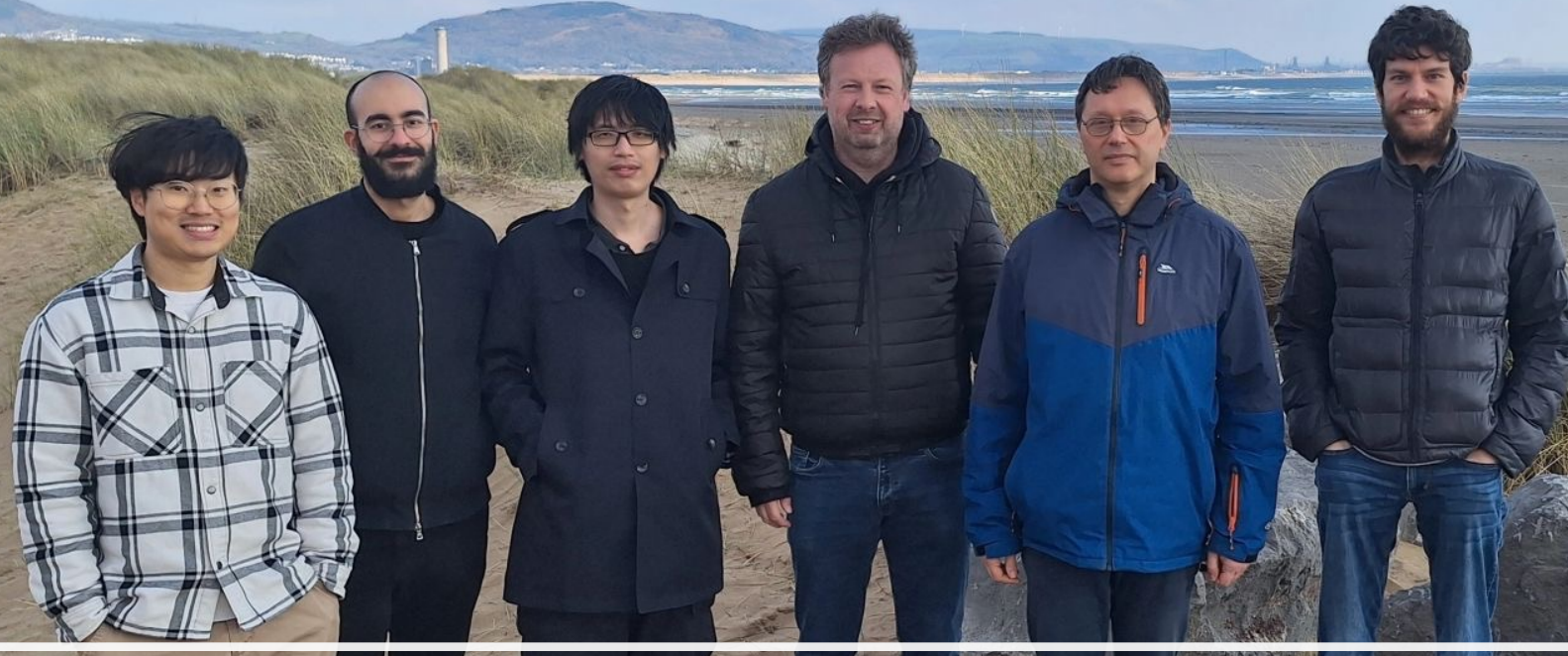


From lattice field theory to machine learning and back

Gert Aarts

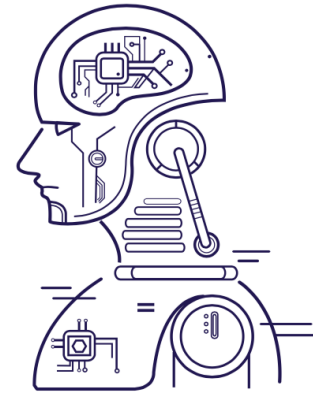


Swansea ML-LFT group



Chanju Park, Diaa Habibi, Shiyang Chen, GA, Biagio Lucini, Matteo Favoni

AIMLAC



UKRI CENTRE FOR DOCTORAL TRAINING
IN ARTIFICIAL INTELLIGENCE, MACHINE
LEARNING AND ADVANCED COMPUTING

- 5 cohorts of O(11) PhD students each
- first cohort has graduated
- final cohort of 17 students started in Oct 2023

- 60% particle physics/astronomy
- 25% computer science/maths
- 15% health

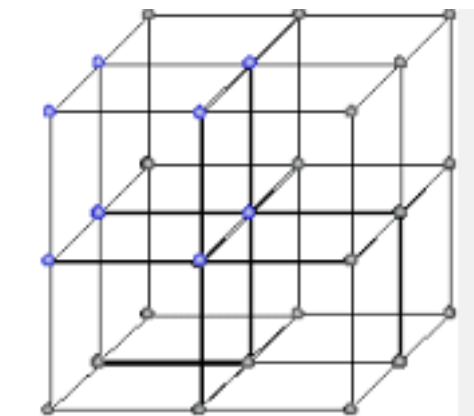


Lattice field theory (LFT) \leftrightarrow ML

- many fascinating connections between ML and statistical physics/field theory
- what can we learn that is relevant for \rightarrow or \leftarrow direction

lattice field theory: very loosely speaking

- many fluctuating degrees of freedom on a lattice/graph \leftrightarrow Markov random fields
- thermalisation and non-equilibrium dynamics \leftrightarrow learning
- stochastic quantisation and Langevin dynamics \leftrightarrow diffusion models
- renormalisation group
- ...



Conceptual and practical questions

can experience in quantum field theory help in understanding ML and vice versa?

- quantum field-theoretic machine learning

Dimitrios Bachtis, GA, Biagio Lucini, Phys. Rev. D 103 (2021) 074510 [[2102.09449](#) [hep-lat]]

- scalar field restricted Boltzmann machines as an ultraviolet regulator

Chanju Park, Biagio Lucini, GA, Phys. Rev. D 109 (2024) 034521 [[2309.15002](#) [hep-lat]]

- stochastic quantisation and diffusion models

Lingxiao Wang, GA, Kai Zhou, JHEP (to appear) [[2309.17082](#) [hep-lat]]

NeurIPS 2023 [2311.03578](#) [hep-lat]

