.

127 A photograph of the detector setup in the PCMAG magnet is shown in
128 Figure 3. The stage positions of the TPC module with respect to the beam
129 and the Mimosa26 planes can be adjusted.

130 The experimental and environmental parameters such as temperature,
131 pressure, gas flow and oxygen content were measured and logged by a Win-

132 dows operated slow control system. The experimental parameters are sum-
 133 marised in Table 1. The chips were cooled by circulating Glycol through
 134 the cooling channels in the module carrier plate. The cooling blocks of the
 135 multiplexers were further cooled by blowing pressurised air on them.

Table 1: Overview of the experimental parameters. The ranges indicate the variation over the data taking period

Number of analysed runs at $B=0$ (1) T	6 (8)
Run duration	10-90 minutes
Number of triggers per run	3-100 k
E_{drift}	280 V/cm
V_{grid}	340 V
Threshold	550 e ⁻
Gas temperature	303.3-306.6 K
Pressure	1011 – 1023 mbar
Oxygen concentration	240 - 620 ppm
Water vapour concentration	2000 - 7000 ppm

136 The data was produced in four main data streams: one stream produced
 137 by the Mimosa26 telescope, two data streams by the two Timepix multiplex-
 138 ers and one trigger stream. The double scintillator coincidence provided a
 139 trigger signal to the Trigger Logic Unit (TLU) [11] that sends a signal to the
 140 telescope readout and the trigger SPIDR. The data acquisition systems of
 141 the telescope and trigger SPIDR injected a time stamp into their respective
 142 data streams. Hits from the Mimosa26 planes were collected with a sliding
 143 window of $-115 \mu\text{s}$ to $230 \mu\text{s}$ around the trigger time. The data acquisition
 144 of the multiplexer and the trigger SPIDR were synchronised at the start of

145 the run. By comparing the time stamps in these streams, telescope tracks
146 and TPC tracks could be matched. Unfortunately, the SPIDR trigger had
147 - due to a cabling mistake at the output of the TLU - a common 25ns flat
148 time jitter.

149 After a short data taking period one of the chips (nr 11) developed a
150 short circuit and the HV on the grid of the chip was disconnected. After the
151 testbeam data taking period the module was repaired in the clean room in
152 Bonn.

153 4. Analysis

154 4.1. Telescope track reconstruction procedure

155 The data of the telescope is decoded and analysed using the Corryvreckan
156 software package [12]. The track model used for fitting was the General
157 Broken Lines (GBL) software [14]. The code was extended and optimised to
158 fit curved broken lines for the data with magnetic field. The telescope planes
159 were iteratively aligned using the standard alignment software provided by
160 the package. The single point Mimosas26 resolution is 4 μm in x and 6 μm
161 in z (drift direction) [10].

162 Telescope tracks were selected were required to have hits in at least 5 out
163 of the 6 planes and a total χ^2 of better than 25 per degree of freedom. The
164 uncertainties on the telescope track prediction in the middle of the GridPix
165 detector module are dominated by multiple scattering. The amount of mul-
166 tiple scattering was estimated by comparing the predictions from the two
167 telescope arms for 6 GeV/c tracks at $B = 0$ T. The expected uncertainty in
168 x and z is 26 μm on average.

169 *4.2. TPC Track reconstruction procedure*

170 GridPx hits are selected requiring a minimum time over threshold ToT
171 of 0.15 μs . The drift time is defined as the measured time of arrival minus
172 the trigger time recorded in the trigger SPIDR data stream minus a fixed t_0
173 (the drift time at zero drift). The drift time was corrected for time walk [2]
174 using the measured time over threshold (ToT in units of μs) and the formula
175 (1):

$$\delta t = \frac{18.6(ns \mu s)}{\text{ToT} + 0.1577(\mu s)}. \quad (1)$$

176 Furthermore, small time shift corrections - with an odd-even and a $16 \times$ pixels
177 structure - coming from the TPX3 clock distribution were extracted from the
178 data and applied.

