

Theory of heavy ions in the LHC era

Towards a precision analysis of heavy ion collisions

Based on *Trajectum* with Govert Nijs

2010.15130, 2010.15134 with Govert Nijs, Umut Gursoy and Raimond Snellings

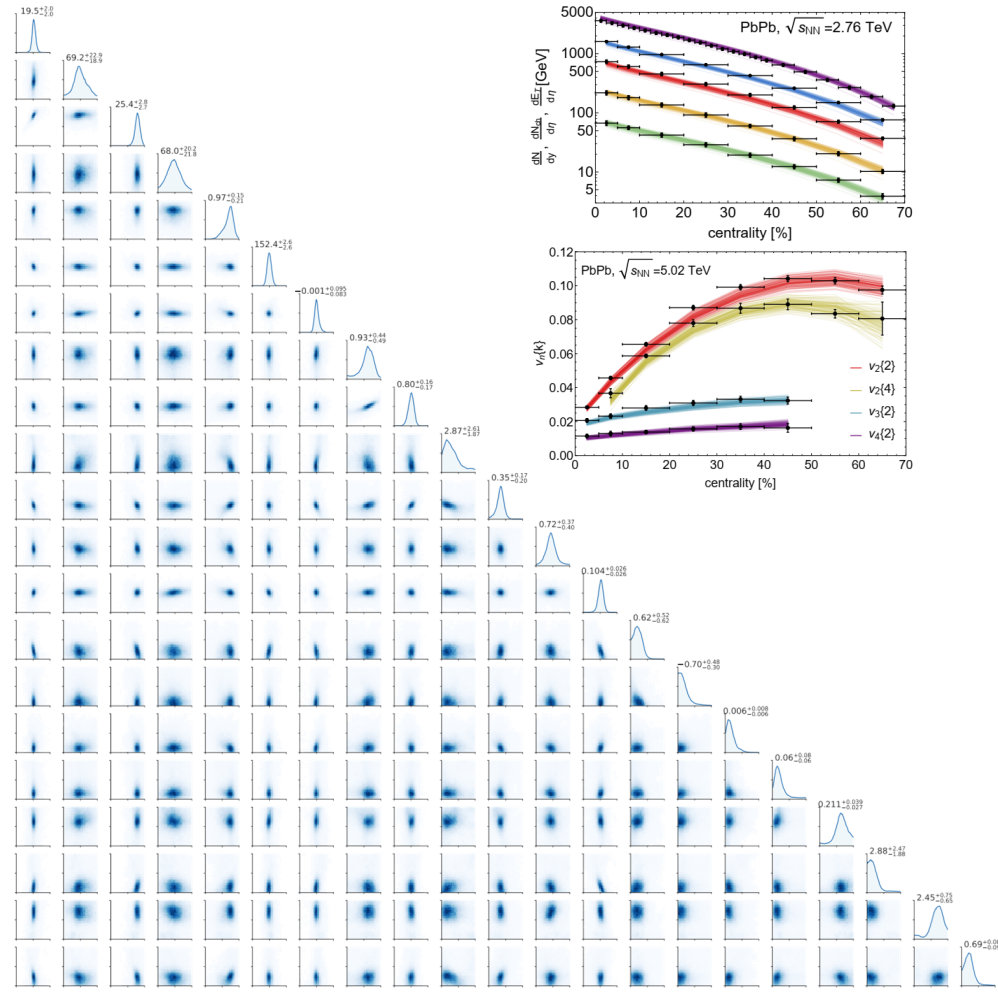


Roman excavations in Utrecht (from *Trajectum*, or bridge) in 1929

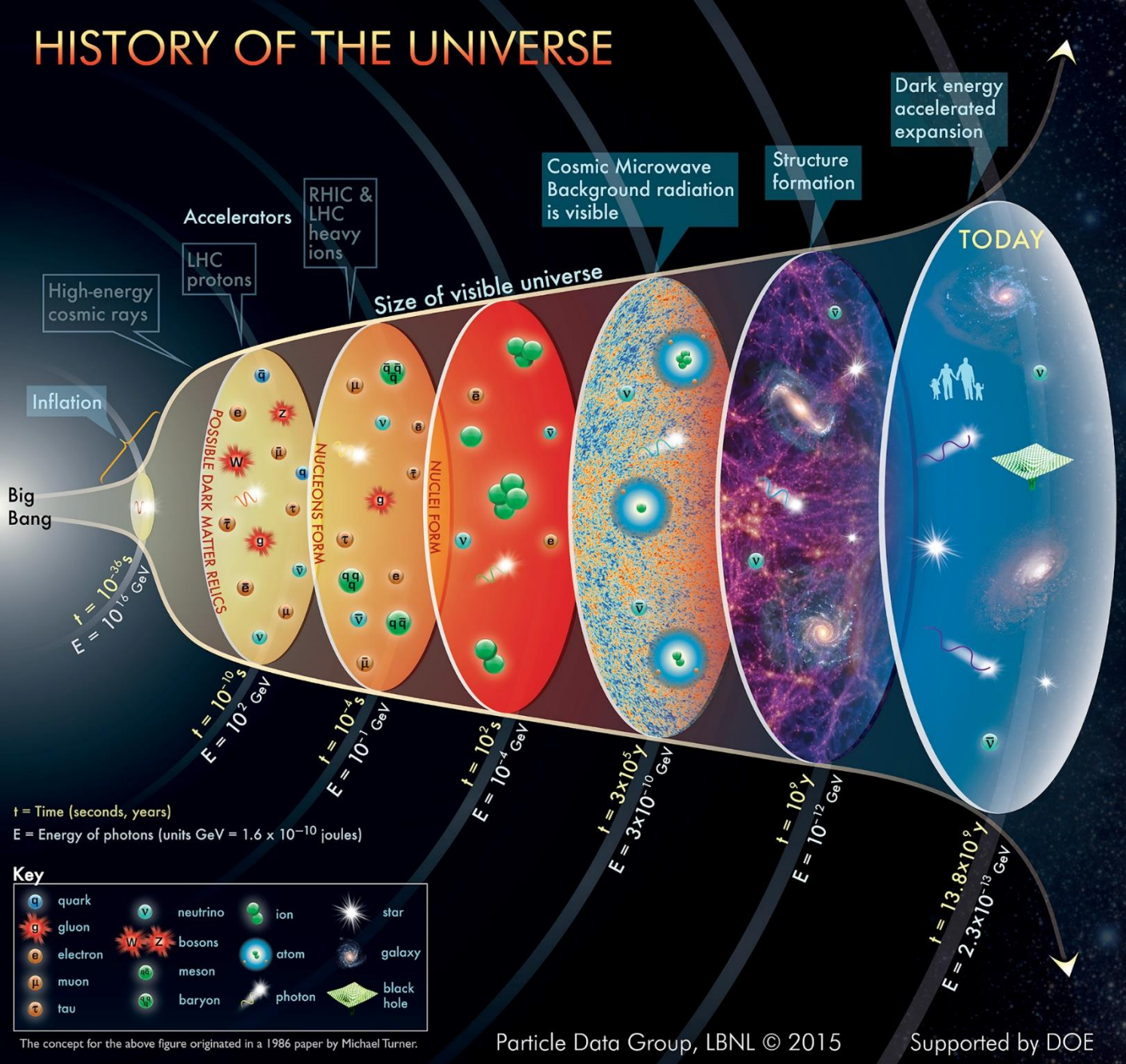
Wilke van der Schee

NIKHEF colloquium, Amsterdam

5 March 2021



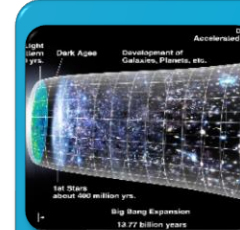
HISTORY OF THE UNIVERSE



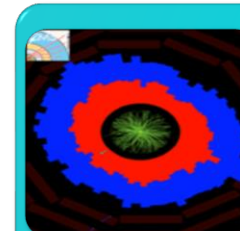
Quark-gluon plasma (QGP)



Quantum-Chromodynamics (QCD)
A fundamental force of nature



Recreating the big bang
At age $1 \mu s$ the entire universe was QGP!



QGP turns out interesting
Strongly coupled quantum matter

The concept for the above figure originated in a 1986 paper by Michael Turner.