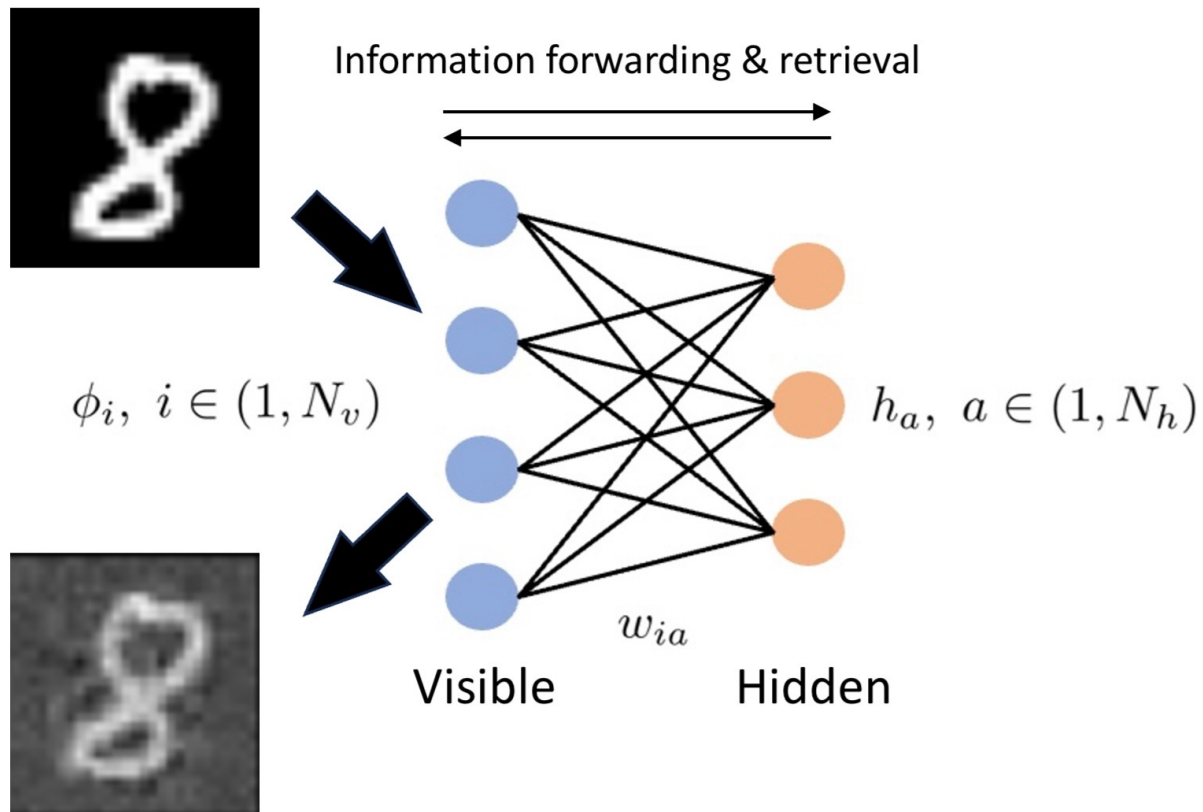
Topic 1

scalar field restricted Boltzmann machines as an ultraviolet regulator

Chanju Park, Biagio Lucini, GA

Phys. Rev. D 109 (2024) 034521 [[2309.15002](#)] [hep-lat]

Restricted Boltzmann Machine: generative network

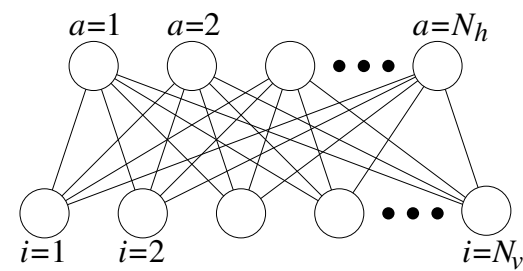


- energy-based method
- probability distribution
- binary or continuous d.o.f.

$$p(\phi, h) = \frac{1}{Z} e^{-S(\phi, h)}$$

$$Z = \int D\phi D h e^{-S(\phi, h)}$$

Scalar field RBM



- treat RBM as a lattice field theory with action

$$S(\phi, h) = \sum_i \frac{1}{2} \mu_i^2 \phi_i^2 + \sum_a \frac{1}{2\sigma^2} (h_a - \eta_a)^2 - \sum_{i,a} \phi_i w_{ia} h_a$$

- only quadratic terms, add interactions later, e.g ϕ^4 terms
- learn weight matrix w_{ia} and bias η_a (put to 0 below)
- induced distribution on visible layer

$$p(\phi) = \int Dh p(\phi, h) = \frac{1}{Z} \exp \left(-\frac{1}{2} \sum_{i,j} \phi_i K_{ij} \phi_j + \sum_i J_i \phi_i \right)$$

Gaussian scalar field RBM

- induced distribution on visible layer $p(\phi) = \int Dh p(\phi, h) = \frac{1}{Z} \exp \left(-\frac{1}{2} \sum_{i,j} \phi_i K_{ij} \phi_j + \sum_i J_i \phi_i \right)$
- scalar field with kinetic (all-to-all) term $K_{ij} = \mu_i^2 \delta_{ij} - \sigma^2 \sum_a w_{ia} w_{aj}^T$
and source $J_i = \sum_a w_{ia} \eta_a$
- unusual Gaussian LFT: what is the weight matrix W and bias η ?
- learn from data or directly from known distribution
- simplest case: target theory = LFT of free scalar field in 1 or 2d, $K^\phi \approx p^2 + m^2$

$$WW^T = \frac{1}{\sigma^2} (\mu^2 \mathbf{1} - K^\phi) \equiv \mathcal{K}$$

Infinite # of solutions ($N_h = N_v$)

infinite # solutions for weight matrix, use that \mathcal{K} is symmetric and positive-definite

1. Cholesky decomposition $\mathcal{K} = LL^T$: $W = L$ triangular
2. diagonalisation $\mathcal{K} = ODO^T = O\sqrt{D}O^T O\sqrt{D}O^T$: $W = W^T = O\sqrt{D}O^T$
3. non-uniqueness: internal symmetry $W \rightarrow WO_R \rightarrow \phi^T Wh \rightarrow \phi^T WO_R h = \phi^T Wh'$

in practice

- realisation depends on initialisation
- non-observable degeneracy due to internal symmetry on hidden layer

$$K_{ij} = \mu_i^2 \delta_{ij} - \sigma^2 \sum_a w_{ia} w_{aj}^T$$

Dependence on N_h and μ^2

- what if $N_h < N_v$? role of hyperparameter μ^2 ?
- use LFT insights: target distribution is scalar field theory
- method: SVD of weight matrix $W = U \Xi V^T$
- quadratic kernel $K = \mu^2 \mathbb{1} - \sigma_h^2 U \Xi \Xi^T U^T = U [\mu^2 \mathbb{1} - \sigma_h^2 \Xi \Xi^T] U^T \equiv U D_K U^T$

- eigenvalues of quadratic kernel

$$D_K = \text{diag}(\underbrace{\mu^2 - \sigma_h^2 \xi_1^2, \mu^2 - \sigma_h^2 \xi_2^2, \dots, \mu^2 - \sigma_h^2 \xi_{N_h}^2}_{N_h}, \underbrace{\mu^2, \dots, \mu^2}_{N_v - N_h})$$

- both N_h and μ^2 act as **ultraviolet regulators**

What if $N_h < N_v$?