179 The z drift coordinate was calculated as the product of the drift time and
180 the drift velocity. This implies that $z_{\text{drift}} = -z$ as defined in figure 1. GridPix
181 hits outside an acceptance window of 30 mm wide in x and 15 mm wide in
182 z were not used in the track finding and reconstruction. Based on a Hough
183 transform an estimate of the TPC track position and angles in the middle of
184 the module (at $y = 1436$ pixels) were obtained. This estimate was used to
185 collect the hits around the TPC track and fit the track parameters. For this
186 fit a linear (for $B = 0$ T data) or a quadratic track (for $B = 1$ T data) model
187 was used. In the fit, the expected uncertainties per hit $\sigma_x y$ and σ_z were
188 used. The expected uncertainties were derived using the parametrisations
189 discussed in section 5. The fit was iterated three times to reject outlier hits
190 at respectively 10, 5 and 2.5 sigma. A TPC track was required to have a
191 least 100 hits in each multiplexer. At least 25% of the total number of hits

Table 2: Table with track/event selection cuts

Track/Event Selection

$$|x_{\text{TPC}} - x_{\text{telescope}}| < 0.3 \text{ mm}$$

$$|z_{\text{TPC}} - z_{\text{telescope}}| < 2 \text{ mm}$$

$$|dx/dy_{\text{TPC}} - dx/dy_{\text{telescope}}| < 4 \text{ mrad}$$

$$|dz/dy_{\text{TPC}} - dz/dy_{\text{telescope}}| < 2 \text{ mrad}$$

192 should be on track and the χ^2 per degree of freedom had to be less than 3 in
 193 xy and zy . All track parameters were expressed at a plane in the middle of
 194 the TPC module.

195 The calibration and alignment of the detector was done using high quality
 196 tracks for which the track selections are summarised in table 2.

197 The drift velocity was calibrated per run by fitting a linear function to
 198 the z (predicted from the telescope track at the measured TPC hit position)
 199 versus the measured drift time in the TPC. For the $B = 0$ T runs it varies
 200 between 61.6 and 63.0 $\mu\text{m}/\text{ns}$. For the $B = 1$ T runs it is between 57.2 and
 201 59.1 $\mu\text{m}/\text{ns}$. The variation comes mainly from the changes in the relative
 202 humidity of the gas volume due to small leaks.

203 The individual TPX3 chips were iteratively aligned fitting a shift in x
 204 (z drift) and two slopes $dx(z \text{ drift})/drow(\text{column})$. The alignment was
 205 done per run, because the detector was moved in x and/or z for each run.
 206 The fitted slopes were also corrected for small shifts and rotations (3D) in
 207 the nominal chip position.

208 An example event from run 6916 without B field with a TPC and a
 209 telescope track is shown in figure 4. The TPC is located between $y = 0$ and

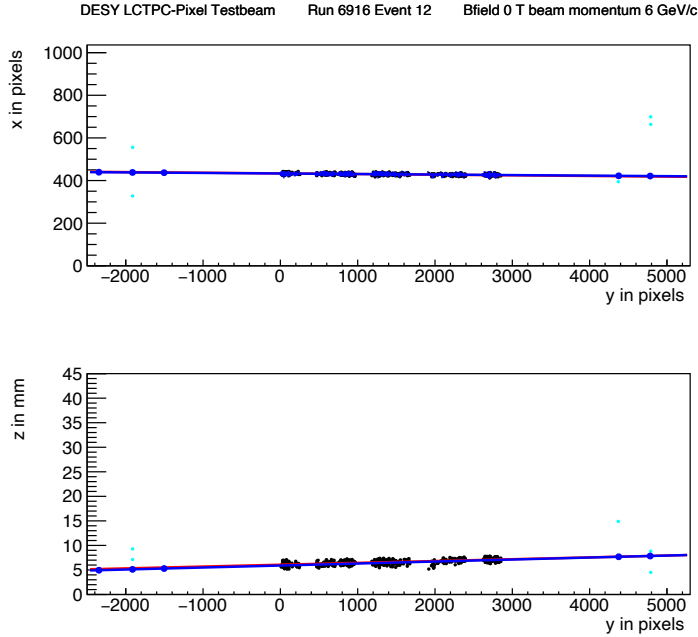


Figure 4: An event display for run 6916 without B field, with in total 1293 TPC hits (black dots) in the precision plane (x, y) and drift plane $(z \text{ drift}, y)$. The fitted TPC track (red line) with 1130 hits on track and the telescope track (blue line) with 5 Mimosa26 planes (blue hits) on track are shown. In green the off track Mimosa26 hits are shown.

210 2872 pixels. Three Mimosa26 planes are located at $y < -1000$ and three at
 211 $y > 4000$ pixels.

