

# Does electron capture decay matter?

---

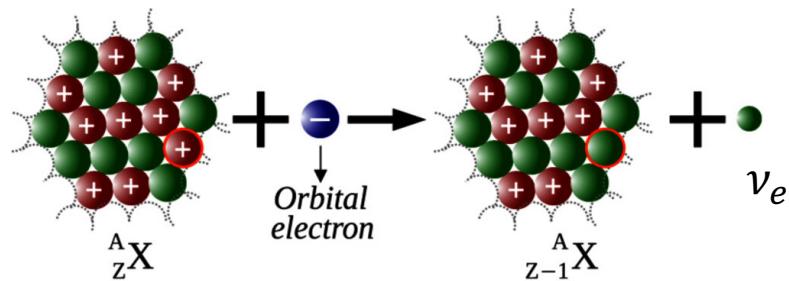
Revisiting Electron Capture decay in the context of high-precision cosmic-ray data

Marta Borchiellini

Kapteyn Astronomical Institute, RUG

in collaboration with D. Maurin and M. Vecchi

# Electron capture decay



- Electron capture (EC) decaying nuclei decay by capturing a K-shell electron
- Most Cosmic-Ray (CR) nuclei are completely ionized.
- EC decay in CR nuclei depends on attachment and stripping processes

# Scientific motivations

---

Direct CRs detection experiments are providing data that allow investigating new (astro)physical phenomena (e. g. dark matter) :

- AMS-02 measured high-precision cosmic-ray fluxes up to Iron (*Aguilar et al. 2021*)
- Recent measurement of the isotopic composition for  $29 < Z < 38$  by ACE-CRIS (*Binns et al. 2022*)

→ We need models as accurate as possible for comparison with data

# Scientific motivations

---

Direct CRs detection experiments are providing data that allow investigating new (astro)physical phenomena (e. g. dark matter) :

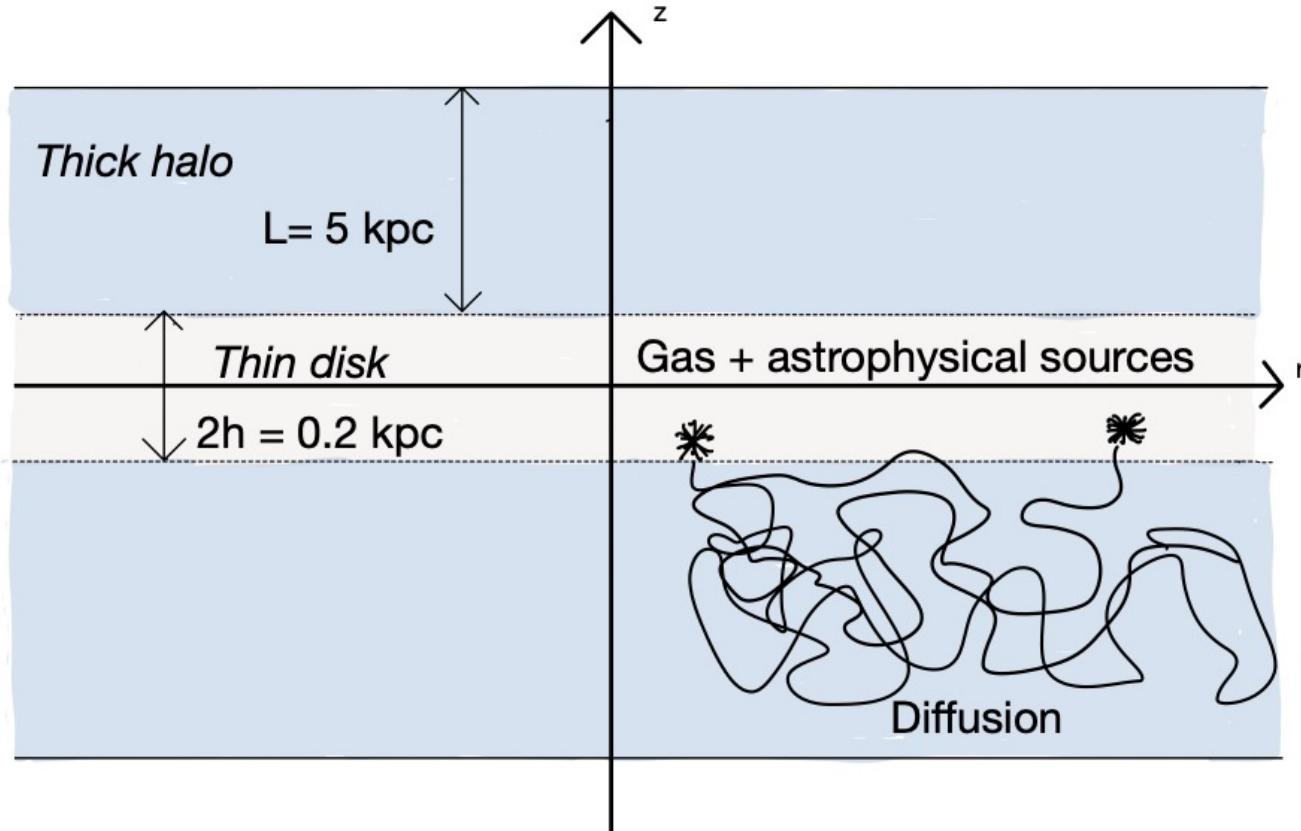
- AMS-02 measured high-precision cosmic-ray fluxes up to Iron (*Aguilar et al. 2021*)
- Recent measurement of the isotopic composition for  $29 < Z < 38$  by ACE-CRIS (*Binns et al. 2022*)

→ We need models as accurate as possible for comparison with data

Is electron capture decay relevant for the interpretation of these cosmic-ray data?