The QCD phase diagram

Strong coupling: first principle only from lattice QCD

- Smooth cross-over from confined hadron gas to deconfined QGP

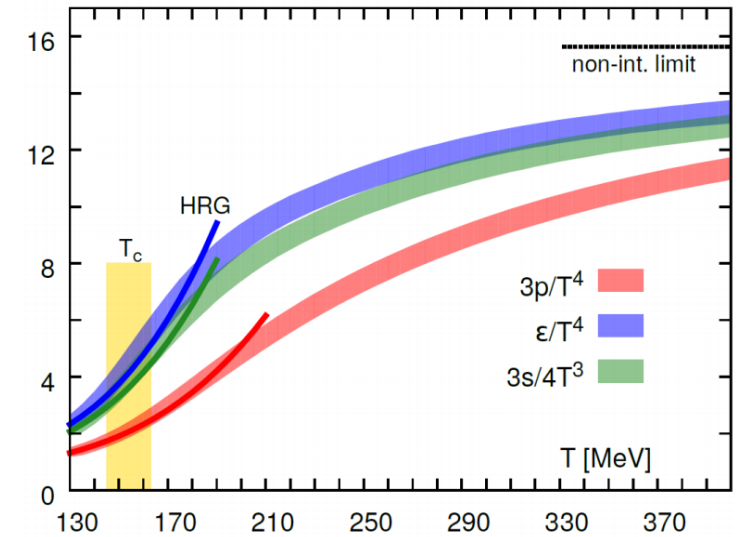
Sign problem: only Euclidean

- Problematic to study baryon chemical potential (neutron stars)
- Problematic to study real-time dynamics (shear viscosity)

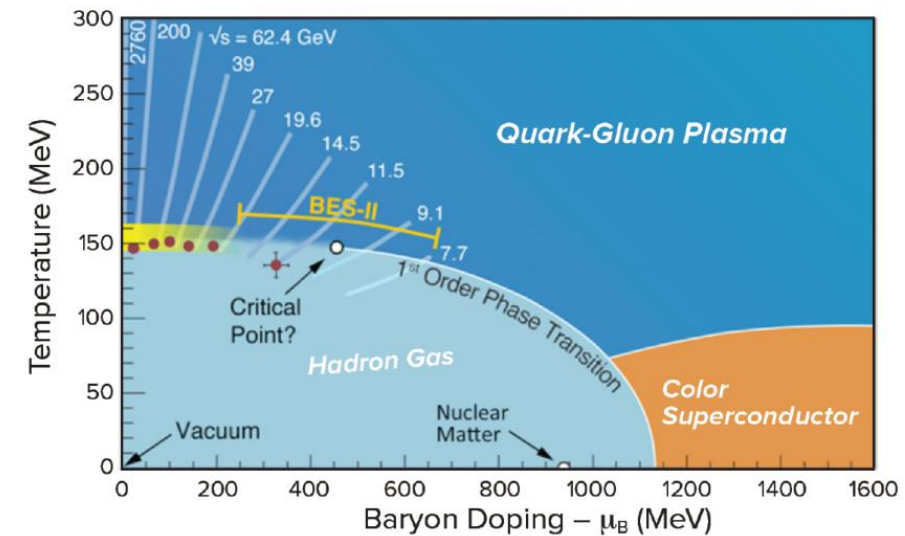
A conjectured critical point in reach of RHIC energies?

- LHC does not reach high enough baryon number densities

Lattice equation of state



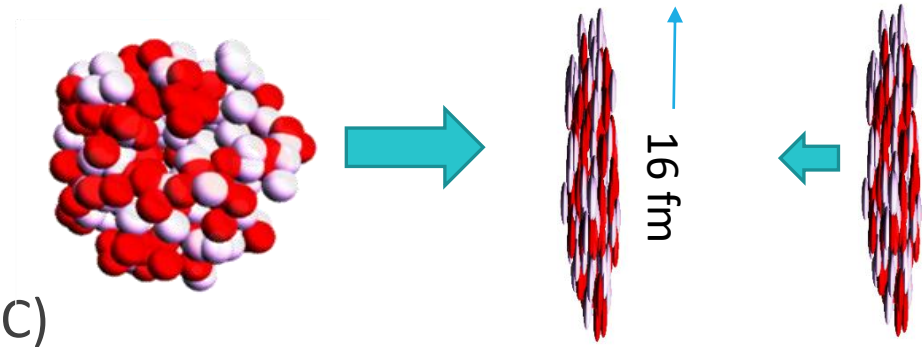
Phase diagram



How to create QGP

Colliding heavy nuclei (Pb, Au) at high energies

Lorentz gamma factor up to 2500 (LHC) or 100 (RHIC)

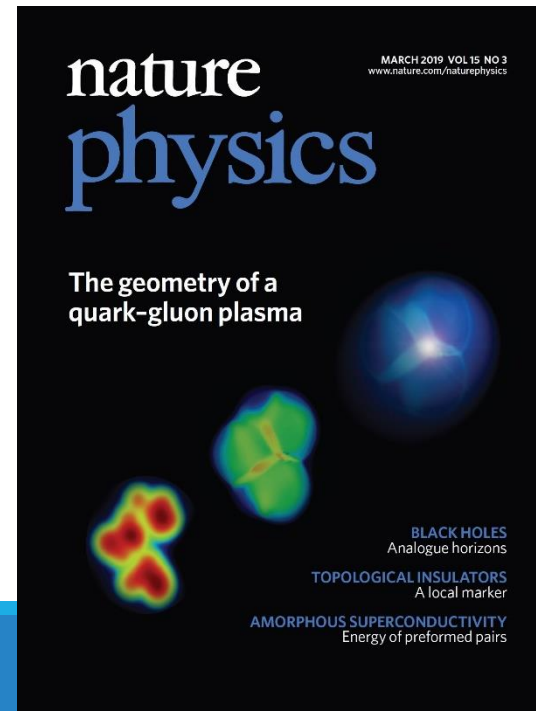
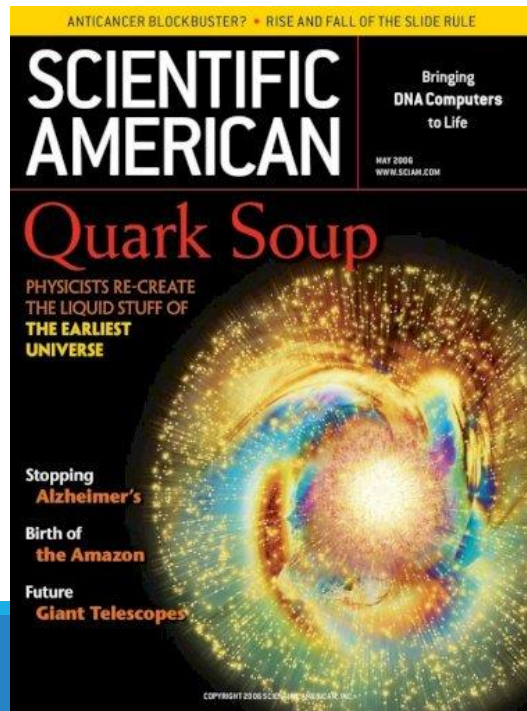