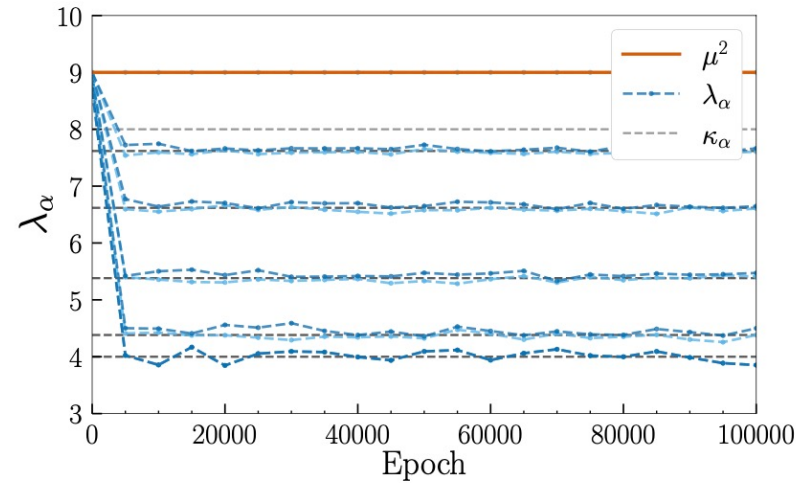
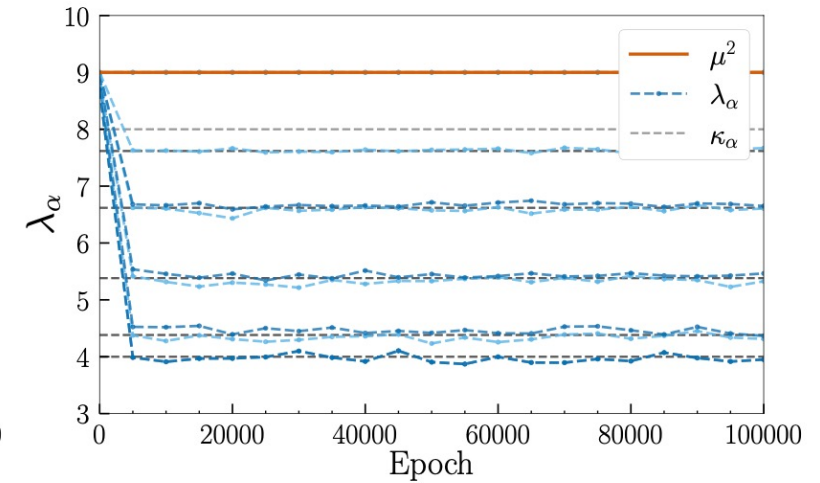
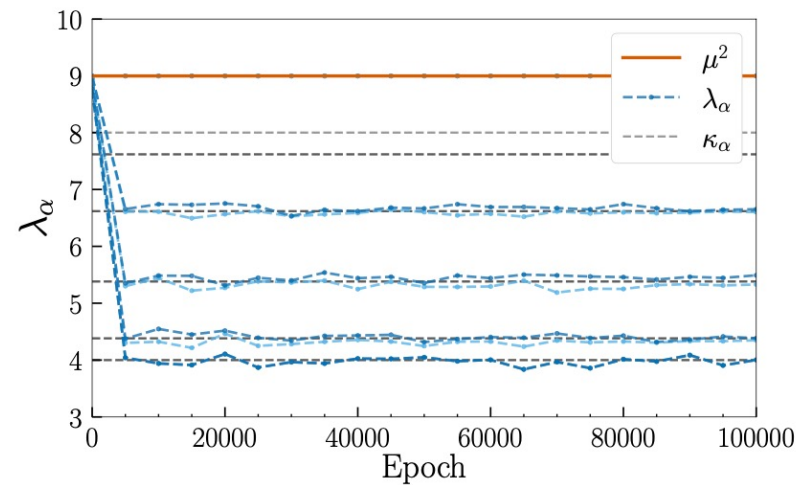
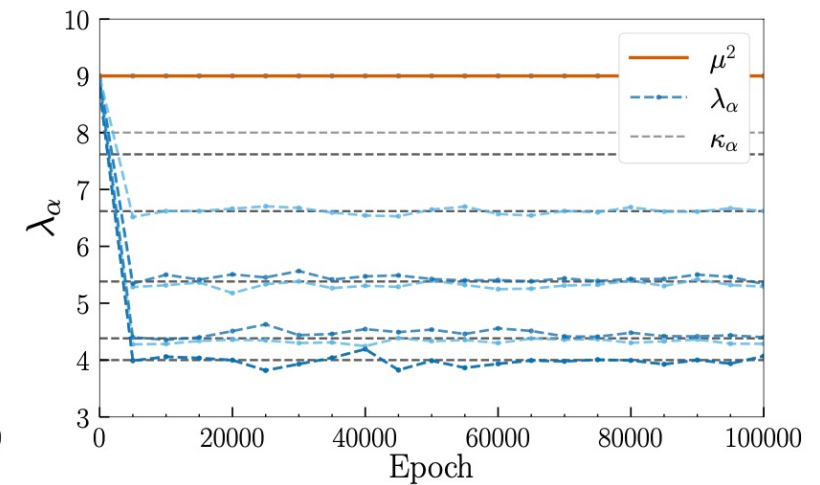
- example: scalar LFT with $N_v = 10$ nodes

- exact spectrum:

$$\kappa \sim p^2 + m^2$$

- reproduced by RBM (λ) from smallest eigenvalue upwards

- higher modes are moved to cut-off scale (μ^2)

(a) $N_h = 9$ (b) $N_h = 8$ (c) $N_h = 7$ (d) $N_h = 6$

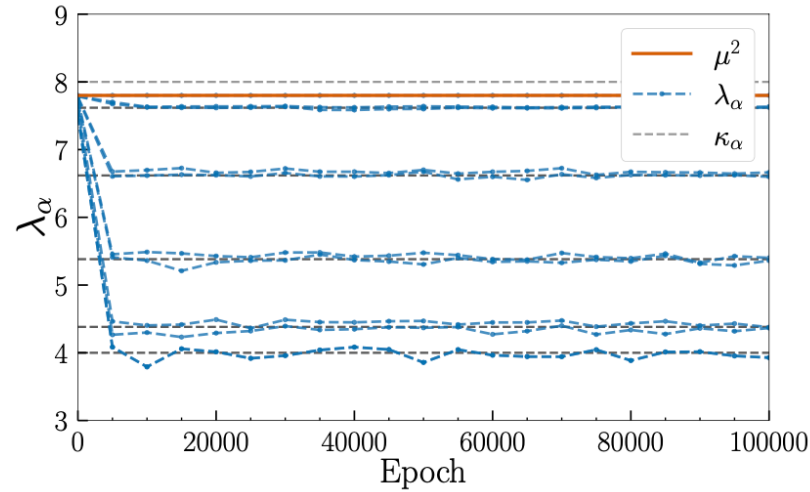
What if RBM mass $\mu^2 < \lambda_{max}$?

