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Universal relations to measure neutron star properties from targeted r-mode searches

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R-modes as source of CGW

- Toroidal mode of fluid oscillation in NS for which the restoring force is the Coriolis force.
- ➤ Unstable in all rotating stars due to the CFS mechanism although dissipation mechanisms can damp and saturate the oscillations.
- ➢ Spindown of young pulsars leading to continuous GW emissions.



Courtesy : C. Hanna and B. Owen



- ➢ GW Searches using LVK data :
 - Crab pulsar (Rajbhandari et al. 2021)
 - PSR J0537-6910 with $n \approx 7$ (Fesik & Papa 2020, Abbott et al. 2021)
- No detection of GW but upper limits on r-mode amplitude were obtained.

R-modes frequency and Universal Relations

- > Newtonian limit ; $k = f_{GW} / f_{rot} = 2m / l(l+1)$
- ➤ Correction due to GR and rapid rotation:

$$\mathbf{f}_{\mathrm{GW}} / \mathbf{f}_{\mathrm{rot}} = \mathbf{A} - \mathbf{B} (\mathbf{f}_{\mathrm{rot}} / \mathbf{f}_{\mathrm{K}})^2$$

▶ For slowly rotating stars, effect of the term B is negligible.

Universal relation	κ	A
(compactness range)		
(Idrisy et al. 2015)(0.11-0.31)	0.614-0.433	1.39-1.57
Tabulated $EOS(38)(0.11-0.31)$	0.601 - 0.412	1.40 - 1.59
Tabulated $EOS(38)(0.10-0.35)$	0.608 - 0.364	1.39 - 1.64
Non-parametric $EOS(40)(0.11-0.31)$	0.596 - 0.415	1.39 - 1.59
Non-parametric $EOS(40)(0.10-0.35)$	0.604 - 0.371	1.40-1.63

range of A : 1.39 < A < 1.59 PSR J0537 : 86 - 99 Hz (86-97 Hz)





Universal relations (Tabulated EoS)

Ghosh et al ApJ 944, 53 (2023)

 $log(\Lambda)$



- Universal relation between r-mode \succ frequency and dimensionless tidal deformability.
- These can be used to estimate the \succ effects of dynamical tides from r-mode excitation during binary NS inspirals. (Gupta et al. arxiv: 2205.01182, Ma et al. **PRD 103, 063020 (2021)**)

 \succ

Distance estimation from binary mergers

- ➤ Inspiralling binaries are "Standard Siren"
- ➢ GW amplitude

$$h_0 = \frac{4\pi^2 G^{5/3}}{c^4} (f_{GW} \mathcal{M})^{5/3} \frac{1}{f_{GW} d}$$

- ► From the orbital decay energy $\dot{f}_{GW} = \frac{96}{5} \pi^{8/3} \left(\frac{G\mathcal{M}}{c^3}\right)^{5/3} f_{GW}^{11/3}$
- \succ Eliminate \mathcal{M} to get luminosity distance

$$d = \frac{5c}{24\pi^2} \frac{1}{h_0} \frac{\dot{f}_{GW}}{f_{GW}^3}$$

 $\succ \qquad \text{Measurement of } H_0 \text{ if location can be fixed by} \\ \text{independent method.}$

Schutz B. F., Nature 323, 310 (1986)

Distance estimation from CGW emission

➤ CGW are not so "Standard Siren"

GW amplitude
$$h_0 = \sqrt{\frac{512\pi^7}{5}} \frac{G}{c^5} (\alpha M R^3 \tilde{J}) \frac{1}{d} f_{GW}^3$$

Assuming Spin-down is dominated by r-mode emission

$$\dot{f}_{GW} = -\frac{4096\pi^7}{25} \frac{G}{c^7} \frac{M^2 R^6 \tilde{J}^2}{I} \alpha^2 f_{GW}^7$$

> Eliminate $\alpha M R^3 \tilde{J}$ to get luminosity distance $\sqrt{I} = \sqrt{f_{GW}} \sqrt{8c^3}$

$$\frac{\sqrt{I}}{d} = h_0 \sqrt{\frac{f_{GW}}{f_{GW}}} \sqrt{\frac{8c^3}{45G}}$$

Distance is always degenerate with moment of inertia.
 Sieniawska & Jones, MNRAS 509, 5179 (2022)

Measuring NS parameters from a r-mode detection



Measuring NS parameters from a r-mode detection



Measuring NS parameters from a r-mode detection



Signal Model and Fisher matrix

 \succ Upto second order in frequency derivatives, the phase is given by

Jaranowski & Królak, PRD 59, 063003 (1999)

$$\psi = \psi_0 + 2\pi \left[ft + \frac{1}{2}\dot{f}t^2 + \frac{1}{6}\ddot{f}t^3 \right]$$

> The Fisher Covariance matrix is given by the inverse of the below matrix

$$\Gamma = \frac{g^{-1}}{\rho^2} \text{ where } g_{ij} = \left\langle \frac{\partial \psi}{\partial f^{(i)}} \frac{\partial \psi}{\partial f^{(j)}} \right\rangle - \left\langle \frac{\partial \psi}{\partial f^{(i)}} \right\rangle \left\langle \frac{\partial \psi}{\partial f^{(j)}} \right\rangle$$

 $\sim \rho^2$ is the signal-to-noise ratio(SNR) assuming optimal match between the true signal and the best-fit template

$$\rho^{2} = \int_{0}^{\infty} \frac{4|\tilde{h}(f)|^{2}}{S_{n}(f)} df$$

> The error for any physical quantity can be calculated from the Fisher matrix information. For example,

$$\sigma \left(\frac{d}{\sqrt{I}}\right)^2 = \frac{45G}{8c^3} \frac{1}{(\pi\rho h_0)^2} \left[\frac{75}{T}\frac{\dot{f}}{f^3} + \frac{1620}{T^4f\dot{f}} + \frac{675}{T^3f^2} + \pi^2\frac{\dot{f}}{f}\right]$$

SNR Estimates

- \succ A canonical NS at a distance of 1 kpc
- ➢ Observing period of 2 years with a-LIGO and ET sensitivity curves.



Error Estimations for MoI(I) and distance(d)

- \succ Error estimates does not change much with different EOSs.
- > For most reasonable values of κ , MoI can be measured upto 10% accuracy for $f_{GW} > 100$ Hz.
- For detectable signals we can measure the distance up to an accuracy of 20% for $f_{GW} > 200$ Hz and $\alpha > 10^{-4}$ and these estimates scale accordingly with changing distance



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- Most promising candidate PSR J0537-6910 has $n \approx 7$.
- Possible to re-do energy budget calculation given knowledge of n. (Lu et al. MNRAS, 521 (2023))

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➤ Cosmological corrections not taken into account.

- Valid for galactic pulsars.
- Presence of solid crust (Yuri 2001), stratification (Andersson 2023), magnetic field (Luciano 2000) or superfluidity (Lindblom 2000) in the core that might affect the r-mode frequency.
 Effect negligible compared to GR and fast rotation corrections.

Summary and Future Works

- ➤ Universal relation between r-mode frequency and the NS compactness determines the search parameter space.
- ► The URs are updated with recent multi-messenger constraints on the NS EoS and consistent with their Newtonian limits.
- > These universal relations along with the knowledge of EOS can be used to break the degeneracy between distance and moment of inertia.
- ► These universal relations with prior knowledge of distance from EM observation can give strict constraints on the dense matter EOS.
- > Plans to extend this simplistic signal model to much more realistic analysis using a Bayesian formalism to infer NS properties from LIGO strain data.

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Thanks for listening !!