New collider searches for dark photons

Anh Vu Phan NNV annual meeting, subatomic physics section (Nov 04, 2022) Based on work with Joerg Jaeckel

Radboud University

- Dark photon model
- Displaced vertex at Belle II
- Searches at photon colliders (LUXE and Gamma Factory)

Dark photon model

New collider searches for dark photons



Motivation

- Dark matter should exist.
- Fail to find dark matter?



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- Fail to find dark matter \rightarrow look for portal particles.
- Some BSM models (e.g., string theory) predict additional U(1) symmetries.

Dark photon (X_{μ})



 \rightarrow Behaves like regular photon, but with coupling ϵe to SM fermions and is massive.

For more on dark photon model: Holdom (1986) Fabbrichesi, et al. (2020)

Existing limits on dark photons



Places to search for feebly coupling dark photons

- e^+e^- colliders
 - Displaced vertex search at Belle II ← most promising
- Photon colliders
 - Beam dump experiment at LUXE
 - Bump hunt search at Gamma Factory

Displaced vertex at Belle II

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Belle II

- e^+e^- collision at $\sqrt{s} = 10.58$ GeV.
- Why Belle II?
 - Good statistics
 - Low background
 - Algorithm suitable for finding displaced vertices



Belle II experiment

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Matvienko (QUARKS-2018)



Sensitivity



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Where to find displaced vertices?

Cross-sectional view of Belle II Adapted from Pestotnik (2015)



Region categorization based on

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$\frac{\gamma \text{ conversion (main background)}}{\gamma \text{ from } e^+e^- \rightarrow \gamma\gamma \text{ can convert in matter } \gamma \rightarrow f^+f^-}$



γ conversion (main background) Main background for R > 0.9 cm.



γ conversion (main background)

For 0.2 cm < R < 0.9 cm, background is due to misreconstruction.





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Compute γ **conversion cross section** $\gamma(q) + N(p) \rightarrow f^{-}(l_{-}) + f^{+}(l_{+}) + N(p')$



Suppressed

Background estimation



Phase space coverage



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Searches at photon colliders

New collider searches for dark photons



LUXE

• Proposed $e^-\gamma$ collider at DESY, Hamburg

Adapted from Abramowicz *et al.* (2019)



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For more on LUXE: Abramowicz et al (2021)

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LUXE

• $e^-\gamma$ collider



Beam dump at LUXE



Beam dump at LUXE



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ODEL > DISPLACED VERTEX AT BELLE II > SEARCHES AT PHOTON COLLIDERS

Gamma Factory proposed at CERN





Bump hunt at Gamma Factory



Bump hunt at Gamma Factory



Adapted from image by Fermilab

Bump hunt at Gamma Factory



Regions above the lines are covered by Gamma Factory at 95% C.L. Assume integrated luminosity of 10⁴⁰ cm⁻²

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THANK YOU FOR LISTENING

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Signal

- Signal: photon + displaced e^+e^- pair.
- At high energy, $\frac{d\sigma_{e^+e^- \to \gamma X}}{d\cos\theta} = \frac{2\pi\alpha^2\chi^2}{s} \left(\frac{1+\cos^2\theta}{\sin^2\theta}\right).$
- Decay length (N_f : number of possible decay products) $\lambda_X = \frac{\hbar c E_X}{m_X \Gamma_X}$, with $\Gamma_X = \frac{1}{3} \alpha \chi^2 m_X N_f$.
- Number of signal events $S = LT\sigma_{e^+e^- \to \gamma X} (e^{-r_{min}/\lambda_X} - e^{-r_{max}/\lambda_X}).$
- For Belle II, we search inside beam pipe (i.e. $r_{min} = 0.1 \ cm$, $r_{max} = 0.9 \ cm$).