Hottest fluid:
 10^{12} K

Smallest fluid:
 ~ 2 fm living 10^{-23} s

Most perfect/strange:
 $\eta/s \sim 0.08$

Most vortical fluid:
 $\omega \sim 10^{22}/s$



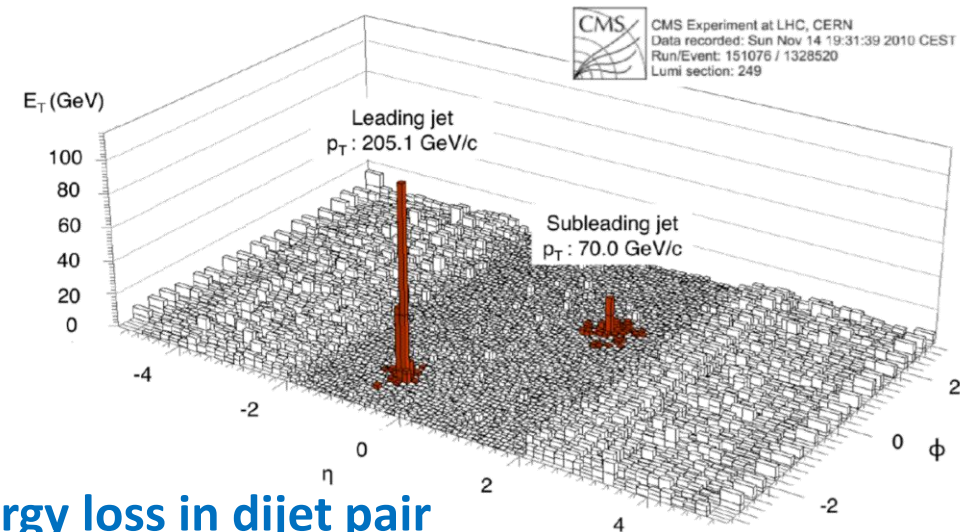
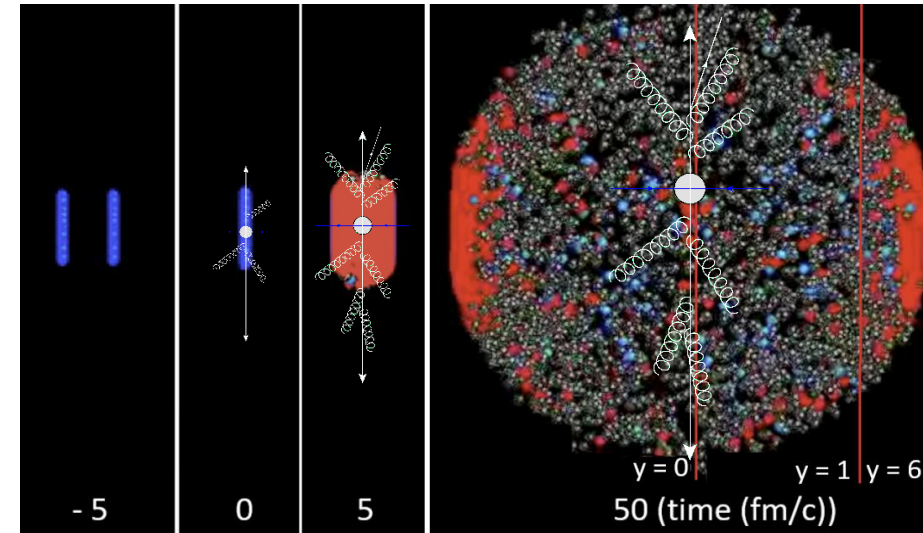
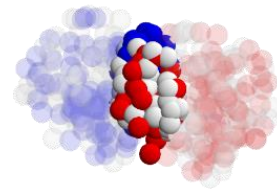
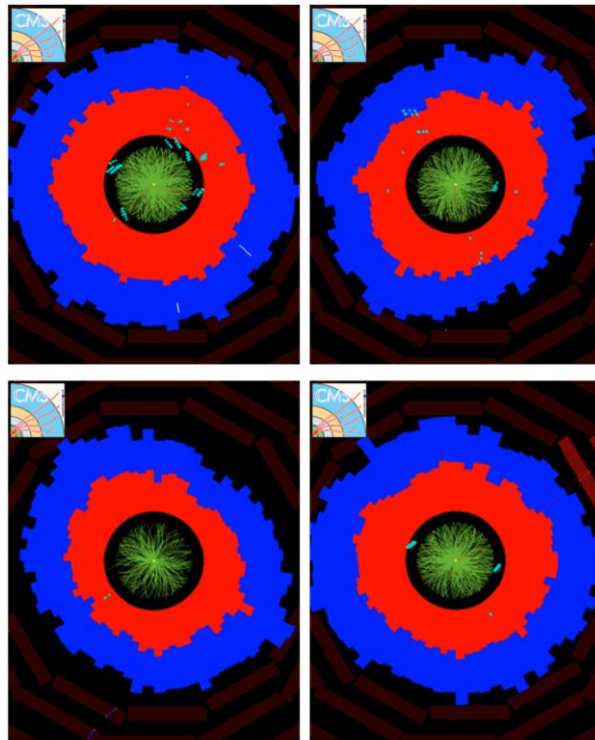
CERN accelerator complex



Quark-gluon plasma is strongly coupled

Initial stage - QGP - hadronic phase

Anisotropic flow (small viscosity)



Jet energy loss in dijet pair

Strangeness: from pQCD to thermal

1. Ratio of strange baryons versus pions

- Pythia fits low multiplicity
- But constant towards higher multiplicity (!)

Thermodynamical string fragmentation



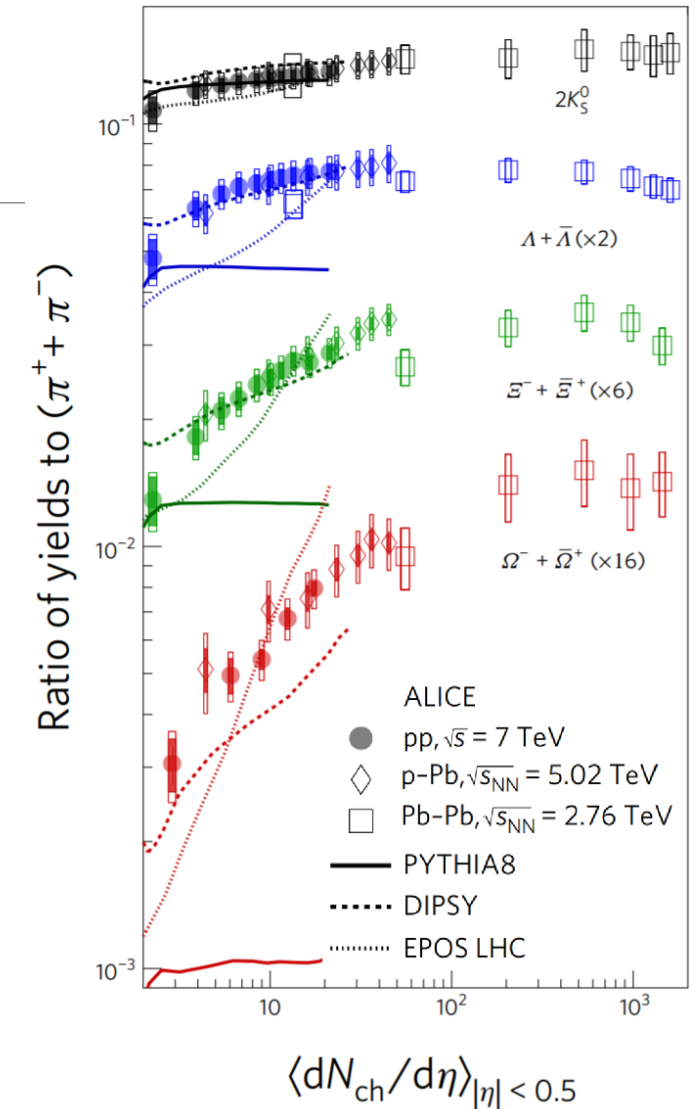
Nadine Fischer^{a,b} and Torbjörn Sjöstrand^a

January 31, 2017

ABSTRACT: The observation of heavy-ion-like behaviour in pp collisions at the LHC suggests that more physics mechanisms are at play than traditionally assumed.

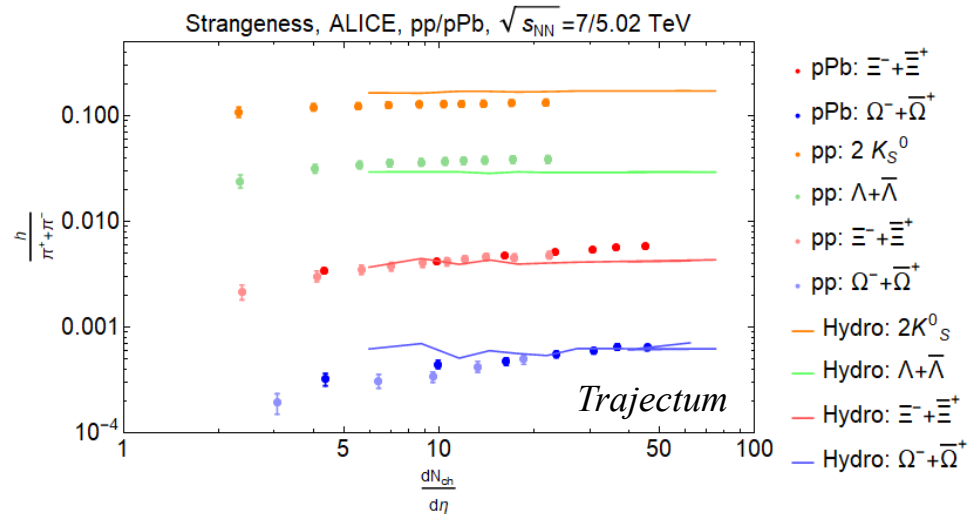
2. Saturates for high multiplicity pPb / PbPb

