

NNV section for (astro)particle physics fall meeting - 03/11/23

Probing the Width of the Higgs Boson Through Offshell Decays

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General Introduction

At the LHC, and in particular ATLAS and CMS, we study the Higgs Boson by colliding **protons** at energies high enough to produce the Higgs momentarily



m₄₁ [GeV]

Ni

[1]

- Just over 10 years ago we managed to find it, both in the ATLAS and CMS experiments, but large questions about its properties still remain
 - In this analysis, we aim to answer one such question, namely: "Can we measure the decay width of this elusive particle?"
 - Trailer: Yes! Through "offshell" Higgs decays, but it is not trivial

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Probing new Physics Through the Width

- While the width of the Higgs is an interesting property to measure in and of itself, this analysis **might also hint towards, or constrain, new physics phenomena**
 - As the only known mechanism to give particles mass, the Higgs should decay into all massive particles, so **new massive particles should affect its total decay width**



Introduction - Offshell Decays?

 With the collisions of these protons (which consist of gluons for a large part) we create the Higgs through the following process, among others



- Due to it being virtual, the Higgs does not have to obey the energy-momentum relation, and can decay at a higher energy than its pole mass ("offshell")
 - Intuitively, the Heisenberg uncertainty relation allows for this brief violation $\Delta t \Delta E \ge \hbar/2$
- The Breit-Wigner shape of the Higgs is very narrow
 - The $t\bar{t}$ and WW intermediate states enhance offshell Higgs decay at $2m_t$ and $2m_W$
 - Offshell accounts for $\mathcal{O}(10\%)$ of all $gg \to H$ decays!



Introduction - Width from Offshell Higgs?

• Now how do we get the width from this? Consider cross section of $gg \rightarrow H \rightarrow WW$



Introduction - Quantum Interference

- One of the major "backgrounds" in this analysis is $gg \rightarrow WW$ without the Higgs intermediate state
 - Processes with the same initial and final states exhibit quantum interference effects
 - In the cross section, i.e. "probability for an interaction to occur", we get

 $\sigma_{gg \to (H \to)WW} \propto |M_{gg \to (H \to)WW}|^2 = |M_{gg \to H \to WW}|^2 + |M_{gg \to WW}|^2 + M_{gg \to H \to WW}M_{gg \to WW}$



Introduction - Quantum Interference

• We parametrize our analysis using the **offshell signal strength** $\mu_{off} \propto \Gamma_{H}^{tot}$

• Which tells us our signal strength as a function of **the different contributions**

$$\frac{100}{10} + \mu_{off}^{ggF} * S + sqrt(\mu_{off}^{ggF}) * I + \mu_{off}^{ggF} + S + sqrt(\mu_{off}^{ggF} + S + sqrt(\mu_{off}^{ggF}) * I + \mu_{off}^{ggF} + S + sqrt(\mu_{$$

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$$\sigma_{gg \to (H \to)WW} = \mu_{offshell}\sigma_S + \sigma_B + \sqrt{\mu_{offshell}}\sigma_I$$



Analysis Strategy

Analysis Topologies



• We consider the different flavour ($WW \rightarrow e\nu\mu\nu$) and same flavour ($WW \rightarrow e\nue\nu/\mu\nu\mu\nu$) decay modes

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This separation helps us isolate specific processes and backgrounds

Analysis Composition



Analysis Optimisation

• Cancellation of destructive interference with signal **reduces sensitivity**

- We need a fitting variable that **discriminates between signal and interference**
- Mass of WW system provides this naturally, but since neutrino's cannot be measured in detector, we cannot fully reconstruct WW mass
 - Instead, we construct V_3 (=0. $x \times M_{ll} + M_T$), as a proxy to M_{WW}



• We also **train and cut on a DNN** to reduce all other (non-interfering) backgrounds



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Results

Higgs Offshell Measurement

- We reject the hypothesis that the Higgs \widehat{H} does not decay offshell with an expected significance of 1.4 σ
 - The expected upper limit of µ_{off} is 4.5 at the 95% confidence interval level
 ATLAS run 1 H → WW: 20.3



- The $H \rightarrow ZZ$ [3] analysis achieved an expected upper limit of 2.4
 - While the H → WW decay mode has a 7 times larger cross section, H → ZZ is more sensitive due to having fully reconstructable final states
 - This in turn increases resolution and decreases systematic uncertainties



[3] Aad, Georges, et al. Evidence of off-shell Higgs boson production from ZZ leptonic decay channels and constraints on its total width with the ATLAS detector. No. arXiv: 2304.01532. ATLAS-HIGG-2018-32-003, 2023.