# Galaxy model

Good first-order description of the Milky Way for Galactic CR fluxes.



- 1D, observer at  $z=0$
- Thin disk: gas (density  $n_{\text{ISM}}$ ) and CR sources
- Thick halo: diffusion and confinement of CR

# CR propagation

*Transport equations:* coupled differential equations for the CRs number density.

$$\begin{cases} -D \frac{\partial^2 n_0}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z) q \\ -D \frac{\partial^2 n_1}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{cases}$$

CR number density

$$n = n_0 + n_1$$

$n_0$  = fully ionized

$n_1$  = one electron attached

Assuming:

- Steady state
- no energy losses
- 2 populated charged states

# CR propagation

---

$$\begin{cases} -D \frac{\partial^2 n_0}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z) q \\ -D \frac{\partial^2 n_1}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{cases}$$

# CR propagation

Diffusion (random walk) on  
magnetic inhomogeneites  
(disk and halo)

$$\left\{ \begin{array}{l} -D \frac{\partial^2 n_0}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z) q \\ -D \frac{\partial^2 n_1}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{array} \right.$$

# CR propagation

Diffusion (random walk) on magnetic inhomogeneities  
(disk and halo)

$$\left\{ \begin{array}{l} -D \frac{\partial^2 n_0}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z)q \\ -D \frac{\partial^2 n_1}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{array} \right.$$

Generic source term (disk)

# CR propagation

Diffusion (random walk) on magnetic inhomogeneities (disk and halo)

$$\left\{ \begin{array}{l} -D \frac{\partial^2 n_0}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z)q \\ -D \frac{\partial^2 n_1}{\partial z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{array} \right.$$

Generic source term (disk)

Inelastic interaction rate on gas (disk)

$$\Gamma^i = n_{ISM} v \sigma_{inel}$$

# CR propagation

Diffusion (random walk) on magnetic inhomogeneities (disk and halo)

Attachment rate of  $e^-$  (disk)

$$\Gamma^a = n_{ISM} v \sigma_{att}$$

Generic source term (disk)

$$\begin{cases} -D \frac{\partial^2 n_0}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z)q \\ -D \frac{\partial^2 n_1}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{cases}$$

Inelastic interaction rate on gas (disk)

$$\Gamma^i = n_{ISM} v \sigma_{inel}$$

# CR propagation

Diffusion (random walk) on magnetic inhomogeneities (disk and halo)

Attachment rate of  $e^-$  (disk)

$$\Gamma^a = n_{ISM} v \sigma_{att}$$

Generic source term (disk)

$$\begin{cases} -D \frac{\partial^2 n_0}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z)q \\ -D \frac{\partial^2 n_1}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{cases}$$

Inelastic interaction rate on gas (disk)

$$\Gamma^i = n_{ISM} v \sigma_{inel}$$

Stripping rate of  $e^-$  (disk)

$$\Gamma^s = n_{ISM} v \sigma_{strip}$$

# CR propagation

Diffusion (random walk) on magnetic inhomogeneities (disk and halo)

Attachment rate of  $e^-$  (disk)

$$\Gamma^a = n_{ISM} v \sigma_{att}$$

Generic source term (disk)

$$\begin{cases} -D \frac{\partial^2 n_0}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_0 + \Gamma^a n_0 - \Gamma^s n_1\} = 2h\delta(z)q \\ -D \frac{\partial^2 n_1}{\partial^2 z^2} + 2h\delta(z)\{\Gamma^i n_1 - \Gamma^a n_0 + \Gamma^s n_1\} + \Gamma^{EC} n_1 = 0 \end{cases}$$

Inelastic interaction rate on gas (disk)

$$\Gamma^i = n_{ISM} v \sigma_{inel}$$

Stripping rate of  $e^-$  (disk)

$$\Gamma^s = n_{ISM} v \sigma_{strip}$$

EC decay rate  $\Gamma^{EC}$  (only if  $e^-$  attached)