- example: scalar LFT with $N_v = 10$ nodes

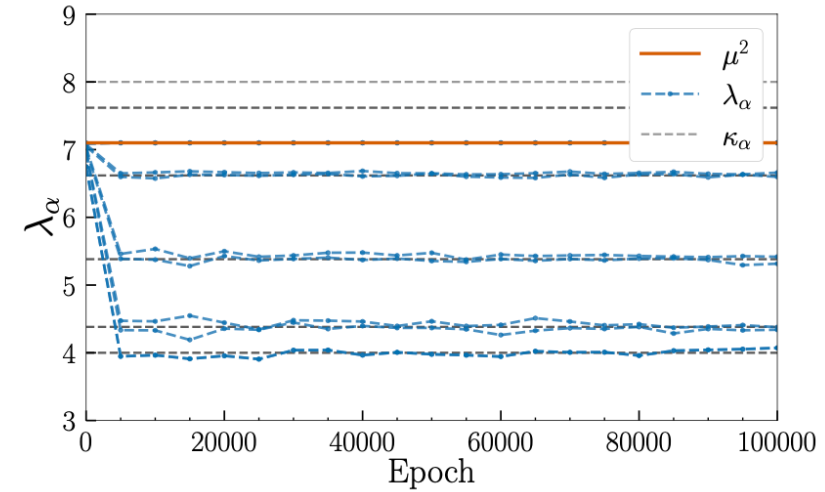
- exact spectrum:

$$\kappa \sim p^2 + m^2$$

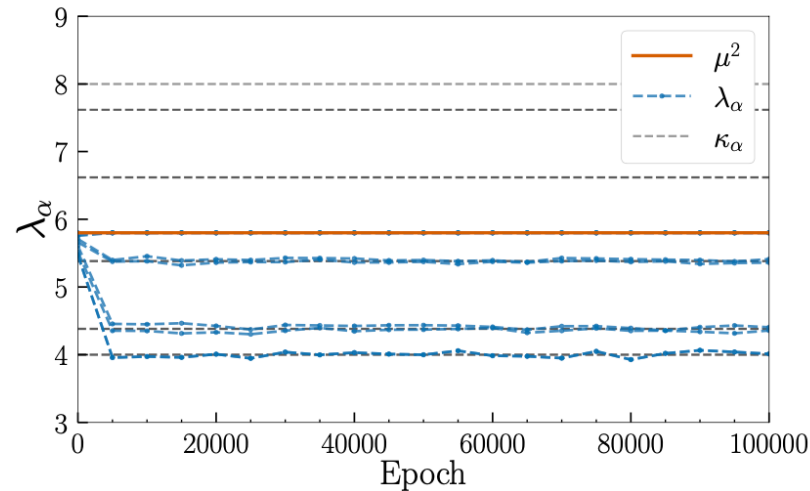
- reproduced by RBM (λ) from smallest eigenvalue upwards
- higher modes are suppressed at cut-off scale (μ^2)