212 5. Hit resolutions

213 The single electron hit resolutions in xy and z will be extracted from
 214 the residuals with respect to the fitted track. The track residual in xy is
 215 the closest point of the track in the xy plane to the hit at the center of the
 216 pixel. The residual in z is calculated at this point of closest approach. In
 217 order to study the single electron resolution for the data with and without

218 magnetic field, additional selections on the telescope and TPC tracks were
 219 applied. Firstly, due to the trigger time jitter of 25 ns (corresponding to 1.5
 220 mm drift), the prediction of the telescope track in z must be used as the
 221 reference for z . Secondly, the z hits of the TPC track were fitted to correct
 222 for the common time shift and the z residuals were calculated with respect to
 223 the fitted TPC track. In the xy plane the residuals of TPC hits with respect
 224 to the telescope track were used to extract the single electron resolution in
 225 xy . For the resolution studies runs at three different z stage positions of the
 226 TPC were selected where the beam gave hits in the central chips. The data
 227 of 14 central chips (9, 12, 21, 20, 17, 16, 2, 3, 6, 7, 30, 31, 26 and 27) was
 228 used. Two chips (8 and 13) were left out because of the E field deformations
 229 caused by the short circuit in chip 11.

230 *5.1. Hit resolutions in the pixel plane*

231 The residual of the hits in the pixel plane (xy) was measured as a function
 232 of the predicted drift position (z_{drift}). Only hits are used crossing the fiducial
 233 region defined by the central core of the beam and staying 20 pixels away in
 234 x from the chip edges. The spread on the residual in xy for an ionisation
 235 electron is given by:

$$\sigma_{xy}^2 = \sigma_{\text{track}}^2 + \frac{d_{\text{pixel}}^2}{12} + D_T^2(z_{\text{drift}} - z_0), \quad (2)$$

236 where σ_{track} is the uncertainty from the track prediction, d_{pixel} is the pixel
 237 pitch size, z_0 is the position of the grid, and D_T is the transverse diffusion
 238 coefficient. The last two terms correspond to the single electron detector
 239 resolution (squared). The resolution at zero drift distance $d_{\text{pixel}}/\sqrt{12}$ was
 240 fixed to 15.9 μm and σ_{track} to 30 μm for $B = 0$ T and 42 μm for $B =$

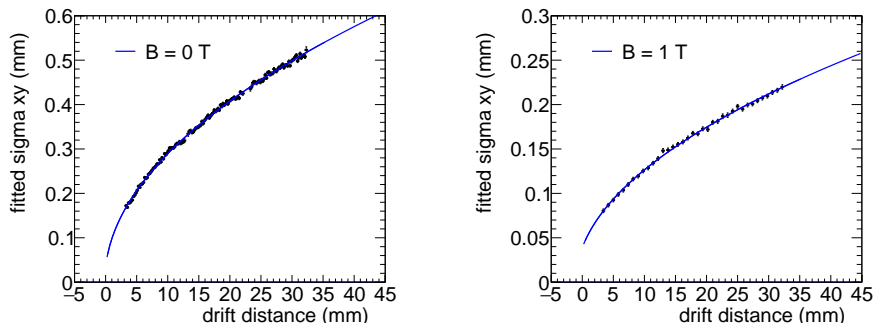


Figure 5: Measured spread on the residuals in the pixel plane (black points) fitted with equation (2) (blue line).

241 1 T data. The uncertainty on the track prediction was measured and is
 242 larger than the Mimosa plane resolution because of multiple scattering in
 243 the sensors and in the entrance and exit windows.

244 The expression (2) - leaving z_0 and D_T as free parameters - is fitted to
 245 the $B = 0$ T data shown in Figure 5. The fit gives a transverse diffusion
 246 coefficient D_T of $(287 \pm 0.5) \mu\text{m}/\sqrt{\text{cm}}$. The measured value is in agreement
 247 with the value of $287 \mu\text{m}/\sqrt{\text{cm}} \pm 4\%$ predicted by the gas simulation software
 248 Magboltz 11.9 [15]. The values of the diffusion coefficients depend on the
 249 humidity that was not precisely measured during the testbeam. The humidity
 250 strongly affects the drift velocity. Therefore the drift velocity prediction from
 251 Magboltz was used to determine the water content per run and predictions
 252 for the diffusion coefficients could be obtained.