# Characteristic timescales

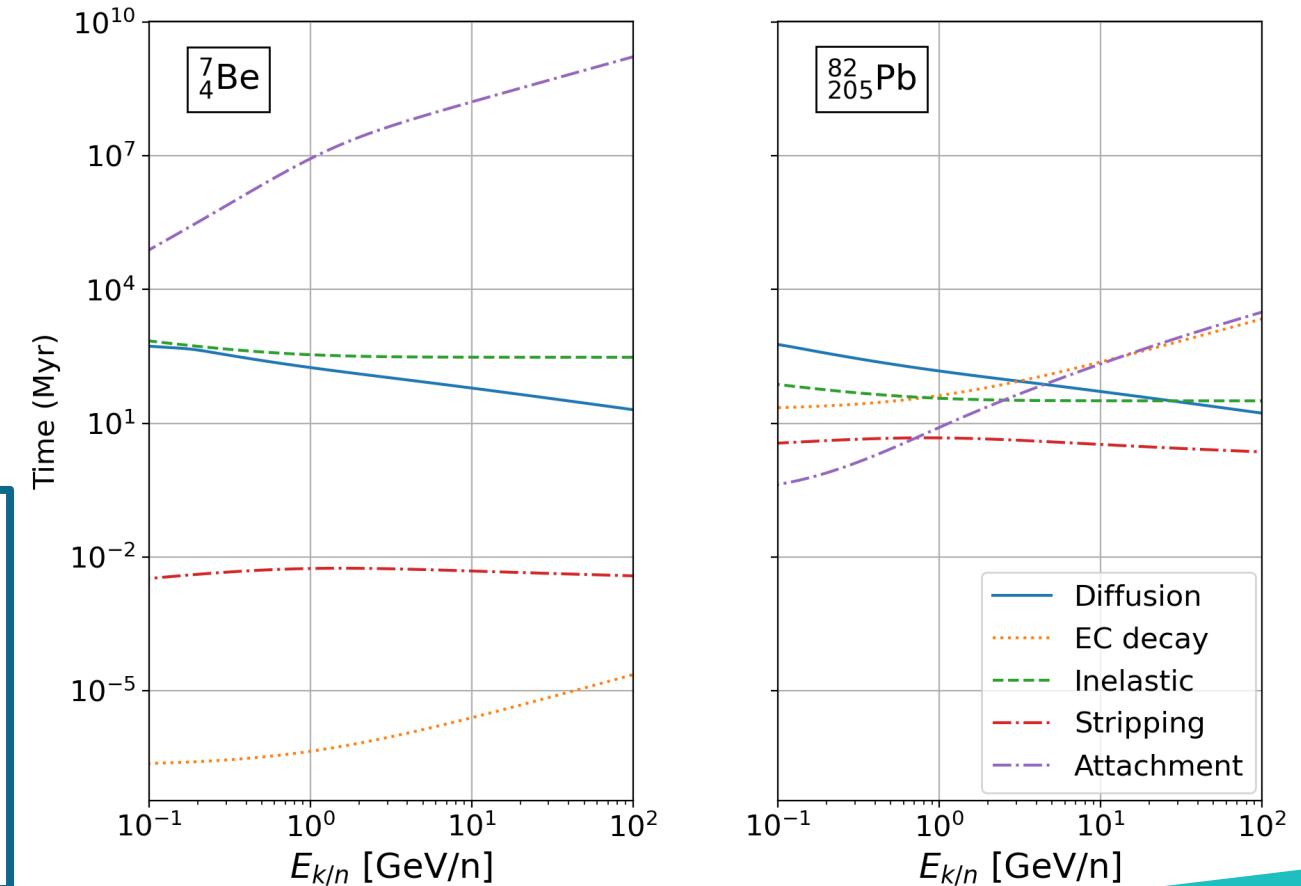
Diffusion	$t_D = \frac{L^2}{2D}$	$D \propto E^{0.5}$
Inelastic scattering	$t_{inel} = \frac{1}{n_{ISM} v \sigma_{inel}}$	$\sigma_{inel} \propto A^{2/3}$
Attachment	$t_a = \frac{1}{n_{ISM} v \sigma_{att}}$	$\sigma_{att} \propto \sigma(E)Z^2$
Stripping	$t_s = \frac{1}{n_{ISM} v \sigma_{strip}}$	$\sigma_{strip} \propto \sigma(E)Z^{-2}$
EC decay	$t_{EC} = \gamma \tau_{EC}$	$t_{EC} \propto E$

The lower the time,  
more dominant is the  
corresponding process

# Characteristic timescales

- Diffusion dominates above a few GeV/n
- Attachment more efficient than stripping for large Z

→ Net effect of EC decay depends on interplay between  $t_{\text{att}}$ ,  $t_{\text{strip}}$  and  $t_{\text{EC}}$   
→ EC decay expected to be relevant at low E and large Z



# Solving the transport equations

The trasport equations have been solved analitically:

- In thin disk approximation
- at  $z=0$ , to allow comparison with data

$$\left\{ \begin{array}{l} N_0 = \frac{q}{\frac{D}{hL} + \Gamma^i + \Gamma^a - \frac{\Gamma^s \Gamma^a}{f_1 + \Gamma^i + \Gamma^s}} \\ N_1 = \frac{\Gamma^a N_0}{f_1 + \Gamma^i + \Gamma^s} \end{array} \right. \quad \text{with} \quad f_1 = \sqrt{\frac{D \Gamma^{EC}}{h^2}} \coth \left( \sqrt{\frac{\Gamma^{EC}}{D}} L \right)$$

Flux:  $J = vN/4\pi$

# EC decaying isotopes

---

We used a selection of EC decaying isotopes from *Letaw et al., 1984, ApJS, 56, 36.*

EC decaying isotopes can be classified in two categories:

Isotope	$\tau_{1/2}$ (Myr)	Isotopic fraction
$^7_4\text{Be}$	$1.46 \cdot 10^{-7}$	0.55
$^{37}_{18}\text{Ar}$	$9.58 \cdot 10^{-8}$	0.30
$^{41}_{20}\text{Ca}$	$1.00 \cdot 10^{-1}$	0.07
$^{53}_{25}\text{Mn}$	3.70	0.35
$^{67}_{31}\text{Ga}$	$8.93 \cdot 10^{-9}$	0.07
$^{73}_{33}\text{As}$	$2.20 \cdot 10^{-7}$	0.36