Onshell Combination and Width

• By combining with the onshell measurement we can calculate $\frac{\Gamma_H}{\Gamma_H^{SM}} = \frac{\mu_{offshell}}{\mu_{onshell}}$





Conclusions and Prospects

Run 3 Prospects

- Currently the LHC is in run 3, which is planned until end of 2025
 - The goal was to collect 300 fb⁻¹, but due to a difficult year at the LHC we are now approximately 30 fb⁻¹ behind schedule
 - In run 2 we collected 140 fb⁻¹
- Assuming an identical analysis, with 280 fb⁻ of collected data and 50% reduced theory systematic uncertainties
 - Improvement of μ_{off} limits by approx. 20%, similar improvement expected in the width
 - Potential advanced analysis techniques might improve this further



Future Developments

- Combination with $H \rightarrow ZZ$ offshell analysis expected in near future for **best limit possible with current data** on Γ_H
 - Potential for 3σ exclusion of the "no offshell Higgs decays" hypothesis!
- Developing analysis techniques for run 3, in particular in the context of advanced machine learning techniques such as inference-aware learning



Alex Martone

Sacha Mugnier







Backup

Results Split by Channel

- In terms of **decay modes** the different flavour channels are relatively sensitive
 - Higher signal yields and purities
 - Reduced systematic uncertainties in different flavour region compared to same flavour region



Results Split by Channel

In terms of production modes the VBF channels are relatively sensitive
 Theory uncertainties in VBF are reduced compared to ggF



Fitting Variable

- Cancellation of destructive interference with signal reduces sensitivity
 - We need a variable that discriminates between ^{²/₂}
 signal and interference, such that they don't cancel¹⁰⁰
 - Mass of WW system provides this naturally
- However, since neutrino's cannot be measured in detector, we cannot fully reconstruct WW mass
 - Instead, we construct V_3 (=0. $x \times M_{ll} + M_T$), as a proxy to M_{WW}
- Alternative attempted approach was to **train a DNN to regress** M_{WW} , but this did not outperform simply V_3 in its accuracy to M_{WW}



DNN Strategy



DNN split up into three bins: low, med, and high signal purity

- In each, M_{WW} proxy used in fit to capture unique shape effects from interference and signal interplay
- Essentially this comprises unrolling the DNN
- Low signal DNN bin included effectively as a



Spin and Helicity of decaying WW bosons









Same Flavour Analysis

Events normalized ATLAS Other H VVV Mis-Id Internal Other VV Z/y $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$ Тор 1.2 $H \rightarrow WW^* \rightarrow evev/\mu v \mu v$ ww W Uncertainty Asimov 0.8 0.6 0.4 0.2 **Relatively large contribution of** .05 $Z\ell\ell$ background, other -,414,414,40,414,4 *╾*∶*┾∖₩*₽*∖*┾*∖*┾*\₩*₩**┼/┼₩ 1.95 backgrounds controlled with 0.9 CR ZII CR ZII SR DNN bin 2 SR DNN bin 2 SR DNN bin 3 SR DNN bin 3 SR DNN bin 2 SR DNN bin 3 SR DNN bin 1 SR DNN bin CR ZI SR DNN bin different flavour control regions 20 1400 0-jet 1-jet 2-jet Data 2010 (Vs = 7 TeV) $Z \rightarrow ee$ Lents/ 1000 Events/ QCD $Z \to \ell \ell$ 800 600 400 200F q 70 80 90 Selected for in CR with $|m_{\ell\ell} - m_Z| < 15 \text{ GeV}_2$ Ni

Notoriously bad modelling of $Z \rightarrow \ell \ell$ sample generated with Sherpa hampers this regions' sensitivity

L dt = 36 pb

Central $Z \rightarrow ee$

ATLAS

100

110 rh_{ee} [GeV]

Likelihood Scan Shape



ef

Nik

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