253 A fit to the $B = 1$ T data, also shown in Figure 5, gives a transverse
 254 diffusion coefficient D_T of $(121 \pm 0.5) \mu\text{m}/\sqrt{\text{cm}}$. The measured value is in
 255 agreement with the value of $119 \mu\text{m}/\sqrt{\text{cm}} \pm 2\%$ predicted by Magboltz.

256 *5.2. Hit resolution in the drift plane*

257 The spread on the residuals in z of the ionisation electrons σ_z is given by:

$$\sigma_z^2 = \sigma_{\text{track}}^2 + \sigma_{z_0}^2 + D_L^2(z_{\text{drift}} - z_0), \quad (3)$$

258 where σ_{track} is the expected track uncertainty, σ_{z_0} the detector resolution at
 259 zero drift distance and D_L the longitudinal diffusion constant. The last two
 260 terms in the equation correspond to the single electron detector resolution
 261 (squared). Only tracks crossing the fiducial region were accepted and hits
 262 with a ToT value above $0.6 \mu\text{s}$ were selected. Because of the time jitter,
 263 the fitted TPC track is used for the drift residuals. For z_{drift} the telescope
 264 prediction at the hit was used. The expected uncertainty on TPC track
 265 prediction is propagated and amounts to $50 \mu\text{m}$ at $z = z_0$. The systematic
 266 uncertainty on σ_{track} is estimated to be $25 \mu\text{m}$.

267 The expression (3) - leaving σ_{z_0} and D_L as free parameters - is fitted
 268 to the $B = 0 \text{ T}$ data shown in Figure 6. The value of z_0 was fixed to the
 269 result of the fit in the xy plane. The value of σ_{z_0} was measured to be 129
 270 μm . The longitudinal diffusion coefficient D_L was determined to be $(251$
 271 $\pm 1 \text{ (stat)} \pm 14 \text{ (sys)}) \mu\text{m}/\sqrt{\text{cm}}$, which is higher than the expected value
 272 $236 \pm 3 \mu\text{m}/\sqrt{\text{cm}}$ from a Magboltz calculation [15]. The quoted systematic
 273 uncertainty on D_L is rather large and obtained from a fit using $\sigma_{\text{track}} = 25$
 274 μm .

275 A fit to the $B = 1 \text{ T}$ data shown in Figure 6 gives a longitudinal diffusion
 276 coefficient D_L of $(224 \pm 2 \text{ (stat)} \pm 14 \text{ (sys)}) \mu\text{m}/\sqrt{\text{cm}}$. The measured value
 277 is lower than the value of $(245 \pm 4) \mu\text{m}/\sqrt{\text{cm}}$ predicted by Magboltz. The
 278 fitted value of σ_{z_0} was $114 \mu\text{m}$.

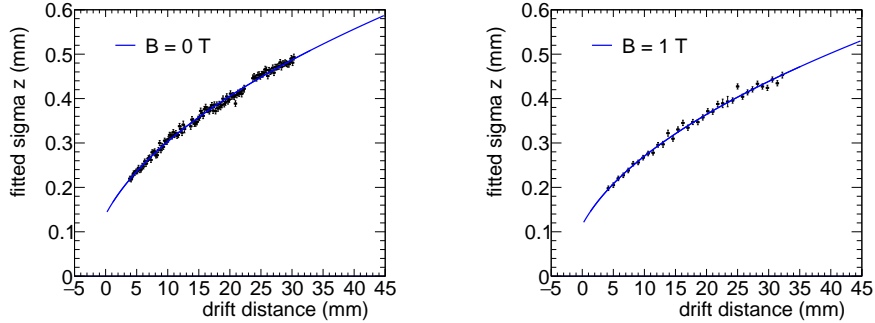


Figure 6: Measured spread on the residuals in the drift plane for hits with a ToT above $0.60 \mu\text{s}$. The data are fitted with the expression of equation (3).

279 5.3. Deformations in the pixel and drift plane

280 It is important to measure possible deformations in the pixel (xy) and
 281 drift (z) plane to quantify the tracking precision. For the construction of
 282 a large Pixel TPC, deformations in the pixel plane deformation should be
 283 controlled to better than typically $20 \mu\text{m}$ because these affect the momentum
 284 resolution. The mean residuals in the pixel and drift planes are shown in
 285 Figure 7 for the $B = 0 \text{ T}$ data set using a large set of runs to cover the whole
 286 module. The residuals were calculated with respect to the telescope track
 287 prediction. Because of limited statistics, bins were grouped into 8×16 pixels.
 288 Bins with less than 100 hits are left out and residuals larger (smaller) than
 289 $+(-)100 \mu\text{m}$ are shown in red (blue).