# EC decaying isotopes

We used a selection of EC decaying isotopes from *Letaw et al., 1984, ApJS, 56, 36.*

EC decaying isotopes can be classified in two categories:

- Short-lived isotopes:

$$\tau_{\text{EC}} < 10^{-3} \text{ Myr}$$

Isotope	$\tau_{1/2}$ (Myr)	Isotopic fraction
$^7_4\text{Be}$	$1.46 \cdot 10^{-7}$	0.55
$^{37}_{18}\text{Ar}$	$9.58 \cdot 10^{-8}$	0.30
$^{41}_{20}\text{Ca}$	$1.00 \cdot 10^{-1}$	0.07
$^{53}_{25}\text{Mn}$	3.70	0.35
$^{67}_{31}\text{Ga}$	$8.93 \cdot 10^{-9}$	0.07
$^{73}_{33}\text{As}$	$2.20 \cdot 10^{-7}$	0.36

# EC decaying isotopes

We used a selection of EC decaying isotopes from *Letaw et al., 1984, ApJS, 56, 36*

EC decaying isotopes can be classified in two categories:

- Short-lived isotopes:  
 $\tau_{\text{EC}} < 10^{-3}$  Myr
- Intermediate-lived isotopes:  
 $10^{-3} < \tau_{\text{EC}} < 10^2$  Myr

NB: Escape from the Galaxy before decaying for  $\tau_{\text{EC}} > 10^2$  Myr

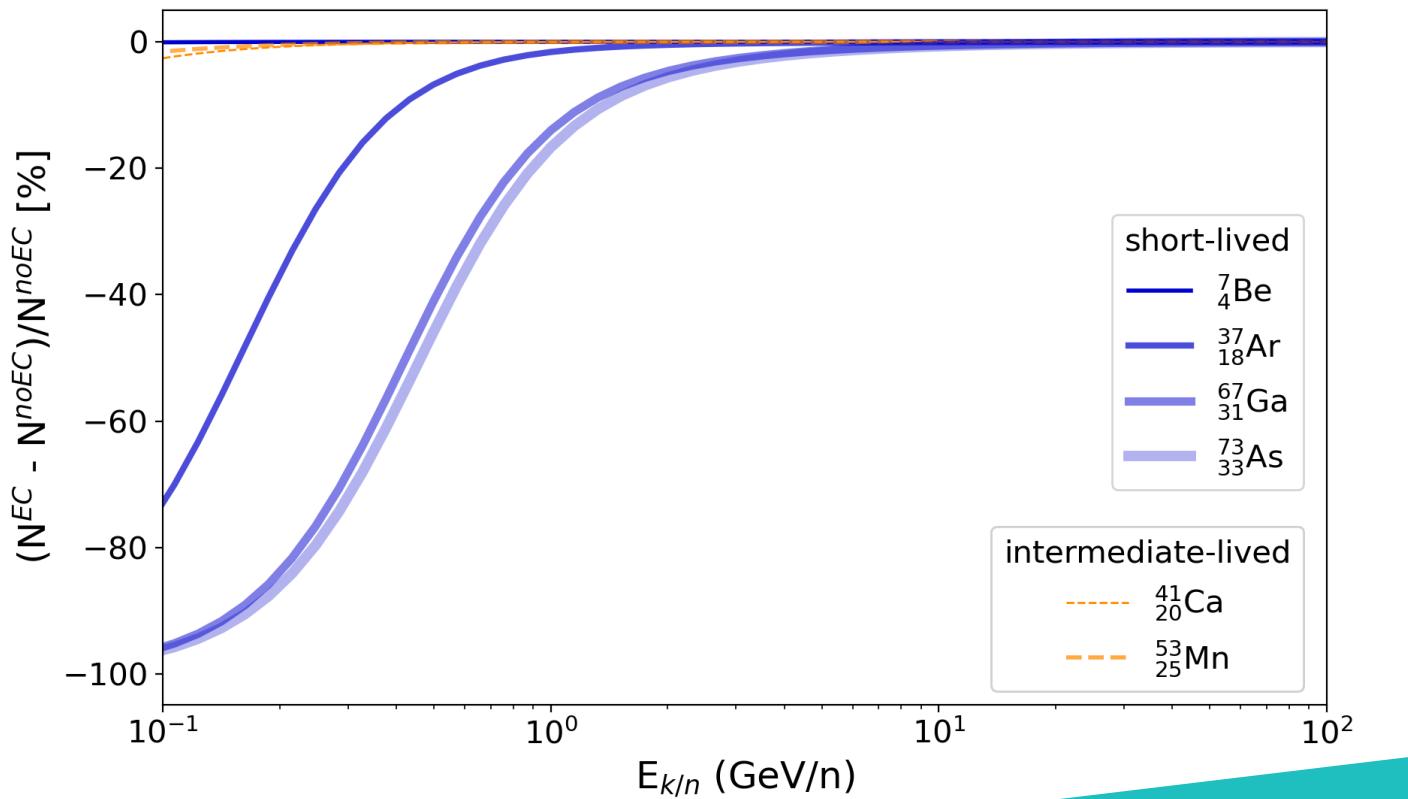
Isotope	$\tau_{1/2}$ (Myr)	Isotopic fraction
$^7_4\text{Be}$	$1.46 \cdot 10^{-7}$	0.55
$^{37}_{18}\text{Ar}$	$9.58 \cdot 10^{-8}$	0.30
$^{41}_{20}\text{Ca}$	$1.00 \cdot 10^{-1}$	0.07
$^{53}_{25}\text{Mn}$	3.70	0.35
$^{67}_{31}\text{Ga}$	$8.93 \cdot 10^{-9}$	0.07
$^{73}_{33}\text{As}$	$2.20 \cdot 10^{-7}$	0.36