Background: $e^+ e^- \rightarrow \gamma \gamma$, $\gamma + matter \rightarrow e^+ e^-$

• The number of photon conversion to e^+e^- satisfying our cuts is



Background: $e^+ e^- \rightarrow \gamma \gamma$, $\gamma + matter \rightarrow e^+ e^-$

• For many layers of matter, conversion probability is

$$p_{\gamma \to e^+e^-} = \frac{dN_{\gamma \to e^+e^-}}{N_0} = \frac{d\sigma_1}{\sigma_1} (1 - e^{-n_1\sigma_1z_1}) + e^{-n_1\sigma_1z_1} \frac{d\sigma_2}{\sigma_2} (1 - e^{-n_2\sigma_2z_2}) + \cdots$$
Probability γ converts in layer 1
Probability γ converts in layer 1
Probability γ converts in layer 2
$$n_1, z_1 \quad n_3, z_3$$

$$\rho = e^+$$

$$n_2, Z_2$$

Compute γ **conversion cross section** $\gamma(q) + N(p) \rightarrow f^{-}(l_{-}) + f^{+}(l_{+}) + N(p')$

• Amplitudes

$$\mathcal{M}_{BH} = \frac{ie}{t} \varepsilon_{\mu}(q) \,\overline{N}(p') F_{1}(t) \gamma_{\nu} N(p) \, T^{\mu\nu}_{\gamma^{*}\gamma \rightarrow f\bar{f}}$$

$$t = (p - p')^{2} \qquad \text{Nucleus form} \qquad \text{For } f = \text{leptons, } T^{\mu\nu}_{\gamma^{*}\gamma \rightarrow f\bar{f}} \text{ is amplitude for } \gamma\gamma \rightarrow l^{+}l^{-}$$



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$$\mathcal{M}_{TCS} = \frac{ie^{3}}{m_{ff}^{2}} J_{\nu}(q') \varepsilon_{\mu}(q) \,\overline{N}(p') \left(\frac{1}{\alpha} T_{fTCS}^{\mu\nu}\right) N(p) \qquad \text{Gryniul}$$

$$m_{ff}^{2} = (l_{+} + l_{-})^{2} \qquad q' = l_{+} + l_{-}$$



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$$T_{fTCS}^{\mu\nu} \approx \left(-g^{\mu\nu} + \frac{q'^{\mu}q^{\nu}}{q \cdot q'}\right) f(\nu) \qquad \text{Gryniuk (2020)}$$

$$\stackrel{\text{Incoming photon energy}}{\bigwedge} \qquad \text{Near-real, near forward} \qquad \text{Spin-averaged forward} \\ \lim(m_{ff}^{2})|t| \ll s) \qquad \text{Spin-averaged forward} \\ \operatorname{real Compton amplitude} \qquad \text{Timelike Compton scattering (TCS) diagram}$$

N

Compute γ **conversion cross section** $\gamma(q) + N(p) \rightarrow f^{-}(l_{-}) + f^{+}(l_{+}) + N(p')$

• Amplitudes

$$\mathcal{M}_{BH} = \frac{ie}{t} \varepsilon_{\mu}(q) \,\overline{N}(p') F_{1}(t) \gamma_{\nu} N(p) \, T_{\gamma^{*}\gamma \to f\bar{f}}^{\mu\nu}$$
$$\mathcal{M}_{TCS} = \frac{ie^{3}}{m_{ff}^{2}} J_{\nu}(q') \varepsilon_{\mu}(q) \,\overline{N}(p') \left(\frac{1}{\alpha} T_{fTCS}^{\mu\nu}\right) N(p)$$
$$T_{fTCS}^{\mu\nu} \approx \left(-g^{\mu\nu} + \frac{q'^{\mu}q^{\nu}}{q \cdot q'}\right) f(\nu)$$

Total unpolarized photoabsorption cross section of nucleus N.

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Gryniuk, Hagelstein, and Pascalutsa (2015) 0

ptical theorem
$$\Rightarrow \operatorname{Im} f(\nu) = \frac{\nu}{4\pi} \sigma_{\gamma N}(\nu)$$

Analyticity $\Rightarrow \operatorname{Re} f(\nu) = -\frac{Z^2 \alpha}{\frac{M_N}{M_N}} + \frac{\nu^2}{2\pi^2} \mathcal{P} \int_0^\infty d\nu' \frac{\sigma_{\gamma N}(\nu')}{\nu'^2 - \nu^2}$

Nucleus N with mass M_N and atomic number Z