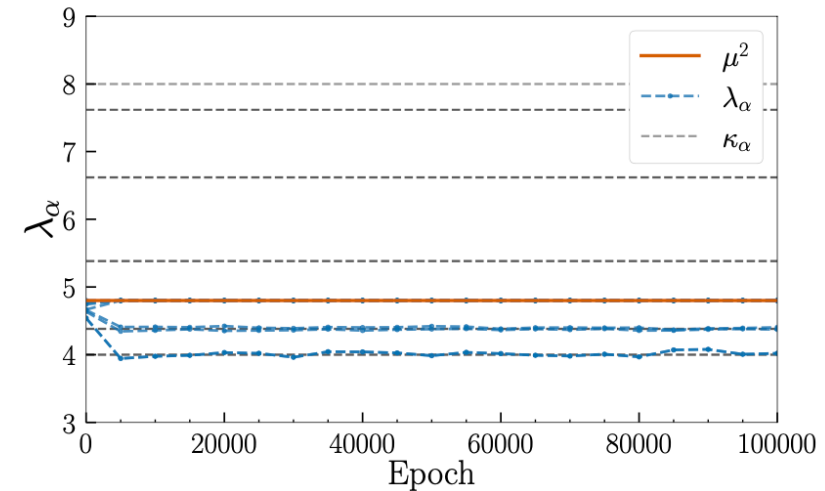
(a) $\mu^2 = 7.8$



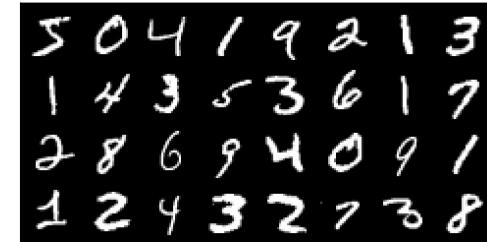
(b) $\mu^2 = 7.1$



(c) $\mu^2 = 5.8$

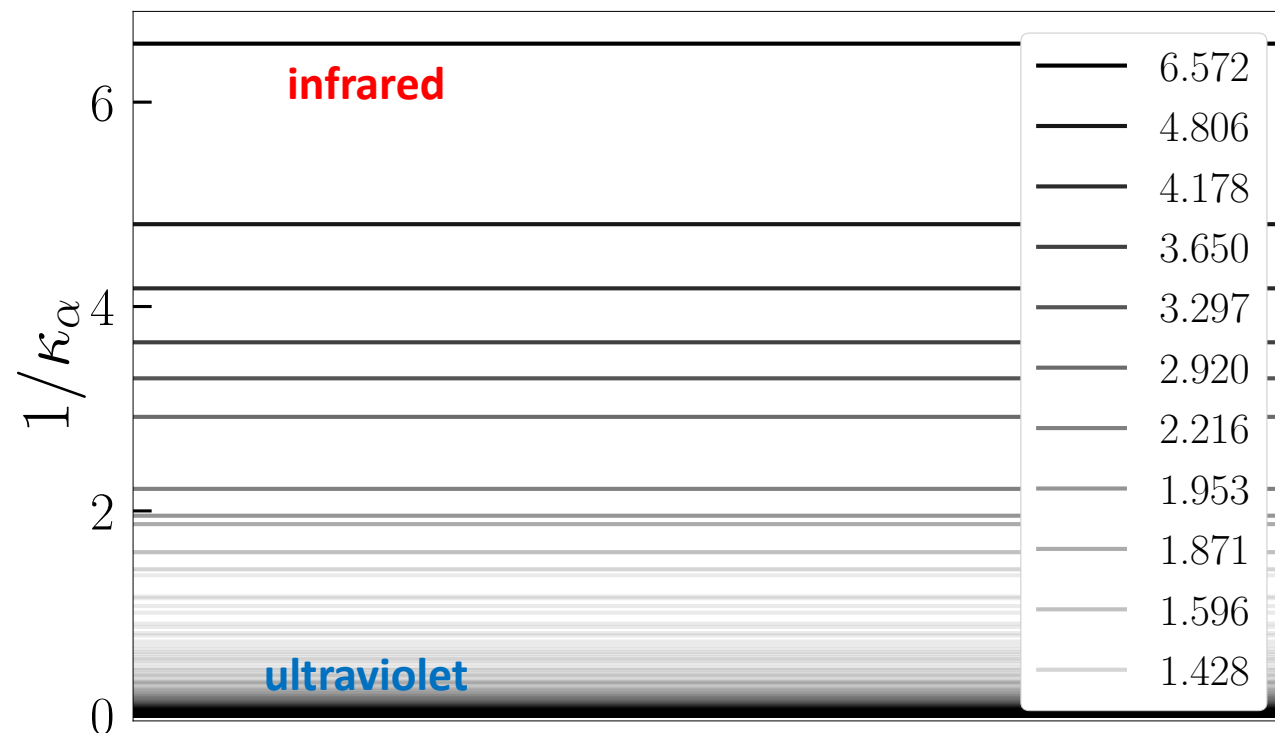


(d) $\mu^2 = 4.8$



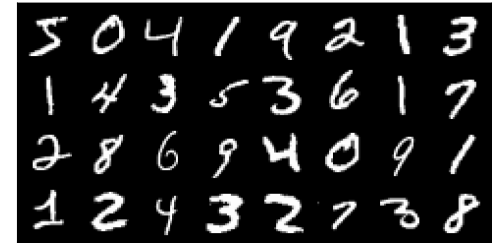
RBM as ultraviolet regulator

- relevant for “real” data sets? MNIST: 28x28 images of digits
- compute spectrum of two-point correlator $K_{ij}^{-1} = \langle \phi_i \phi_j \rangle_{data}$
- inverse spectrum $1/\kappa$
- infrared safe
- ultraviolet divergent

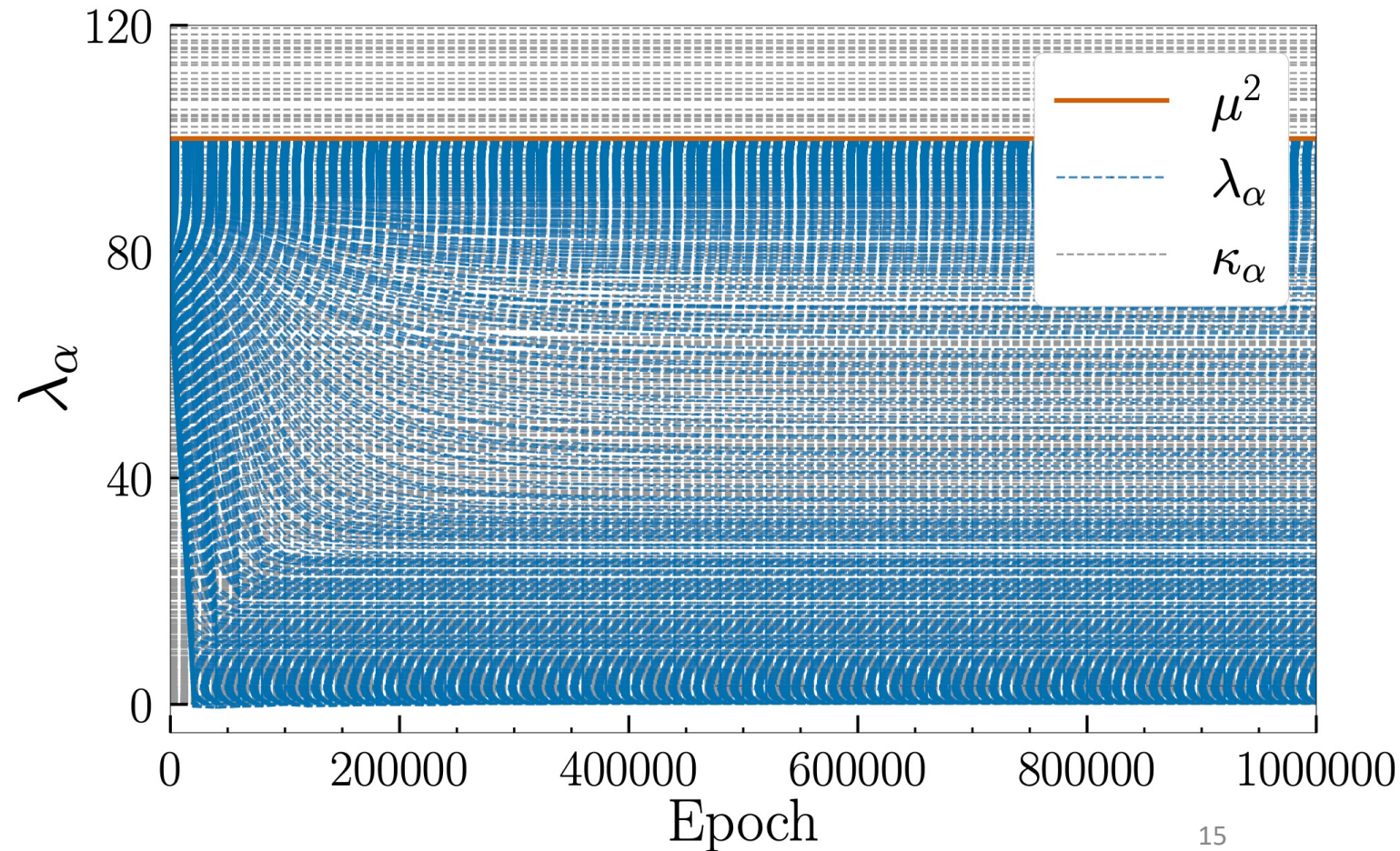


784 eigenvalues

MNIST with fixed RBM mass



- $N_v = N_h = 784$
- fixed RBM mass $\mu^2 = 100$
- spectrum regulated
- infrared modes learned correctly

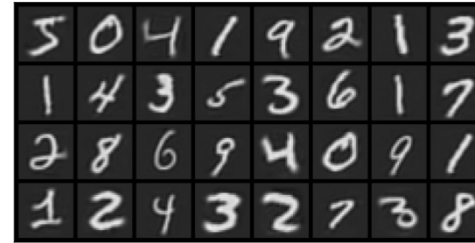


MNIST with $N_h \leq N_v$

what is the effect of including incomplete spectrum?