# Impact on isotopic fluxes

- Short-lived heavy nuclei all decay at low E
- No effect on intermediate-lived isotopes and for  $E > \text{few GeV/n}$

NB: ACE-CRIS data precision is ~ 50% for Ga and As

Percentage of CRs isotopes that decays by EC

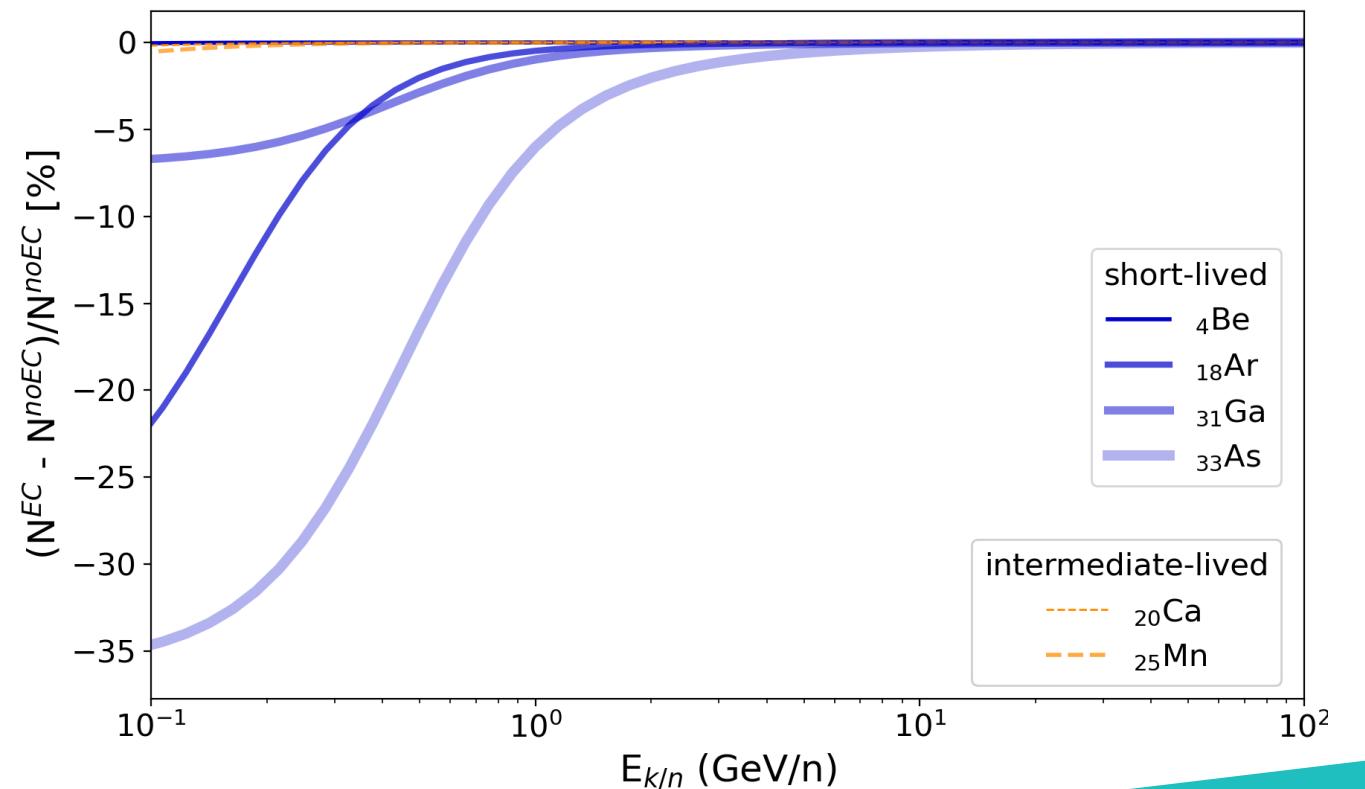


# Impact on elemental fluxes

The impact of EC decay on elemental fluxes is weighted by isotopic abundances

- Short-lived CRs fully decay at low E in Ga and As but not in Ar
- No impact for intermediate-lived nuclei