290 A few critical areas can be observed in figure 7: the region around chip 11
 291 is affected (chips 14, 8 and 13), because the grid of chip 11 was disconnected.
 292 Deformations are present at the four corners of the drift box (chips 1, 10, 19
 293 and 24) and close to the upper corner edge (chip 16) of the drift box. These
 294 come from inhomogenities in the drift field near the supporting pillars, the

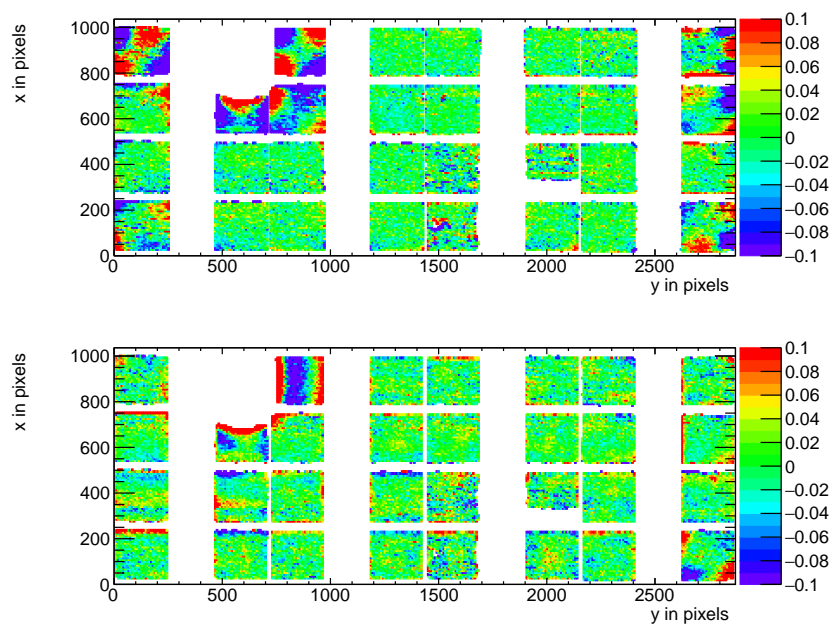


Figure 7: Mean residuals (color coded in mm) in the pixel (top) and drift (bottom) plane for $B = 0$ T data at the expected hit position.

295 field wires are too close to the chip to provide a constant electric field. It
296 was concluded that for the deformation studies the hits of these nine chips
297 have to be removed. The track fit was redone leaving these hits out of the fit,
298 such that they could not bias and affect the results. Note that a bias in the
299 mean residual at the edge of the chips is expected to be present for an ideal
300 detector because of the finite coverage and the diffusion in the drift process.

301 In order to reduce the statistical fluctuations and quantify the tracking
302 precision, the module was regrouped in $(4 \times 256) \times 256$ pixel planes put side
303 by side on the horizontal axis, as shown in figure 8. E.g. the selected chips
304 from the upper left and bottom left quad detectors are combined into the
305 0-256 (x) and 0-256 (y) plane. Bins have a size of 16×16 pixels and bins
306 with less than 1000 entries are not shown. Due to the presence of the dike,
307 pixels at the edge of the chip became covered and inefficient. Therefore, the
308 region of 5 pixels in y near the edge of the chip was removed. For the drift
309 coordinate studies, a region of 10 pixels near the edge of the chip in x and
310 y was removed. The total number of measurements (bins) in xy is 895 and
311 in z 892. One can observe that in the module plane no clear systematic
312 deviations are present and conclude that the guard wire voltages were on
313 average well tuned. Note that in the quad detector module we had no guard
314 wires and deformation corrections had to be applied [2]. The r.m.s. of the
315 distribution of the measured mean residual over the surface in the pixel plane
316 is $11 \mu\text{m}$ and in the drift plane $15 \mu\text{m}$. Similarly, regrouping the module in
317 $256 \times (4 \times 256)$ pixels put them side by side on the vertical axis, yielded a
318 r.m.s. in the pixel plane of $13 \mu\text{m}$ and $13 \mu\text{m}$ in the drift coordinate. The
319 expected statistical error - obtained by propagating the uncertainties on the

