

1 Towards a Pixel TPC part II: performance of a 32 chip  
2 GridPix detector

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9 **Abstract**

10 A Time Projection Chamber (TPC) module with 32 GridPix chips was con-  
11 structed and the performance was measured using data taken in a testbeam at  
12 DESY in 2021. Part I of the results were described in a previous paper. The  
13 analysed data were taken at electron beam momenta of 5 and 6 *GeV/c* and at  
14 magnetic fields of 0 and 1 Tesla(T).

15 *Keywords:* Micromegas, gaseous pixel detector, micro-pattern gaseous  
16 detector, Timepix, GridPix, pixel time projection chamber

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17 **1. Introduction**

18 As a step towards a Pixel Time Projection Chamber for a future collider ex-  
19 periment [5], [6], a module consisting of 32 GridPix chips based on the Timepix3  
20 chip was constructed. The GridPix chips have a very fine granularity of 55x55  
21  $\mu\text{m}^2$  and a high efficiency to detect single ionisation electrons.

22 The 32 GridPix chip module was put in a test beam at DESY and comple-  
23 mented with two sets of Mimosas26 silicon detector planes. The analysed data  
24 were taken at electron beam momenta of 5 and 6 *GeV/c* and at magnetic fields  
25 of 0 and 1 T.

26 A description of the construction of the GridPix module, the test beam setup  
27 and data taking conditions can be found in part I of our paper [3]. The paper

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28 explains the track reconstruction procedure and the precise TPC tracking results  
29 that were obtained.

## 30 **2. Analysis topics**

31 In the following sections the analysis results for different topics will be pre-  
32 sented. Firstly, the particle identification performance using  $dE/dx$  or  $dN/dx$   
33 will be measured. Secondly, the single electron efficiency at high hit rates will  
34 be determined. Thirdly, the characterisation of large hit bursts caused by highly  
35 ionizing particles will be presented. Fourth, the resolution in the precision plane  
36 as a function of the incident track angle will be measured. Finally, the projected  
37  $dE/dx$  performance for a Pixel TPC in the ILD experiment will be presented  
38 and discussed.

### 39 *2.1. Particle Identification using $dE/dx$ or $dN/dx$*

40 The distribution of the number of TPC track hits per chip for the  $B = 0$  T  
41 and for the  $B = 1$  T data sets are a starting point for a measurement of the  
42  $dE/dx$  or  $dN/dx$  performance. As was shown in Part I of the paper [3], the  
43 mean number of hits is measured to be 124 and 89 in the  $B = 0$  T and 1 T data  
44 sets respectively. The most probable values are respectively 87 and 64.

45 In order to measure the track performance of  $dE/dx$  or  $dN/dx$ , the central  
46 chips - defined in ref [3] - were selected and calibrated to give the same mean  
47 number of hits per chip. By combining the hits associated to the track, a new  
48 1 m long track is formed. The 1 m long track has a coverage of 60% because  
49 inactive regions (chip edges and e.g. guard) are included. By applying different  
50 analysis methods, the  $dE/dx$  or  $dN/dx$  resolution can be measured from data.

51 The first method rejects large clusters with more than 6 hits in 5 consecutive  
52 pixels. Finally a  $dE/dx$  truncation at 90% is performed using samples of 20  
53 pixels; so the 10% largest  $dE/dx$  values are removed and  $dE/dx$  re-estimated.  
54 This method doesnot fully exploit the full granularity of the pixel TPC.

55 The second method exploits the distribution of the minimum distance in the  
56 pixel plane between consecutive hits. If only single electron clusters were made

57 in a gas, one would expect an exponentially falling distance distribution. Multi-  
 58 electron clusters will give rise to a peak at low distances tthat is smeared out  
 59 by the transverse diffusion process. The slope of the exponential distribiution  
 60 is proportional to the  $dN/dx$  i.e. the clusters produced by the electron. The  
 61 long Landau tail in the  $dE/dx$  distribution is coming from the multi-electron  
 62 clusters that peak at low distances.