Percentage of CRs nuclei that decays by EC



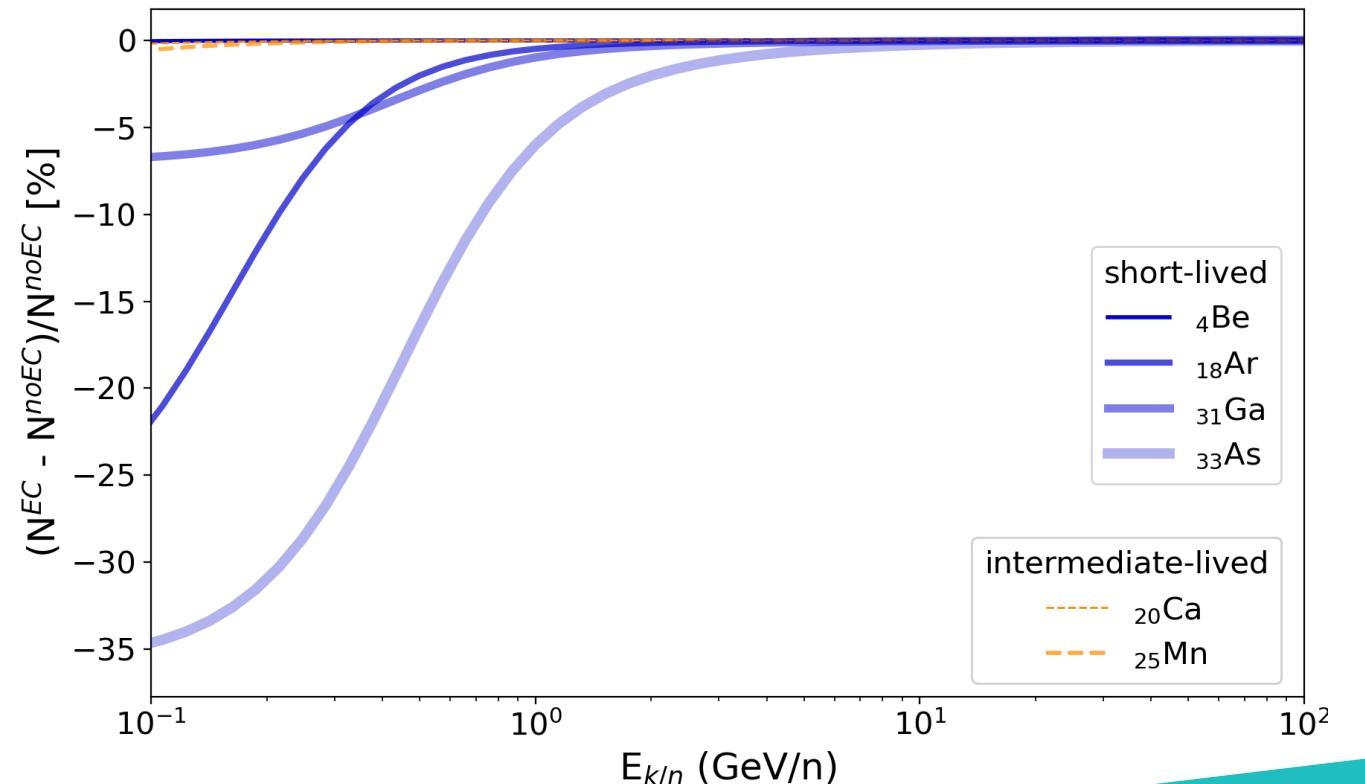
# Impact on elemental fluxes

The impact of EC decay on elemental fluxes is weighted by isotopic abundances

NB: data precision at few hundreds MeV/n is

- Less than 5% (expected) for Ar with AMS-02
- $\sim 7\%$  for Ga and 30% for As with ACE-CRIS

Percentage of CRs nuclei that decays by EC



# Conclusions and perspectives

---

## What we did:

- We computed relevant timescales for GCR fluxes
- We derived solutions for EC decaying isotopes (2-level model)
- We computed the impact of EC decay on isotopic and elemental fluxes

## What we found:

- The net effect of EC decay depends both on  $Z$  and  $\tau_{\text{EC}}$
- Impact on isotopic fluxes  $\sim$  ACE-CRIS precision
- Impact on elemental fluxes slightly larger than AMS-02 precision

# Conclusions and perspectives

---

Still to be done:

- Further improvement of this analytical model
- Implementing EC in the USINE code to account for:
  - energy losses and Solar modulation
  - detailed production of the various isotopes