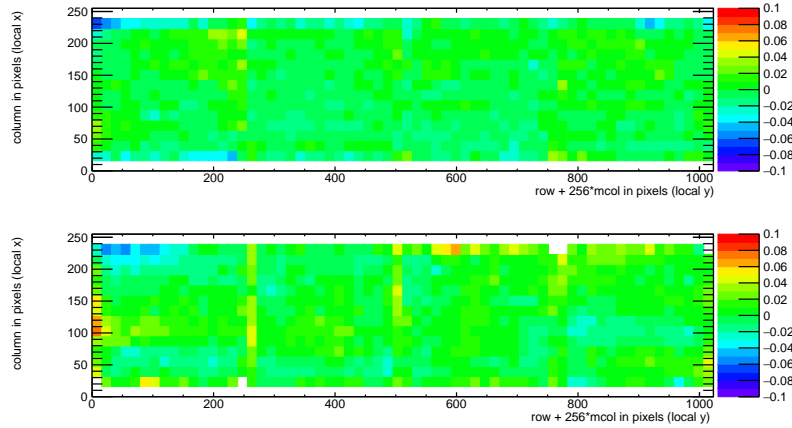


Figure 8: Mean residuals (color coded in mm) in the pixel (top) and drift plane (bottom) for $B = 0$ T data at the expected hit position.

320 residuals - in xy is $4 \mu\text{m}$ and in z $5 \mu\text{m}$.

321 In the $B = 1$ T data set, the electrons will drift mainly along the magnetic
 322 field lines. Deformations are in that case due to e.g. the non-alignment of the
 323 electric and magnetic field, giving $E \times B$ effects. Unfortunately, the statistics
 324 of the telescope tracks that have a matched TPC track was insufficient and
 325 did not cover the full TPC module plane. Therefore the larger statistics of
 326 matched and unmatched TPC tracks was used. TPC tracks were required
 327 to pass angular selection cuts (dx/dy between -40 and -20 mrad and dz/dy
 328 between 0 and 14 mrad) and a momentum cut ($p > 2$ GeV/c and $q < 0$).

329 The mean residuals in the pixel and drift planes are shown in figure 9 for
 330 the $B = 1$ T data set using a large set of runs to cover the whole module. The
 331 (biased) residuals were calculated with respect to the TPC track prediction.
 332 Because of limited statistics bins were grouped into 16×16 pixels. Bins with
 333 less than 100 hits are left out and residuals larger (smaller) than $+(-)100 \mu\text{m}$

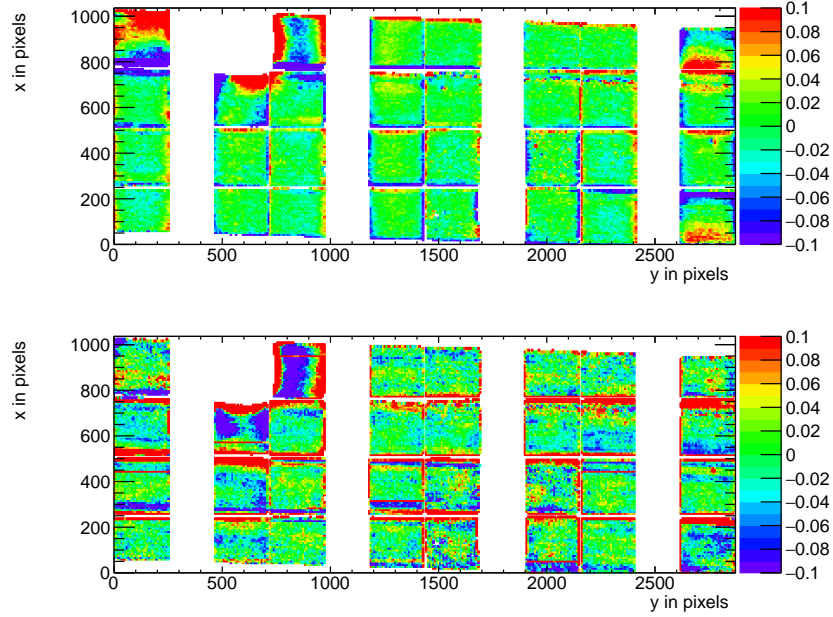


Figure 9: Mean residuals (color coded in mm) in the pixel and drift plane for $B = 1$ T data at the expected hit position.

334 are shown in red (blue).