- Interpretation: thermal strangeness production



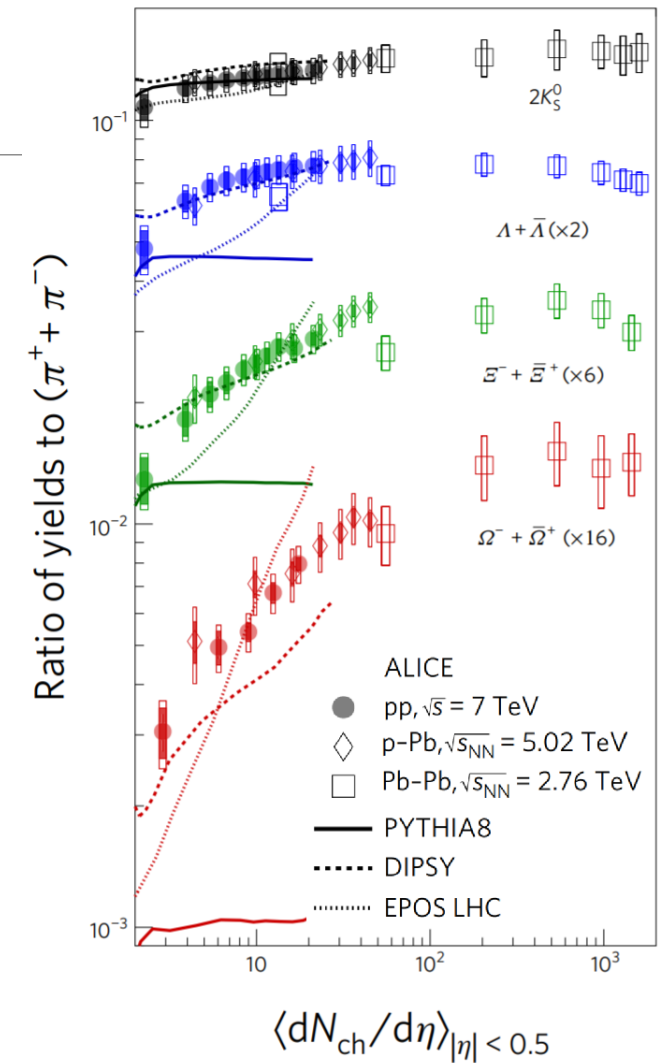
Strangeness: from pQCD to thermal

1. Hydro+hadronic cascade, one parameter ($T_{\text{partic.}}$):



2. Hydro has only small dependence on N_{ch}

- Approximately fits thermal model

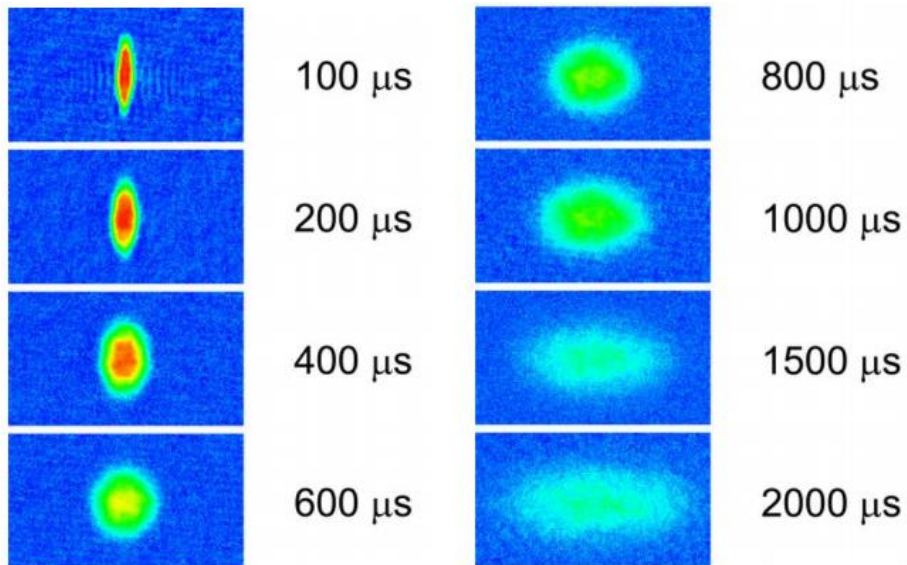


The most perfect liquid?

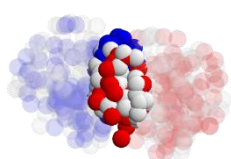
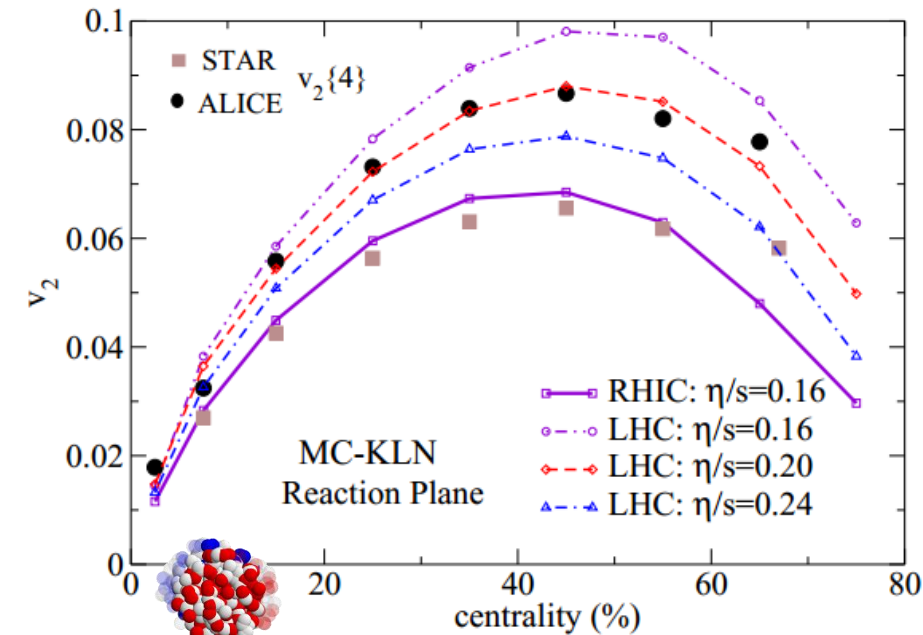
$$\frac{d\bar{N}}{d\varphi} = \frac{\bar{N}}{2\pi} \left(1 + 2 \sum_{n=1}^{\infty} \bar{v}_n \cos(n(\varphi - \bar{\Psi}_n)) \right)$$

Famous viscosity, AdS/CFT or **holography**: $\frac{\eta}{s} = \frac{1}{4\pi} \approx 0.08$

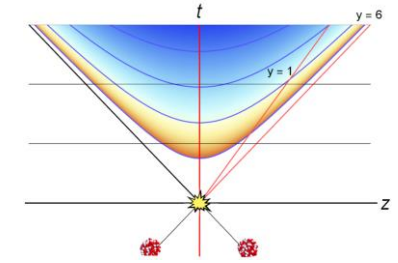
Fermions at unitarity (cold 😊)



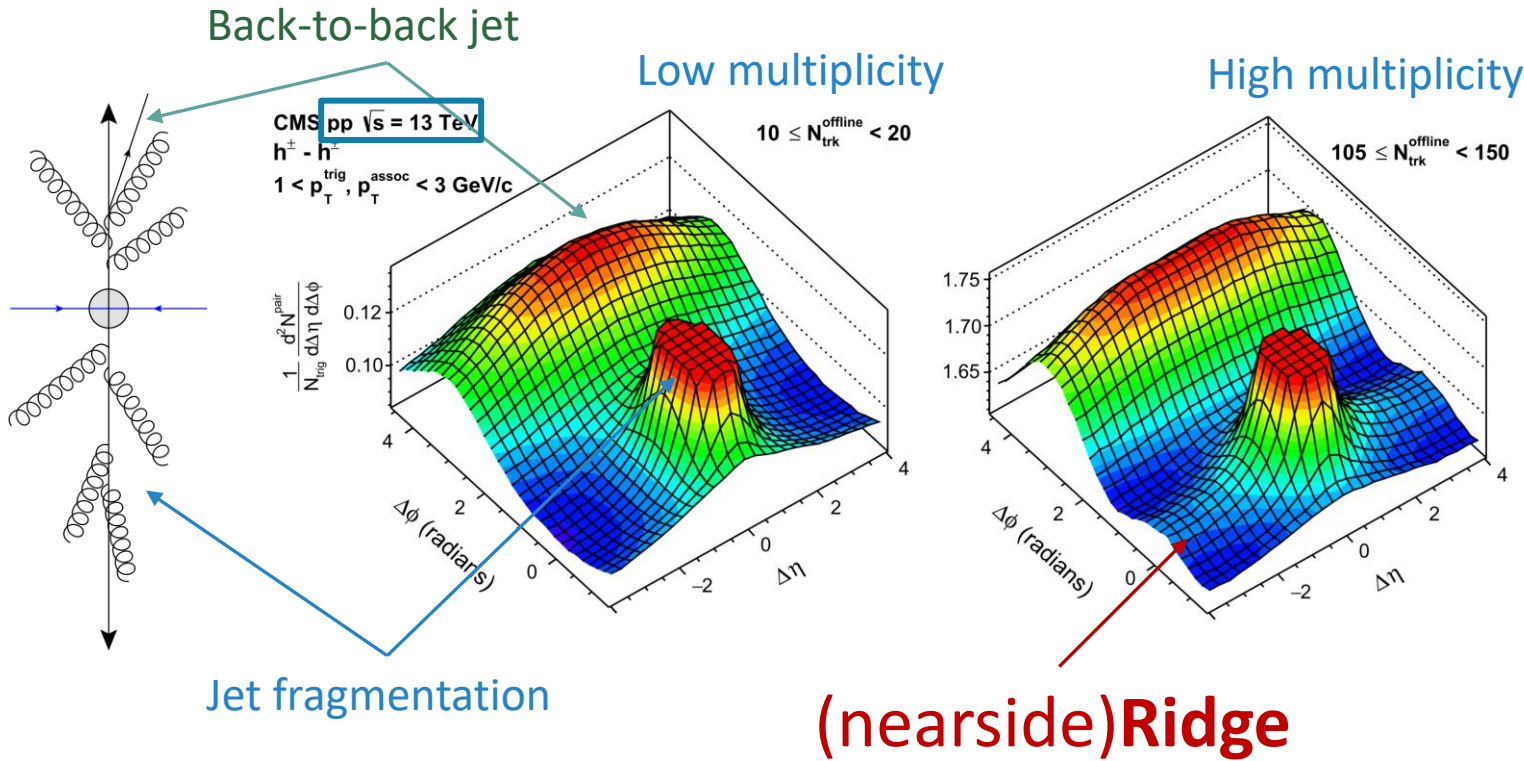
Quark-gluon plasma (hot 😊)