*That's all folks!*

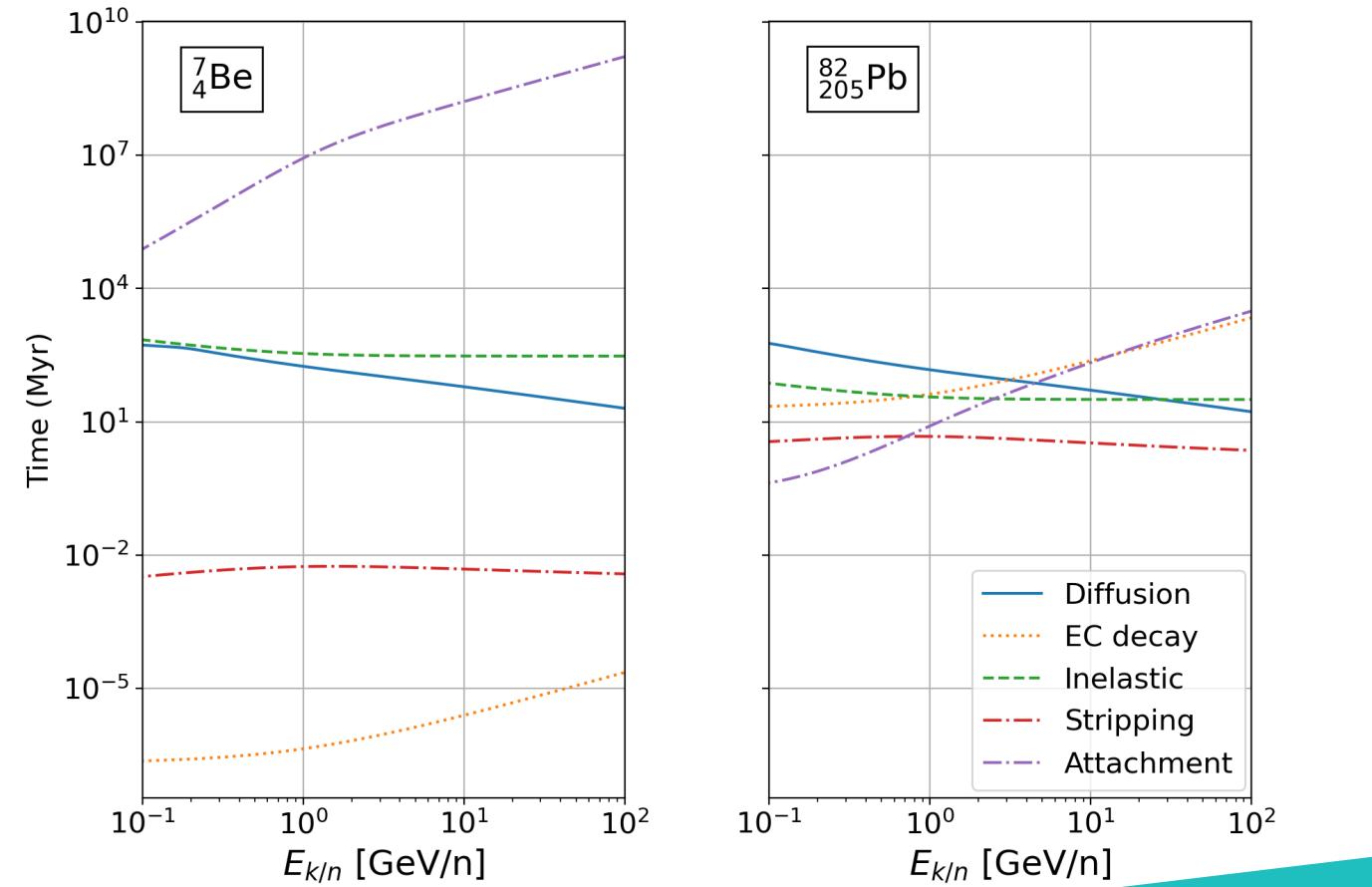
---

# Backup

---

# Characteristic timescales

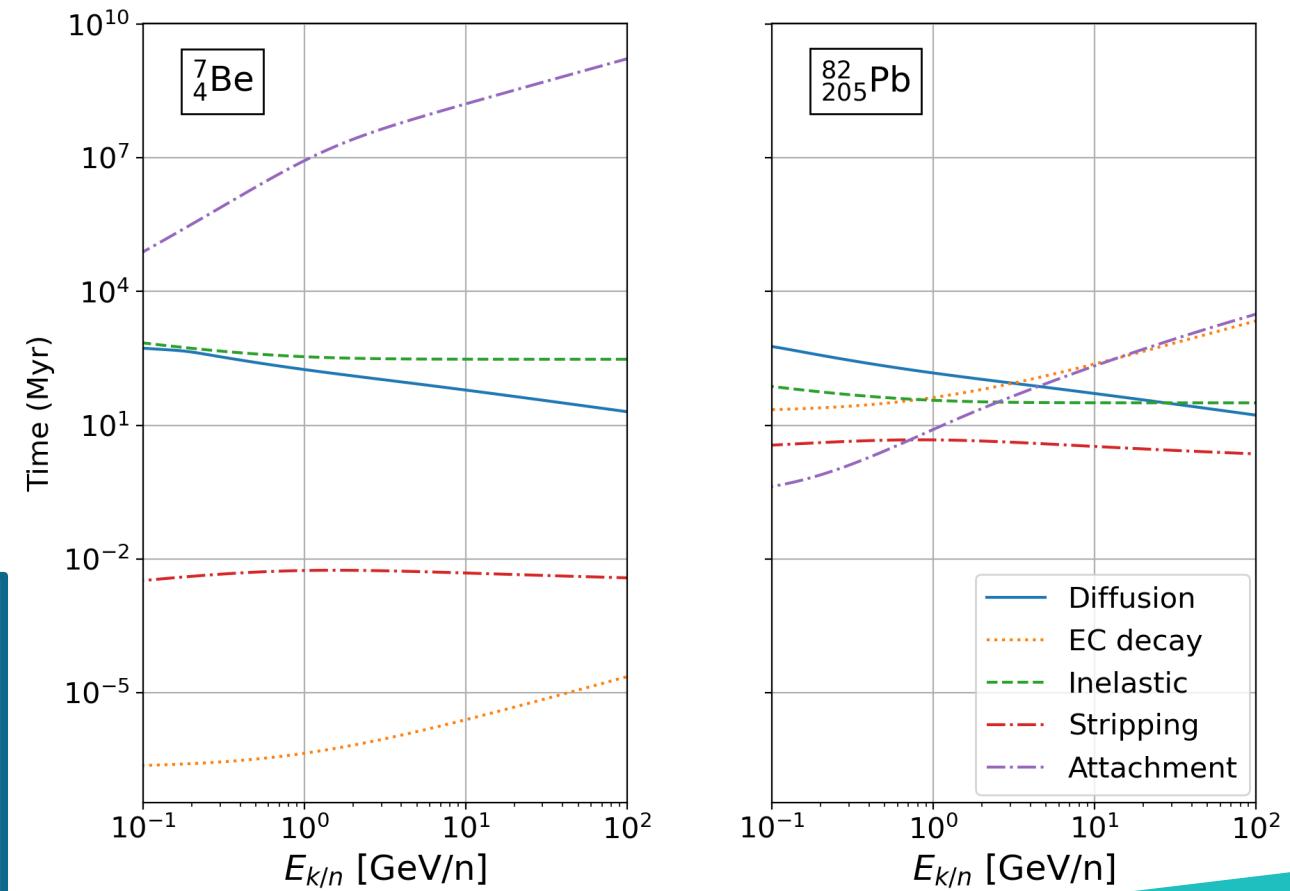
- Diffusion dominates above a few GeV/n
- Attachment more efficient than stripping for large Z