335 In figure 9 the critical areas discussed above - around chip 11, the four
 336 corner chips and chip 16 in the upper corner edge - can be clearly observed.
 337 For the deformation studies the hits of these nine chips have to be removed.
 338 The TPC track fit was redone leaving these hits out of the fit, thus that
 339 they could not bias and affect the results. The TPC plane is well covered,
 340 although one can observe that due to the angle of the beam in the xy plane
 341 the chips in the upper right and lower left corners are not fully covered.

342 In order to reduce the statistical fluctuations and quantify the tracking
 343 precision, the module was regrouped in $(4 \times 256) \times 256$ pixels as described
 344 above, as shown in figure 10. Bins have a size of 16×16 pixels and bins with

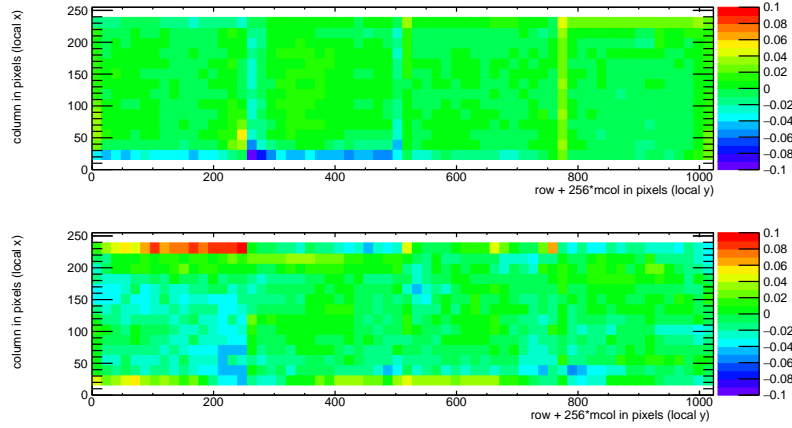


Figure 10: Mean residuals (color coded in mm) in the pixel and drift plane for $B = 1\text{T}$ data at the expected hit position.

345 less than 1000 entries are not shown. Similar to the no-field deformations
 346 studies, acceptance cuts had to be applied. The region of 16 pixels in y
 347 near the edge of the chips was removed. For the drift coordinate studies, in
 348 addition a region of 10 pixels in x near the edge of the chip was removed.
 349 The total number of measurements (bins) in xy is 896 and in z 896. One can
 350 observe that in the module plane no clear systematic deviations are present.
 351 The r.m.s. of the distribution of the measured mean residual over the surface
 352 in the pixel plane is $13\ \mu\text{m}$ and in the drift plane $19\ \mu\text{m}$. Similarly, regrouping
 353 the module in $256 \times (4 \times 256)$, yielded a r.m.s. in the pixel plane of $11\ \mu\text{m}$ and
 354 $20\ \mu\text{m}$ in the drift coordinate. The expected statistical error in xy is $2\ \mu\text{m}$
 355 and in z $3\ \mu\text{m}$.

356 In summary, the deformations studies for the $B = 0$ and $1\ \text{T}$ data demon-
 357 strate that the systematical uncertainties in xy are smaller than $13\ \mu\text{m}$ with
 358 and without magnetic field. The systematical uncertainties in z were smaller

359 than $15 \mu\text{m}$ ($B = 0 \text{ T}$) and $20 \mu\text{m}$ ($B = 1 \text{ T}$).

360 *5.4. Tracking resolution*

361 A selected TPC track in the $B = 0 \text{ T}$ data has on average 1000 hits. The
362 tracking precision in the middle of the TPC (at $y = 1436$ pixels) was derived
363 on a track-by-track basis, by propagating the pixel TPC hit uncertainties. It
364 was found to be on average $9 \mu\text{m}$ in the precision plane and $13 \mu\text{m}$ in z . The
365 angular resolution in dx/dy was on average 0.19 mrad and for dz/dy 0.25
366 mrad . It is clear that the position resolution in the TPC in the precision
367 and drift coordinates is impressive for a track length of (only) 158 mm .
368 The values are smaller than the uncertainty on the track prediction from
369 the silicon telescope of $26 \mu\text{m}$ in x and z on average that is dominated by
370 multiple scattering.

371 **6. Single electron efficiency**

372 The distribution of the number of TPC track hits per chip - without
373 requiring a matched telescope track - are shown in figure 11 for the data
374 without magnetic field and for the $B = 1 \text{ T}$ data. For the $B = 0 \text{ T}$ data the
375 central chips 2,6,7,9,16,17,26 and 27 were selected. For the $B = 1 \text{ T}$ data
376 the same chips plus chips 12,13,20 and 21 were selected.