63 Using a large number of tracks, it is possible to measure from data the  
 64 shape of the minimum distance distribution. At distances above 10 pixels the  
 65 distribution follows an exponential distribution. At lower distance weights for  
 66 the B = 0 T and 1 T data are determined and applied to ensure an exponential  
 67 distribution over the whole range. Finally, per 1 m track, a fit to distance  
 68 distribution in data is performed with the following template function:

$$N(d_{xy}) = N_0 \text{ weight}(d_{xy}) e^{-\text{slope } d_{xy}}. \quad (1)$$

69 where  $d_{xy}$  is the minimum distance in the precision plane (xy). The slope and  
 70  $N_0$  - normalisation - are left free in the per track fit, the weights are fixed using  
 71 the whole data set.

72 The testbeam data provides a  $dE/dx$  or  $dN/dx$  measurement for electrons.  
 73 The data were also used to perform a measurement of response of a MIP particle  
 74 - here defined as a particle that produced 70% of the electron  $dE/dx$ . By  
 75 dropping 30% of the hits associated to the track and applying the two methods,  
 76 the performance could be measured and the linearity of the technique tested.

77 The relative resolution is defined as the r.m.s. of the distribution divided  
 78 by the mean and the results are shown in Table 1. The resolution of the B=  
 79 1 T data is about 40% better than the B= 0 data. This is consistent with the  
 80 smaller fluctuations that are present in the distributions of the number of hits  
 81 per chip in the B= 1T data [3]. The template fit method has in the B= 1T  
 82 data a 20% better performance than the  $dE/dx$  truncation method. One might  
 83 argue that with more diffusion the results from the template fit method will  
 84 move more towards the results of the  $dE/dx$  truncation method. Note however  
 85 that the diffusion contribution to the track resolution in the 1 T data is already

Table 1:  $dE/dx$  or  $dN/dx$  resolution for different methods and data sets

Method	B=0 T resolution	B=1 T resolution
-	%	%
1 $dE/dx$ truncation	6.0	3.6
2 template fit	5.4	2.9

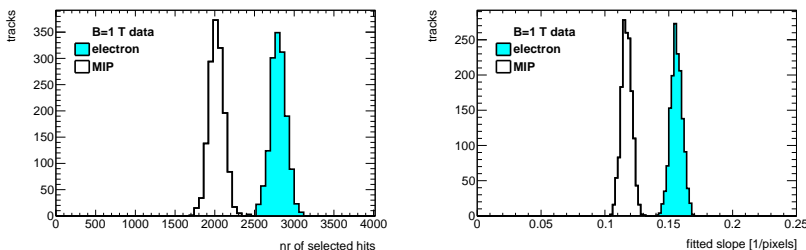


Figure 1: Distribution of the number of selected hits (left) for the  $dE/dx$  truncation method and the fitted slope (right) for the template fit method for an electron (light blue shaded) and MIP 1 m long tracks with 60% coverage for the B = 1 T data.

86 sizeable compared to the pixel size and varies between 85-150  $\mu\text{m}$ .

87 The results for the 1 T data are shown in Figure 1 for electrons and MIPs  
 88 for  $dE/dx$  truncation and template fit methods. The linearity - defined as the  
 89 mean MIP response divided by the mean electron response divided by 0.7 - was  
 90 measured to be 1.03 for method 1 and 1.07 for method 2. This value is slightly  
 91 different from 1, and can be corrected for by scaling the expected values for  
 92 different particles as a function of the measured momentum.

93 The performance result of the 32 chip GridPix detector is impressive. It  
 94 has currently, the best performance (running at atmospheric pressure) of con-  
 95 structed TPCs - and demonstrates the particle identification capabilities of a  
 96 GridPix Pixel TPC.

## 97 2.2. Single electron efficiency at high hit rates

98 The efficiency of the device to detect a hit in a high (low) rate environment  
 99 is measured comparing the mean time over threshold for low and high rate

Table 2: Mean time of threshold and rates for different runs

run	B	ToT 1	ToT 2	triggers	run time	Hits 1	Hits 2	trig rate	Rate 1	Rate 2
	T	$\mu s$	$\mu s$	$10^3$	$10^3 s$	$10^6$	$10^6$	Hz	hits/s	hits/s
6916	0	0.628	0.653	16.8	23.2	6.25	13.1	0.72	269	565
6934	0	-	0.651	73.4	2.41	-	20.5	30.4	-	8479
6935	0	0.620	-	7.39	2.41	6.95	-	30.6	2878	-
6969	1	0.650	0.666	7.94	13.8	1.93	2.16	0.57	139	156
6983	1	0.657	0.678	6.79	2.83	11.6	14.1	24.1	4110	4986