# Characteristic timescales

- Diffusion dominates above a few GeV/n
- $t_{inel}$  smaller than  $t_D$  for low E and large A
- Attachment more efficient than stripping for large Z

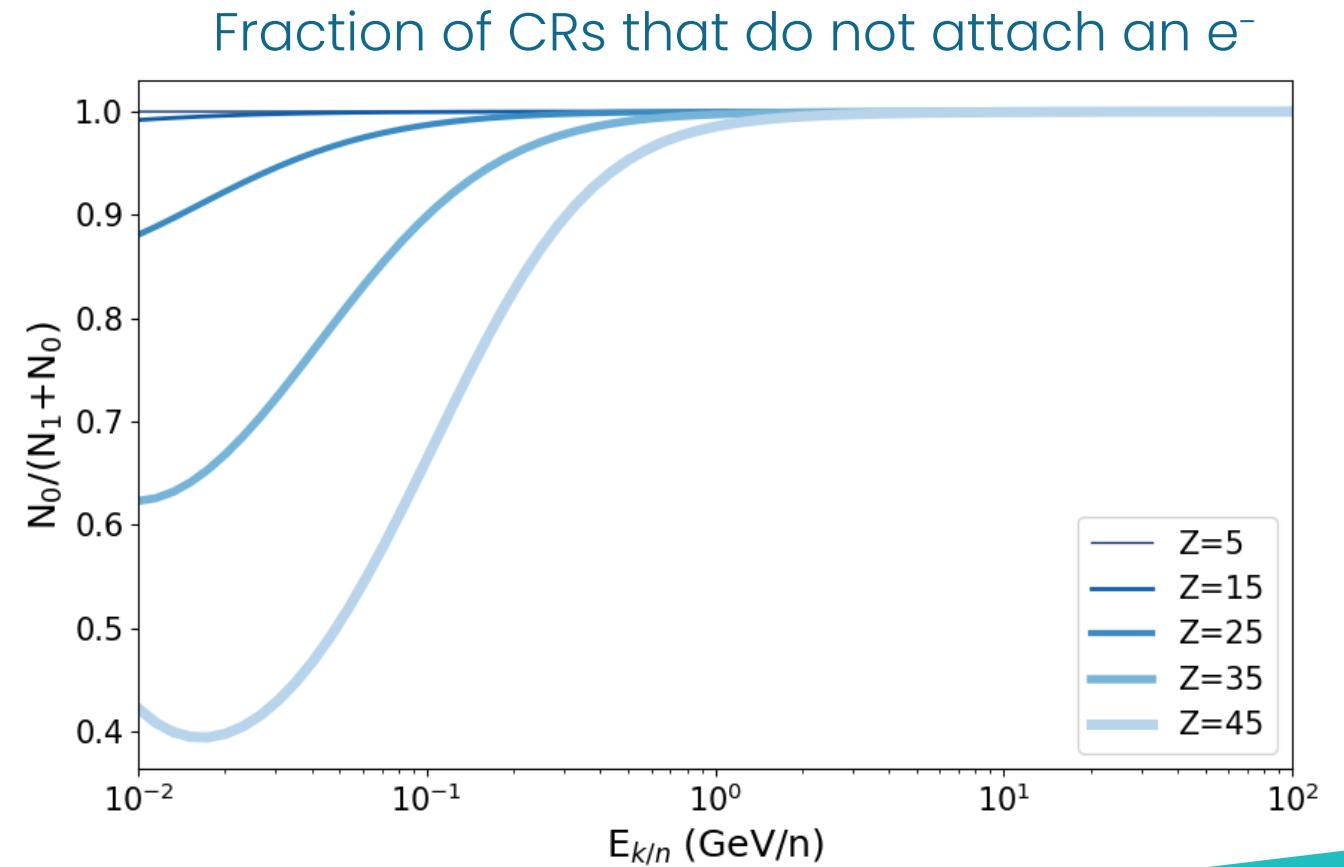
→ Net effect of EC decay depends on interplay between  $t_{att}$ ,  $t_{strip}$  and  $t_{EC}$   
→ EC decay expected to be relevant at low E and large Z



# Attachment and stripping

We want to follow CRs passing from one charge state to the other:

- Stable isotopes:  $\Gamma^{\text{EC}} = 0$
- Charges from 5 to 45



# Attachment and stripping

We want to follow CRs passing from one charge state to the other:

- Stable isotopes:  $\Gamma^{\text{EC}} = 0$
- Charges from 5 to 45

- No  $e^-$  attached above a few  $\text{GeV}/n$
- Heavier CRs attach more  $e^-$  than light CRs