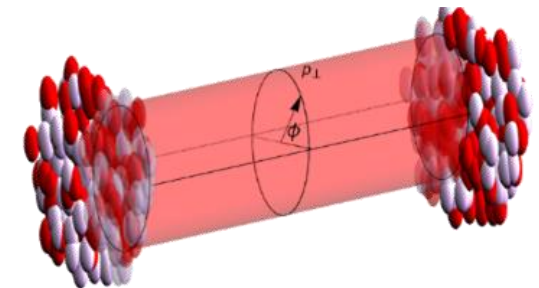
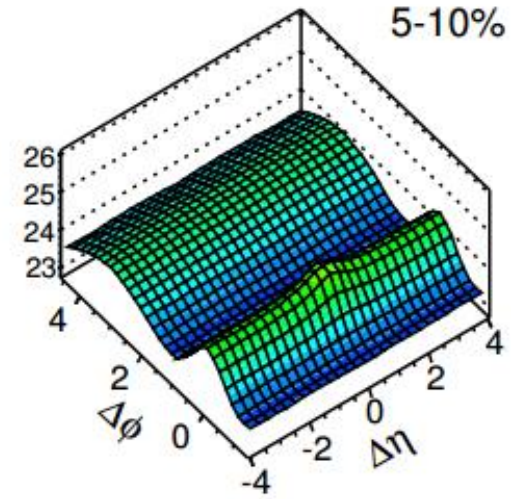
Ridges everywhere: *panta rei*



1. Ridge at $\Delta\phi=0$ and large $\Delta\eta$: an initial or geometric effect

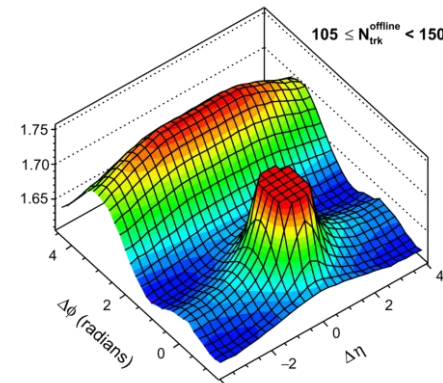
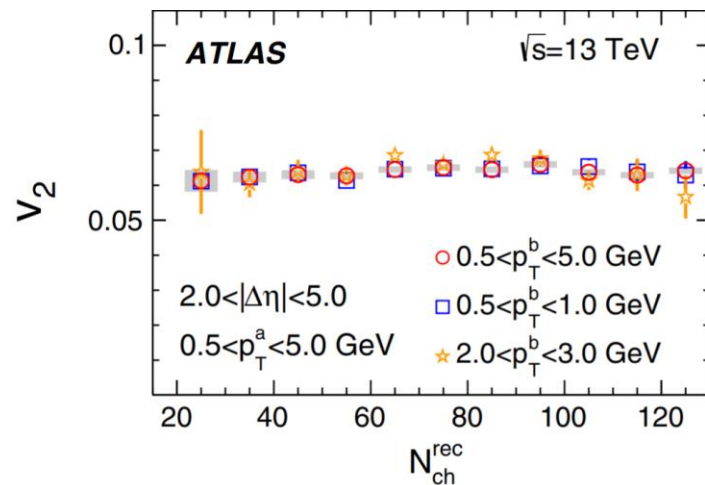


PbPb $\sqrt{s_{\text{NN}}} = 2.76$ TeV



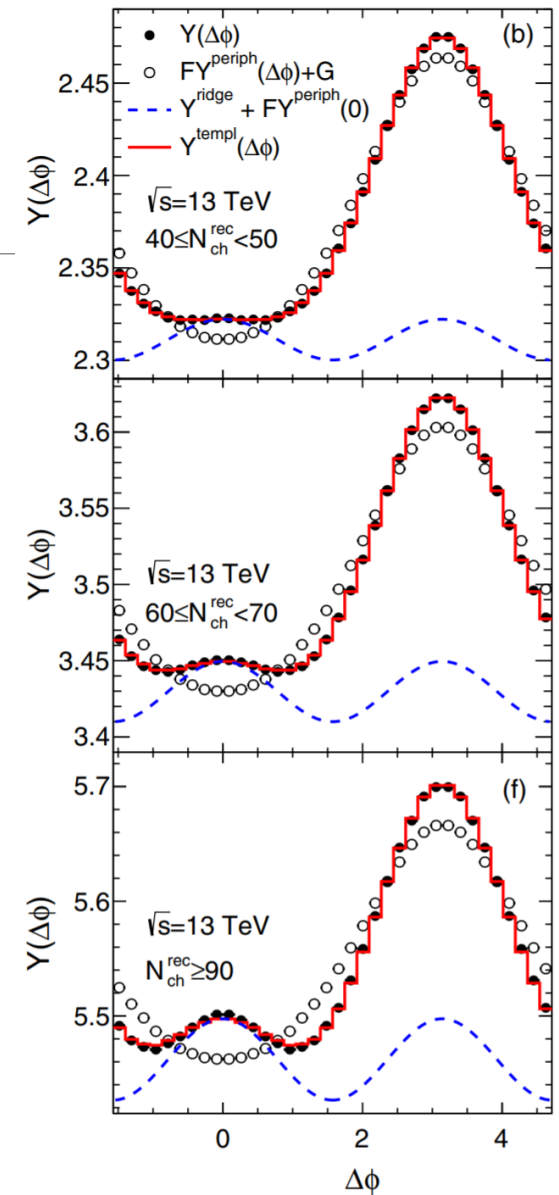
Extract Fourier harmonics of the ridge

1. Essential to split ridge in 'hard' and 'soft' part
2. Template fit allows extrapolation down to $N^{\text{rec}} < 20$
3. Soft v_2 essentially constant versus multiplicity:
 - QGP-like physics in pp collisions?

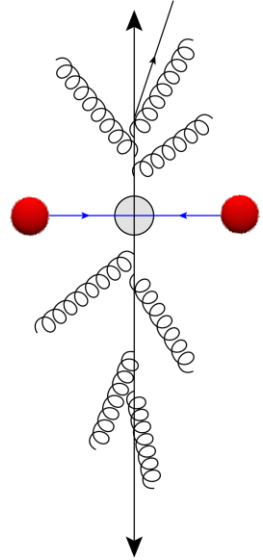


$$Y^{\text{templ}}(\Delta\phi) = FY^{\text{periph}}(\Delta\phi) + Y^{\text{ridge}}(\Delta\phi)$$

$$Y^{\text{ridge}}(\Delta\phi) = G[1 + 2v_{2,2} \cos(2\Delta\phi)]$$

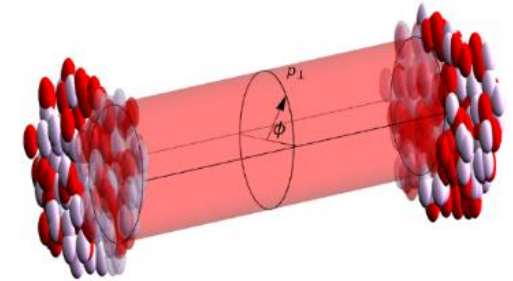


(high energy) ¿HEP versus HIP? (heavy ion)

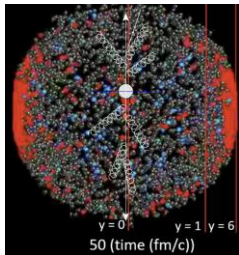


Low multiplicity
Jet-like particle shower
No equilibration

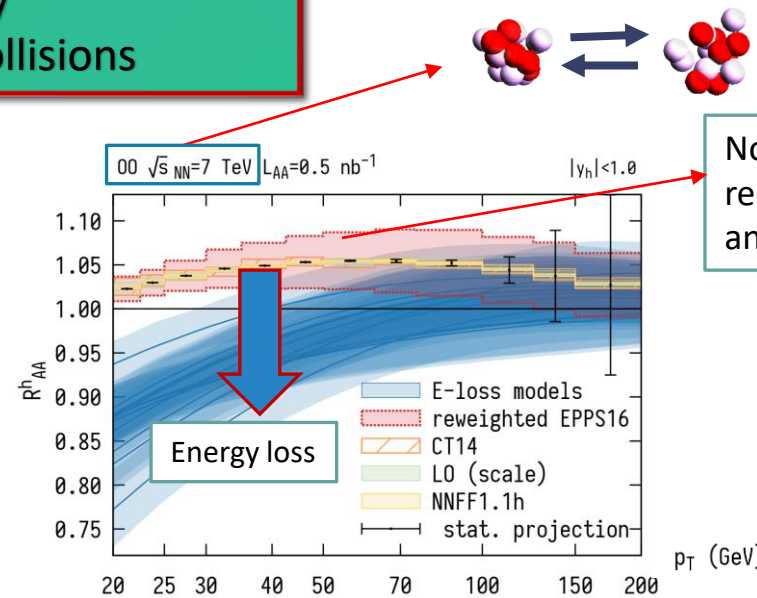
High multiplicity
Relatively few jets
Equilibration: QGP



Jets important in heavy ion/small systems
Often intermediate multiplicity
QGP-type physics part of pp collisions



OO collisions as an example:
Nuclear modification factor: hadron R_{AA}
More energy loss \rightarrow fewer hadrons
Interplay from HEP and HIP



Can heavy ions be understood from a non-Abelian gauge theory?