100 runs at B fields of 0 and 1 T. The mean time over threshold is sensitive to the  
 101 single electron efficiency of the detector. In order to extract a precise result, hits  
 102 associated to TPX3 tracks were used. The track selection is the same as the one  
 103 that was described in the subsection on the particle identification performance.  
 104 The analysed runs for the B=0 T data set were runs 6916, 6934 and 6935 and  
 105 for the B=1 T data set run 6969 and 6983.

106 For each run the mean ToT values were measured for values between 0.15  
 107 and 1.4  $\mu s$ . These cuts were applied to remove the noise and the tail.

108 The results for the measured average time over threshold for different runs  
 109 and hit rates are summarised in Table 2. ToT 1(2) denotes the mean time over  
 110 threshold for upper and lower half of the module and Hits 1(2) corresponds to  
 111 number of recorded raw hits. The mean Rate 1 (2) was calculated dividing the  
 112 total number of raw hits by the total run time. For the B=0 T data two high  
 113 rate runs 6934 and 6935 had to be analysed because the beam crossed either  
 114 the upper or the lower part of the module and therefore no measurement could  
 115 be performed (denoted by -). The statistical uncertainties are - due to the high  
 116 statistics - negligible.

117 The relative change in the mean time over threshold for the B=0 data is  
 118 -1.2% (upper) and -0.3% (lower). In this case the rate goes up to 8.5 kHz for 6  
 119 chips or 1.4 kHz per chip. The relative change in the mean time over threshold  
 120 for the B=1 T data is +1% (upper) and +1.7% (lower) The rate goes up to 5

121 kHz for 6 chips or 1.2 kHz per chip.

122 The relative change in the mean time over thresholds  $\delta\text{ToT}/\text{ToT}$  can be  
123 related to the relative change in the single electron efficiency  $\delta\epsilon/\epsilon$  by:

$$\delta\text{ToT}/\text{ToT} = c \delta\epsilon/\epsilon. \quad (2)$$

124 The slope  $c$  is about 0.5 at the mean working point of  $\text{ToT}=0.65 \mu\text{s}$  and is  
125 determined from the measured efficiency- $\text{ToT}$  curve in [5].

126 This means that the relative efficiency is stable at the level of +0.9% (B=1  
127 T) and -0.6% (B=0T) for hit rates up to 1.2 (1.4) kHz per chip. To conclude,  
128 running at hit rates up to 1.4 kHz per chip gives at most a reduction of 0.6% in  
129 the relative efficiency.

### 130 *2.3. Characterisation of hit bursts*

131 In event displays large hit burst caused by highly ionizing particles (e.g.  
132 alpha's or delta electrons) can be observed. A Pixel TPC is well suited to study  
133 and characterize these typical hit bursts. A pixel TPC also allows to improve  
134 the high momentum tracking by removing these bursts.

135 To study the hits burst the data of run 6969 B=1 T was analysed. No  
136 acceptance or track selection cuts were applied. A burst was selected with more  
137 than 100 hits in a radius of 50 pixels around the burst center within a time  
138 window of 200 ns around the mean time. The mean position in xy and the mean  
139 time of the burst were iteratively estimated. The bursts were characterized by  
140 the number of associated hits, the radius in which 90% of the hits are found  
141 (radius90) and the time in which 90% of the hits are detected (time90).

142 The distributions for radius90 and time90 for different burst sizes are shown  
143 in Figure ??.

144 For high momentum tracking it is important to cut tightly on the track  
145 residuals in xy and z. In particular the cut in z reduces the impact of bursts  
146 in the B=1 T data. One could in addition run a burst finding algorithm and  
147 downweight the hits associated to burst and the selected track. This will remove  
148 biases and improve the track parameter estimation.

149 *2.4. Resolution study*

150 Here, the resolution in the precision plane as a function of the incident track  
151 angle will be measured.

152 *2.5. Projected  $dE/dx$  performance*

153 Here, the projected  $dE/dx$  performance for a Pixel TPC in ILLD will be  
154 presented.

155 **3. Conclusion and outlook**

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