377 The mean number of hits is measured to be 124 and 89 in the $B = 0 \text{ T}$
378 and 1 T data sets respectively. The most probable values are respectively
379 87 and 64. Note that the $B = 0 \text{ T}$ data have a much larger Landau-like
380 tail than the 1 T data. Also the fluctuations in the core of the distribution
381 are larger. The mean time over threshold (ToT) is $0.68 \mu\text{s}$ for the $B = 0 \text{ T}$
382 and $0.86 \mu\text{s}$ at a $B = 1 \text{ T}$ data. A typical ToT distribution can be found

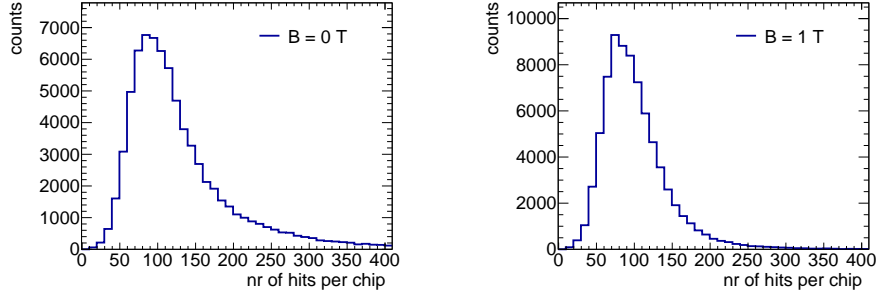


Figure 11: Distribution of the number of TPC track hits per chip for $B = 0$ T (left) $B = 1$ T data.

383 in Figure 5.5 of ref.[4]. The time over threshold is related to the deposited
 384 charge. This means that the deposited charge per pixel is smaller for the 0
 385 T data. The most probable value for the total deposited charge is similar for
 386 both data sets. A possible explanation for this behavior is that because of
 387 the reduced transverse diffusion in the $B = 1$ T data, the possibility of two
 388 primary electrons ending up in a single grid hole is higher. The mean number
 389 of hits is in agreement with the prediction of 106 electron-ion pairs for a 5
 390 and 6 GeV/c electron at $B = 0$ T for the T2K gas by [13], crossing 236 pixels
 391 or 12.98 mm and a detector running at 85% single electron efficiency. The
 392 measured single electron efficiency at this working point is in agreement with
 393 the efficiency vs mean time over threshold curve that was measured using a
 394 Fe source [4].

395 7. Conclusion and outlook

396 A TPC module with 32 GridPix chips was constructed and the perfor-
 397 mance was measured using data taken in a testbeam at DESY in 2021. The

398 TPC could be operated reliably and used a 93.6/5.0/1.4 gas mixture (by vol-
399 ume) of Ar/iC₄H₁₀/CO₂ with a small amount of oxygen and water vapour.
400 The analysed data were taken at electron beam momenta of 5 and 6 GeV/c
401 and at magnetic fields of 0 and 1 T.

402 The result for the transverse diffusion coefficient D_T is $287 \pm 0.5) \mu\text{m}/\sqrt{\text{cm}}$
403 at $B = 0$ T and D_T is $121 \pm 0.5 \mu\text{m}/\sqrt{\text{cm}}$ at $B = 1$ T. The longitudinal
404 diffusion coefficient D_L is measured to be $251 \pm 14) \mu\text{m}/\sqrt{\text{cm}}$ at $B = 0$ T
405 and $224 \pm 14 \mu\text{m}/\sqrt{\text{cm}}$ at $B = 1$ T. Results for the tracking systematical
406 uncertainties in xy were measured to be smaller than $13 \mu\text{m}$ with and with-
407 out magnetic field. The tracking systematical uncertainties in z were smaller
408 than $15 \mu\text{m}$ ($B = 0$ T) and $20 \mu\text{m}$ ($B = 1$ T).

409 The mean number of hits is in agreement with the predictions of [13] and
410 a detector running at 85% single electron efficiency.

411 Not all data were analysed and users are welcome to study them using
412 the data sets on available on the Grid.

413 The GridPix detector will be further tested and developed in view of a
414 TPC that will be installed in a heavy ion experiment at the EIC or other
415 future colliders. A follow up paper is in preparation on the measured dE/dx
416 or dN/dx resolution and other performance topics.

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