(a) $N_h = 784$



(b) $N_h = 225$



(c) $N_h = 64$

removal of ultraviolet modes affects generative power



(d) $N_h = 36$



(e) $N_h = 16$



(f) $N_h = 4$

Topic 2

stochastic quantisation and diffusion models

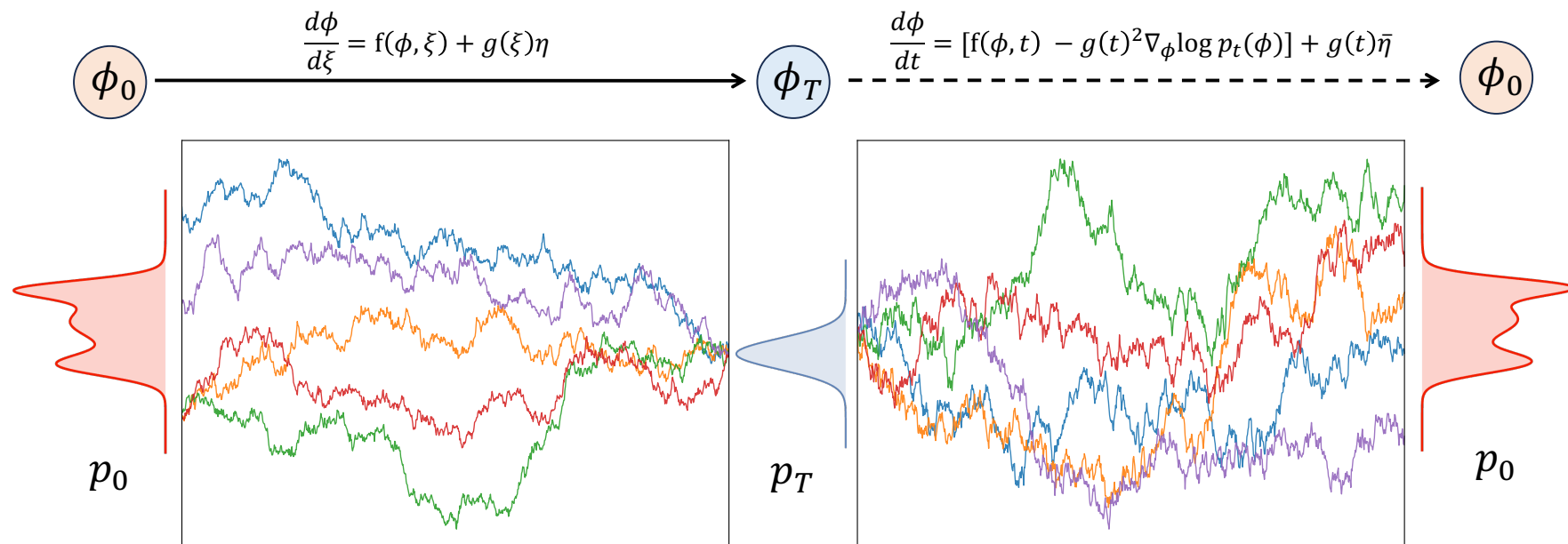
Lingxiao Wang, GA, Kai Zhou

JHEP (to appear) [[2309.17082](#) [hep-lat]]

NeurIPS 2023 [[2311.03578](#) [hep-lat]]

Diffusion models

- solve stochastic process with a particular drift/force/score
- drift is learnt during forward diffusion process, starting from data
- new configurations are generated via backward process using learnt drift



Stochastic quantisation

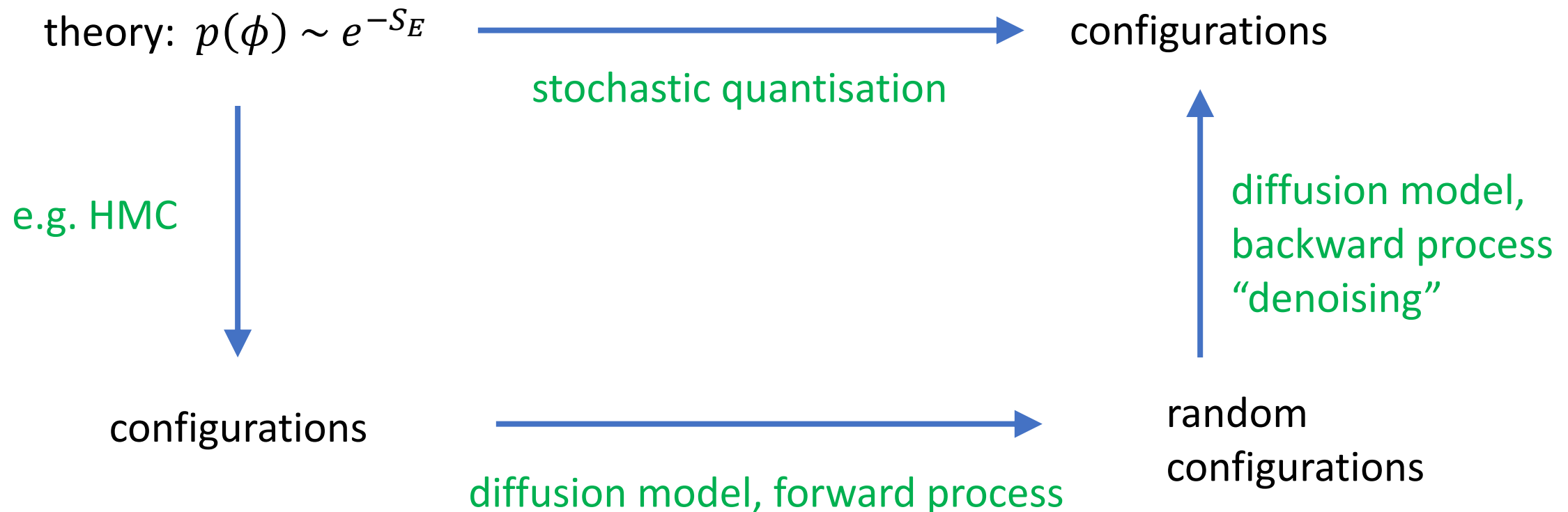
- ideas well-known in quantum field theory: stochastic quantisation (Parisi & Wu 1980)
- path integral quantisation via a stochastic process in fictitious time

$$\frac{\partial \phi(x, \tau)}{\partial \tau} = -\frac{\delta S_E[\phi]}{\delta \phi(x, \tau)} + \eta(x, \tau) \quad \langle \eta(x, \tau) \eta(x', \tau') \rangle = 2\alpha \delta(x - x') \delta(\tau - \tau')$$

- equilibrium solution ($\tau \rightarrow \infty$): distribution $p(\phi) \sim e^{-S_E}$
- convergence guaranteed for real actions due to properties of Fokker-Planck equation
- create samples from Euclidean path integral
- applied to non-abelian gauge theories and QCD in 1980s, but superseded by other methods such as Hybrid Monte Carlo (HMC) [stepsize dependence, efficiency]