1. For low p_T particles: is QGP just sensitive to the thermal sector of QCD?
 - And can we compute **and** measure its **fundamentals**? EOS & viscosities?

2. Towards high p_T and smaller systems:
 - Significant jet-like correlations: non-flow
 - Either suppress (large $\Delta\eta$) or use as a probe (hard probes)

Is QGP strongly coupled?
At which energy scale?
Non-conformal?

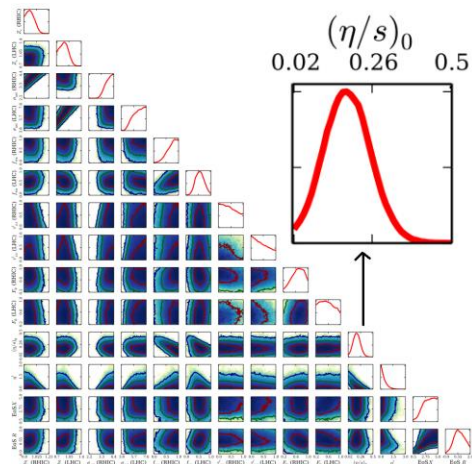
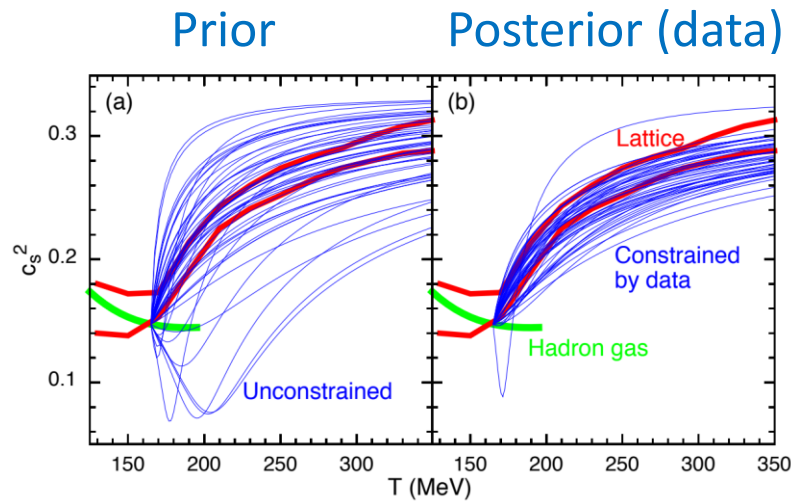
Hydro at large gradients

3. Precise questions to *unravel the fundamentals of QCD*
 - T-dependent shear viscosity, bulk viscosity, second order transport
 - Fast thermalisation? 0.1 fm/c or 1.5 fm/c?
 - Particle ratios: (sizeable) deviations from thermal equilibrium
 - Initial shape: how to convert colliding nucleons to energy density, structure of a proton

What are the d.o.f.?
Partons? Glasma?

First global analyses

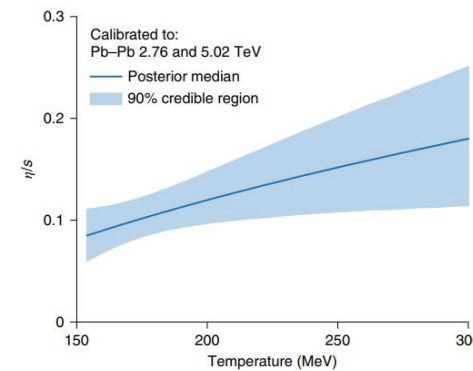
Constraining EOS (Jan 2015)



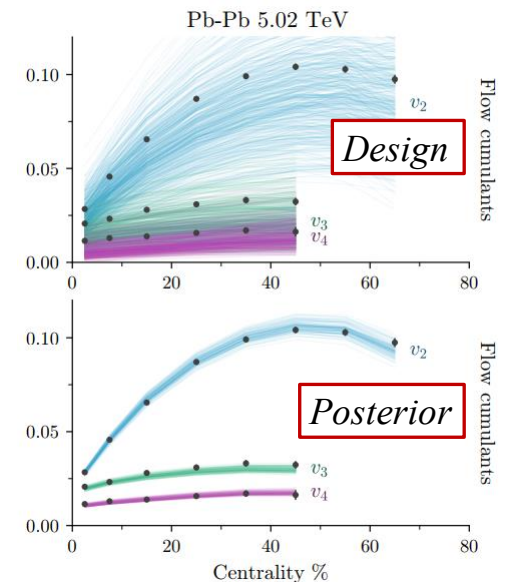
Precise questions require precise understanding of interplay of rich physics in heavy ion collisions

Constraining η/s (2019, Nature Physics)

η/s versus temperature Posterior (η/s +slope)

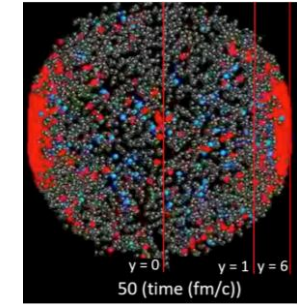
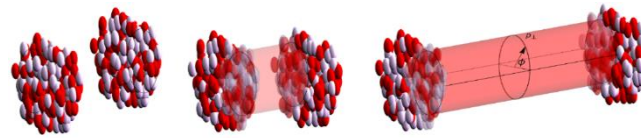


Jonah E. Bernhard, J. Scott Moreland and Steffen A. Bass
Bayesian estimation of the specific shear and bulk viscosity of QGP



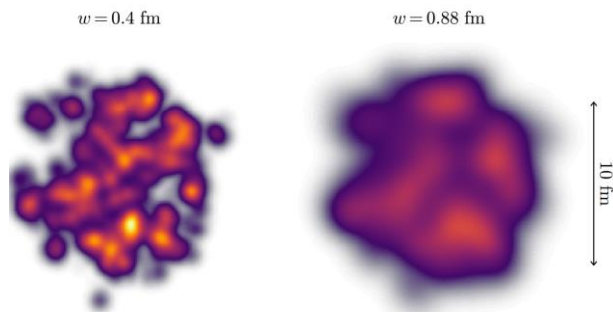
Important: pioneering studies that only included p_T -integrated observables

Standard model of heavy ion collisions



Initial stage (9)

Subnucleonic structure? (7)

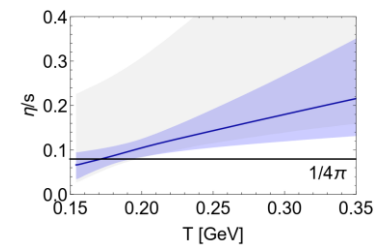


Non-thermal flow? (2)
for time τ with *varying speed (new)*

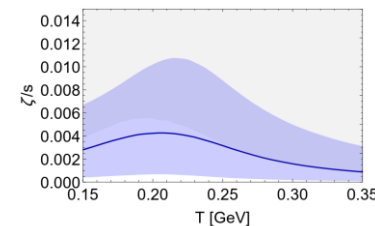
Fluctuations? (1)

Viscous hydrodynamics (9)

Shear viscosity (3)



Bulk viscosity (3)



Second order transports: 3 (*new*)

Cascade of hadrons (1)

Convert quark-gluon plasma at T_{switch} to particles following Boltzmann distribution (particization, 1)

Subtle: viscous corrections

Evolve particles with hadronic code: SMASH

Hydrodynamics: first and second order

1. Constitutive relations for the stress tensor, with $p(\rho)$ EOS from HotQCD

$$T^{\mu\nu} = \rho u^\mu u^\nu - (p + \Pi)\Delta^{\mu\nu} + \pi^{\mu\nu}, \quad \Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$$

With shear and bulk tensors:

$$D\Pi = -\frac{1}{\tau_\Pi} [\Pi + \zeta \nabla \cdot u + \delta_{\Pi\Pi} \nabla \cdot u \Pi - \lambda_{\Pi\pi} \pi^{\mu\nu} \sigma_{\mu\nu}],$$

$$\Delta_\alpha^\mu \Delta_\beta^\nu D\pi^{\alpha\beta} = -\frac{1}{\tau_\pi} \left[\pi^{\mu\nu} - 2\eta \sigma^{\mu\nu} + \delta_{\pi\pi} \pi^{\mu\nu} \nabla \cdot u - \phi_7 \pi_\alpha^{\langle\mu} \pi^{\nu\rangle\alpha} + \tau_{\pi\pi} \pi_\alpha^{\langle\mu} \sigma^{\nu\rangle\alpha} - \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu} \right].$$

2. We vary the green coefficients, η and ζ as a function of temperature,

2nd order according to $\frac{\tau_{\Pi S} T \delta^2}{\zeta}$, $\frac{\tau_\pi S T}{\eta}$ and $\frac{\tau_{\pi\pi}}{\tau_\pi}$

Performing a global analysis

Bayes theorem:

$$\mathcal{P}(\mathbf{x}|\mathbf{y}_{\text{exp}}) = \frac{e^{-\Delta^2/2}}{\sqrt{(2\pi)^n \det(\Sigma(\mathbf{x}))}} \mathcal{P}(\mathbf{x})$$

$$\text{with } \Delta^2 = (\mathbf{y}(\mathbf{x}) - \mathbf{y}_{\text{exp}}) \cdot \Sigma(\mathbf{x})^{-1} \cdot (\mathbf{y}(\mathbf{x}) - \mathbf{y}_{\text{exp}})$$

We have a 20-dimensional parameter space and 514 datapoints

- Run model on 1000 `design` points, spaced on a latin hypercube
- `Interpolate` results by training a *Gaussian Process Emulator*

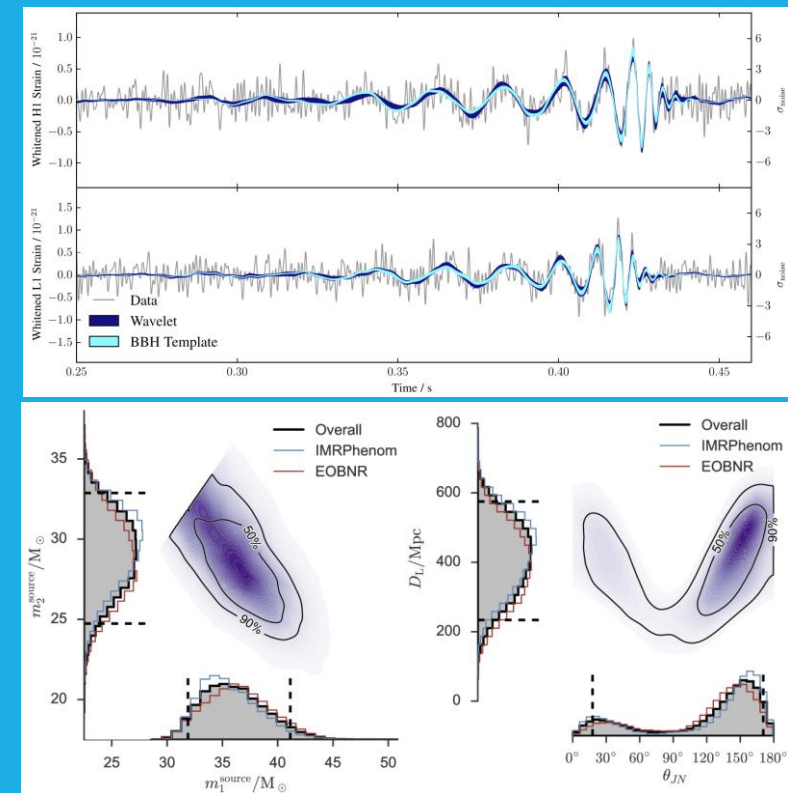
Markov Chain Monte Carlo (emcee2.2)

- Obtain sample of 10^6 likely values

Compare posterior with data

- From emulator (emulator has its own uncertainty estimate)
- A high statistics run at the optimal value (MAP, maximum a posteriori)

Example: gravitational waves





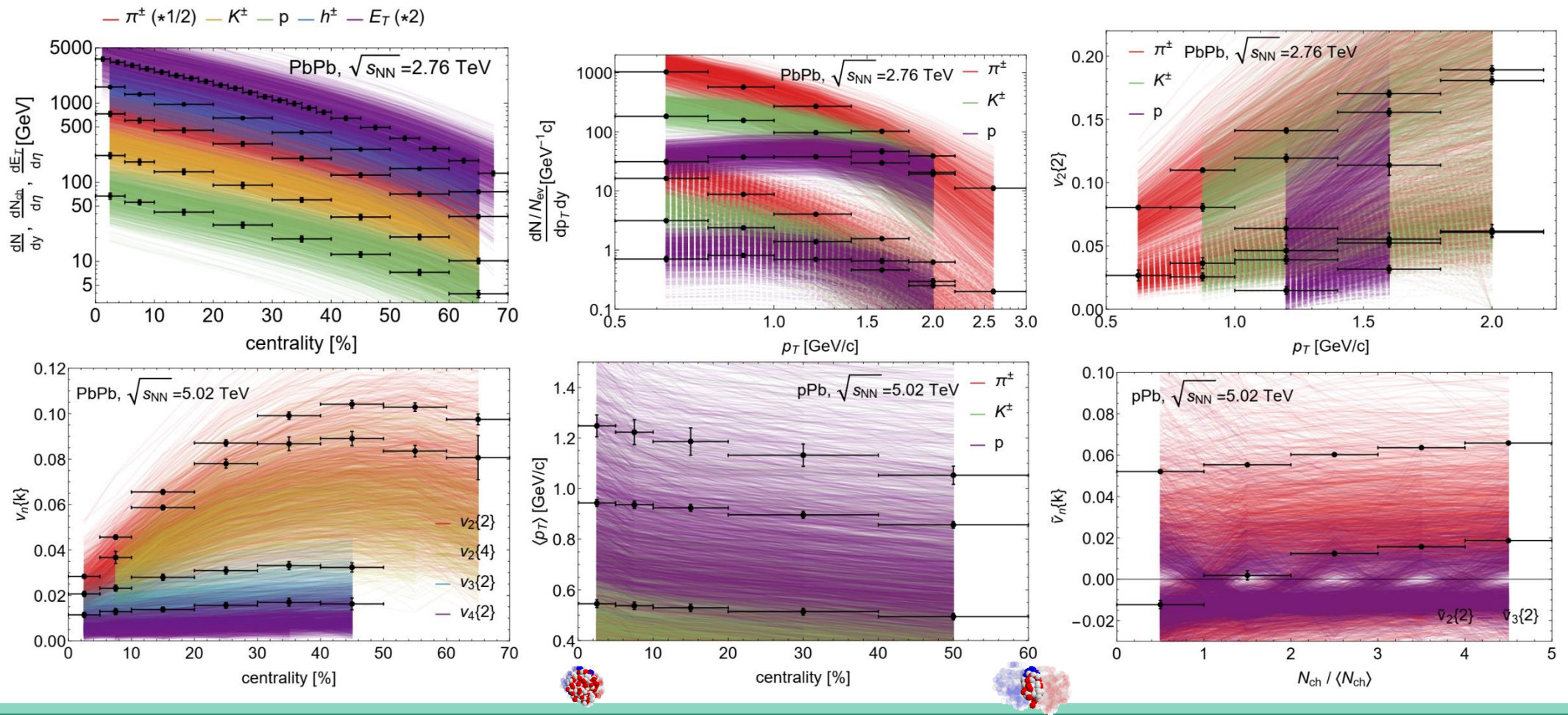
Computing

Making the model

- ▣ $1000 * 15000$ events = 15M events
- ▣ 50 events/hour/core (similar time for hydro + hadronic cascade)
- ▣ 300k CPU hours, both for PbPb, pPb etc