Stochastic quantisation and diffusion models

- diffusion models as an alternative approach to stochastic quantisation

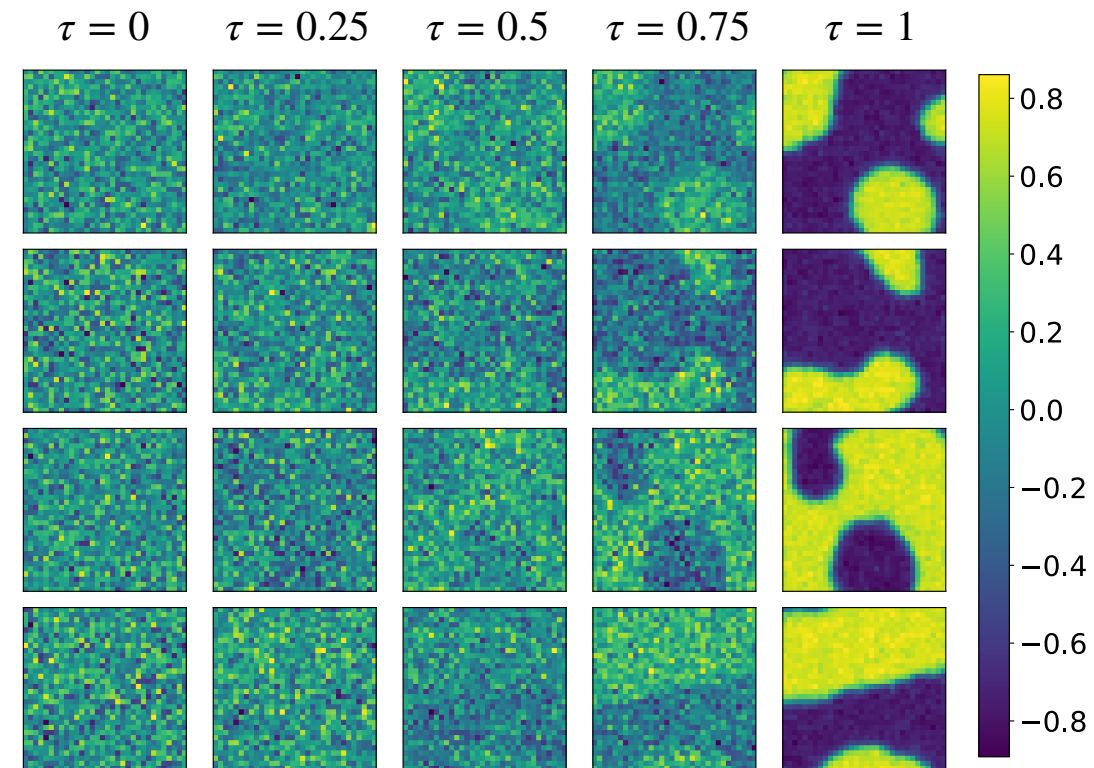


Diffusion model for 2d ϕ^4 scalar theory

- 32^2 lattice, choice of action parameters in symmetric and broken phase
- training data set generated using Hybrid Monte Carlo (HMC)
- variance expanding DM trained using U-Net architecture

generating configurations:

- broken phase
- “denoising” (backward process)
- large-scale clusters emerge, as expected



Diffusion model for 2d ϕ^4 scalar theory

generating configurations in symmetric phase

- compute magnetisation $\langle M \rangle$, susceptibility χ_2 , Binder cumulant U_L
- compare with test HMC data set (with same statistics)

data-set	$\langle M \rangle$	χ_2	U_L
Training (HMC)	0.0012 ± 0.0007	2.5160 ± 0.0457	0.1042 ± 0.0367
Testing (HMC)	0.0018 ± 0.0015	2.4463 ± 0.1099	-0.0198 ± 0.1035
Generated (DM)	0.0017 ± 0.0015	2.4227 ± 0.1035	0.0484 ± 0.0959

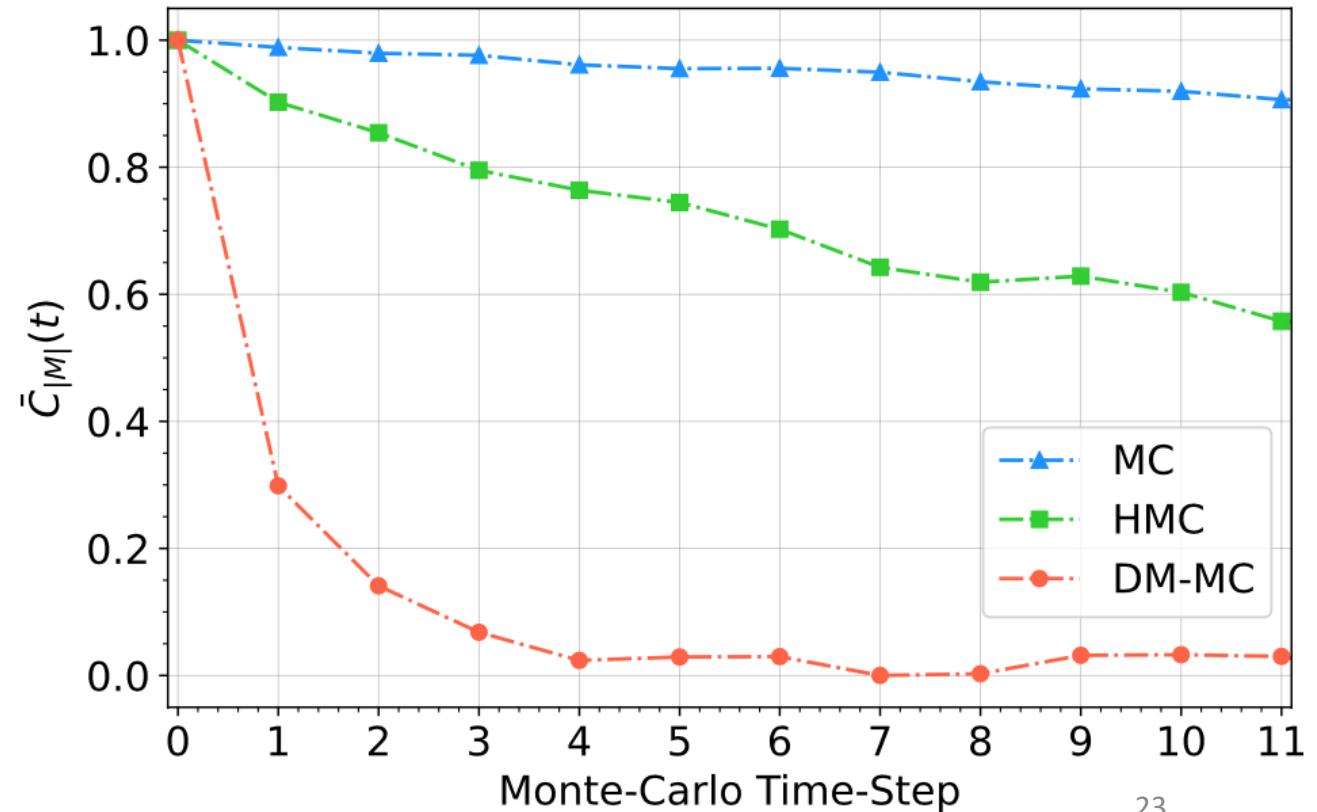
- good agreement is observed

Diffusion model for 2d ϕ^4 scalar theory

- auto-correlation time (first comparison)
- normalised auto-correlation function

overall:

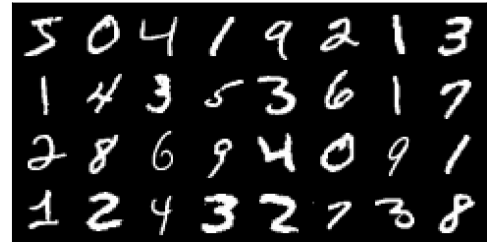
- proof of principle
- expected results obtained
- need to do detailed comparison of precision, speed and scalability



Summary and outlook

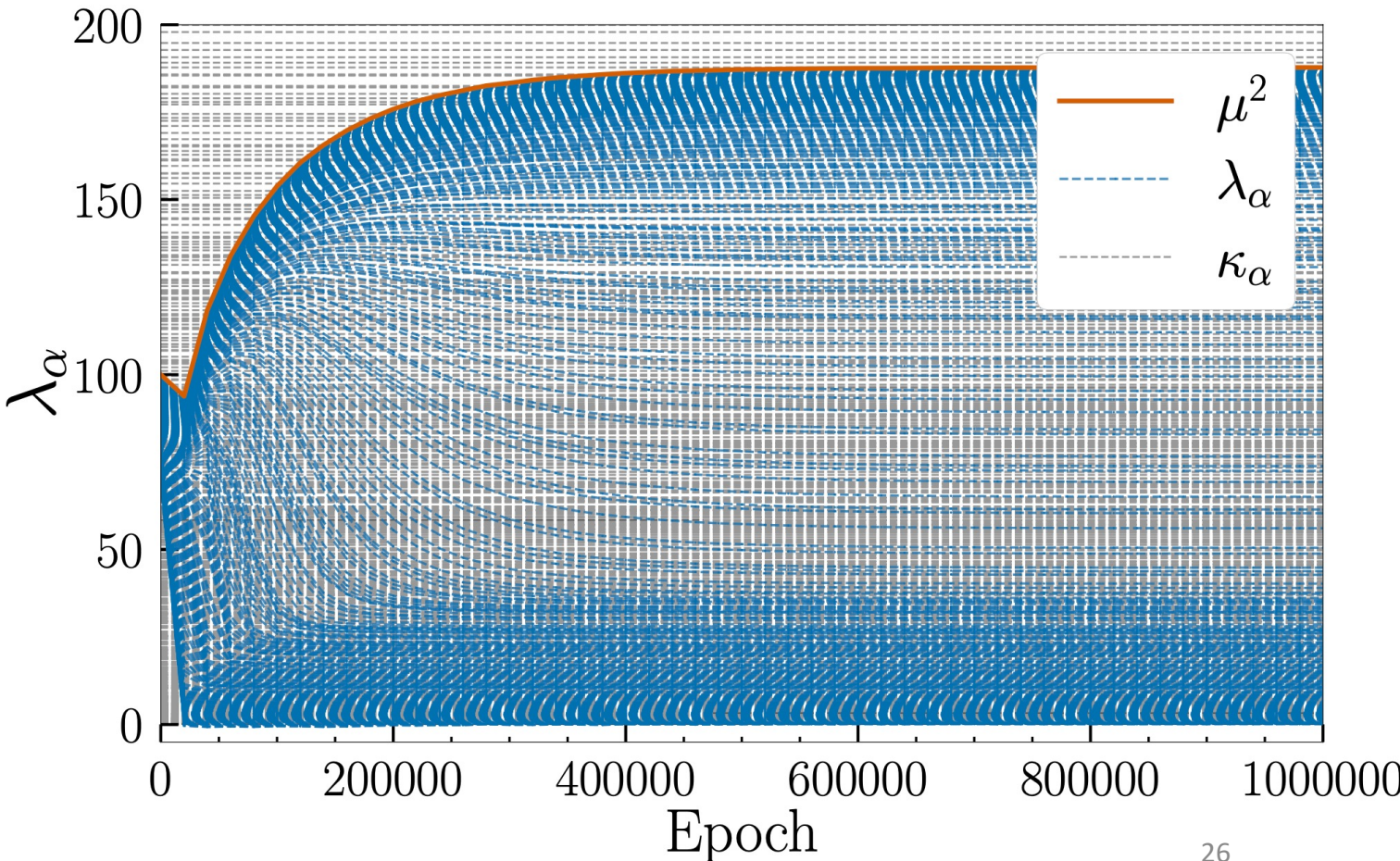
- can experience in quantum field theory help in understanding ML and vice versa?
- two examples of interplay between lattice field theory and ML
 - scalar field RBM as ultraviolet regulator
 - stochastic quantisation and diffusion models
- interplay between statistical/lattice field theory and ML
- many directions to explore

Back-up



MNIST with dynamic RBM mass

- $N_v = N_h = 784$
- dynamical RBM mass μ^2 is learned as well
- spectrum regulated
- ultraviolet cut-off μ^2 increases to include more modes



Interacting scalar field RBM

- Gaussian RBMs can learn Gaussian distributions
- in LFT language: need to include interactions
- various ways to do so, depending on properties of target distribution
- QFT-ML approach: add local potential terms on nodes, e.g. ϕ^4 terms

Quantum field-theoretic machine learning, Bachtis, Aarts, Lucini
Phys. Rev. D 103 (2021) 074510 [2102.09449 [hep-lat]]

- standard RBM approach: use binary hidden layer $h_a = \pm 1$

Scalar-Bernoulli RBM: hidden binary nodes

- induced distribution $p(\phi) = \frac{1}{Z} \exp \left(-S_\phi(\phi) + \sum_a \sum_{n=1}^{\infty} c_n \psi_a^{2n} \right)$ with $\psi_a = \sum_i \phi_i w_{ia} - \eta_a$

- generates all-to-all interactions of all powers of ϕ

- at leading order in W same kinetic term as in Gaussian case

- example of quartic term
(taking $\eta_a = 0$ for simplicity)

$$\sum_a \sum_{i,j,k,l} (\phi_i w_{ia}) (\phi_j w_{ja}) (\phi_k w_{ka}) (\phi_l w_{la})$$

- highly non-local, very different from standard field theories, analysis in preparation

Stochastic quantisation: complex actions

- approach not limited to real-valued distributions/actions
- extend Langevin process to complex manifold: complex Langevin dynamics (Parisi 1981)
- complexify d.o.f.: real scalar \rightarrow complex scalar, $U \in SU(N) \rightarrow U \in SL(N, \mathbb{C})$
- convergence not guaranteed, no general solution of Fokker-Planck equation
- a posteriori justification (GA, Seiler, Stamatescu 2009, Nagata, Nishimura, Shimasaki 2016)
- applied to problems at finite (baryon) chemical potential (GA & Stamatescu 2008, Aarts 2009)
- success in some theories, QCD remains difficult (e.g. Sexty 2019)
- introductory lectures: GA, J.Phys.Conf.Ser. 706 (2016) 2, 022004 [[1512.05145](#)] [[hep-lat](#)]