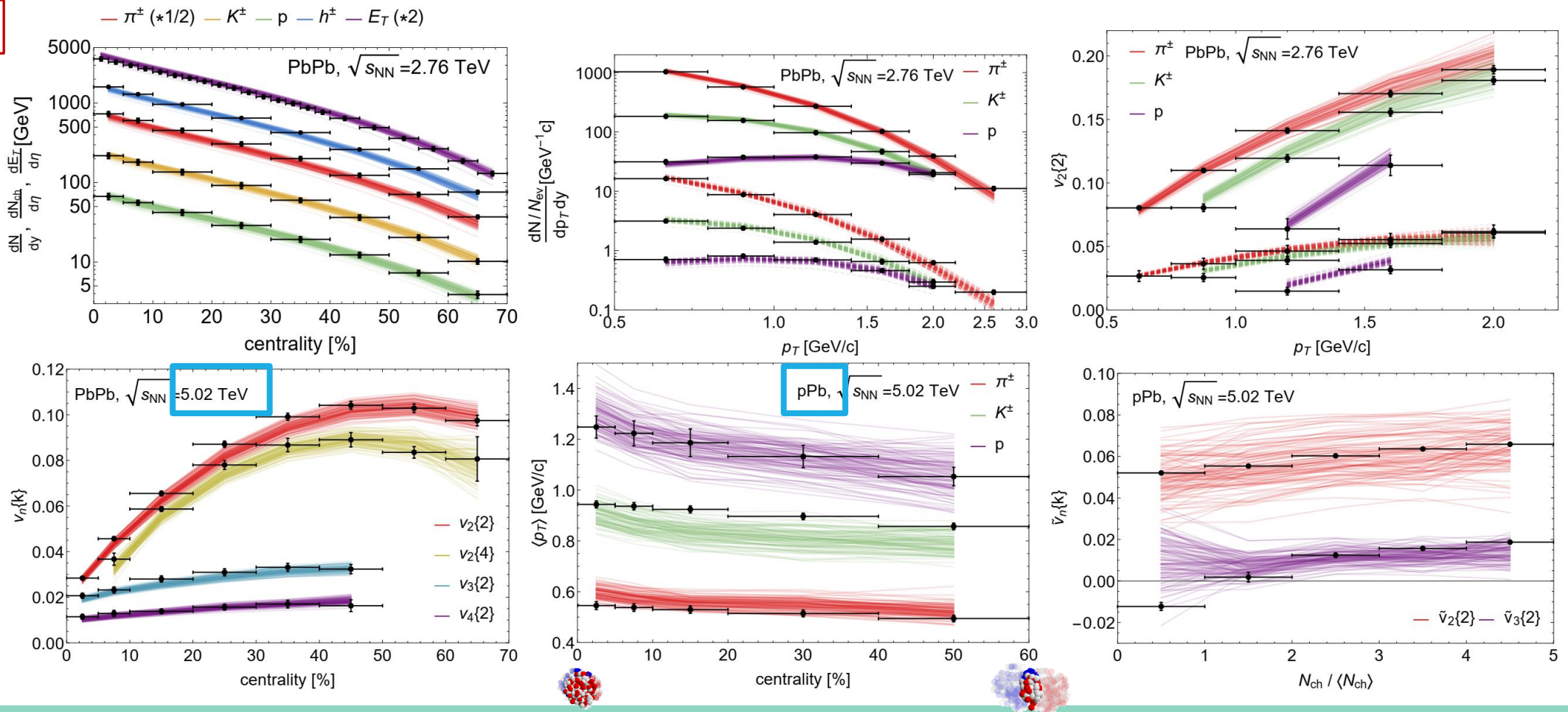
Computing the probability distributions for the parameters

- ▣ MCMC: evaluate emulator at 10M parameter setting: 1000 hours
 - ▣ Not properly parallelised ☹, takes a few weeks



Experimental observables:
a wealth of data

1. Yields, spectra, identified $v_n\{2\}$ versus p_T , pPb and PbPb (514 datapoints)
2. First study with a comprehensive analysis including p_T -differential observables

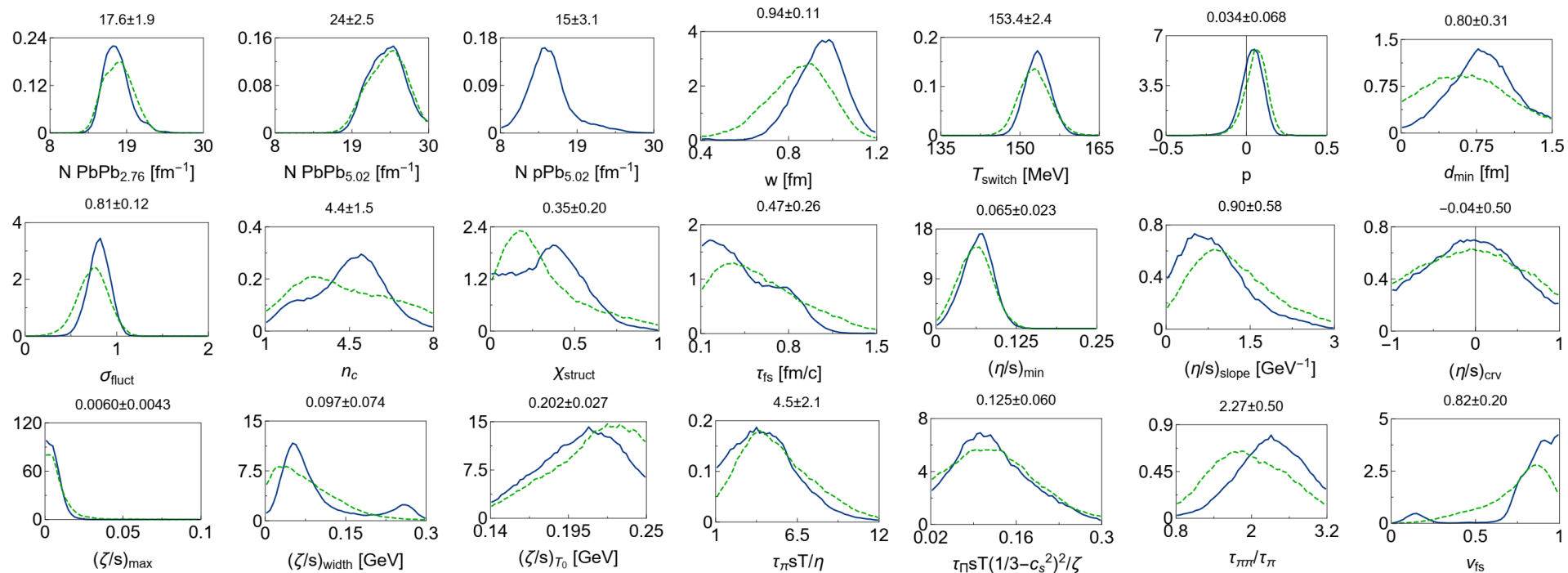


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Posterior distributions

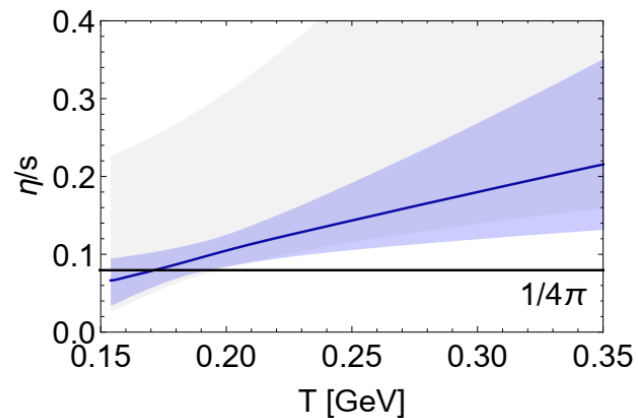
1. Dashed: without p Pb: indeed much flatter for e.g. n_c
2. Strong constraint on nucleon-nucleon fluctuations (also found at Duke)



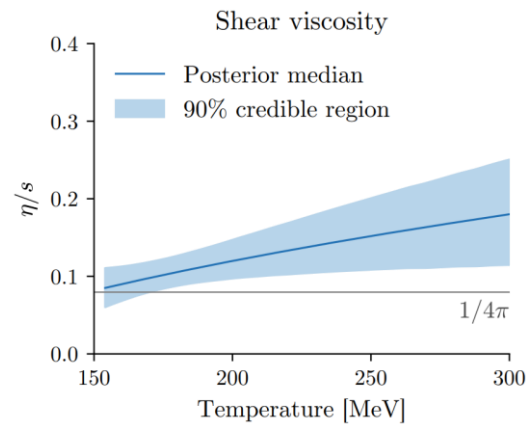
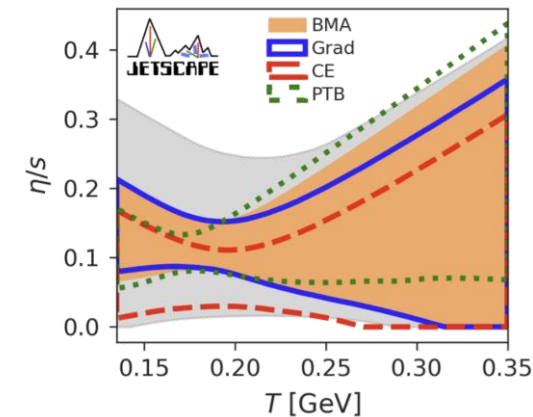
Posterior distributions – shear viscosity

1. Shear viscosity consistent with previous work

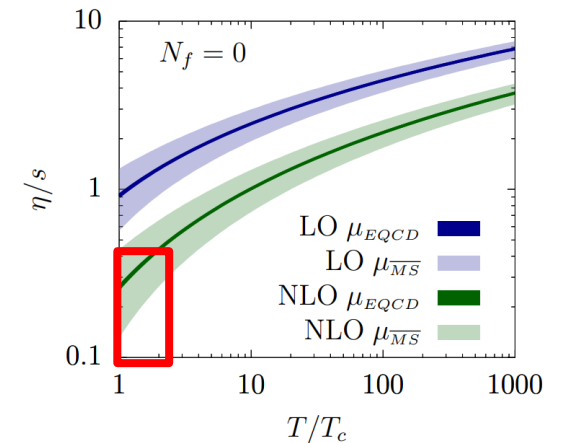
- More data, but also enlarged model \rightarrow similar constraint on η/s
- New JETSCAPE slightly broader band (larger priors, single PbPb energy but including RHIC)
- Consistent with state-of-the-art pQCD computations



Current work (2020)

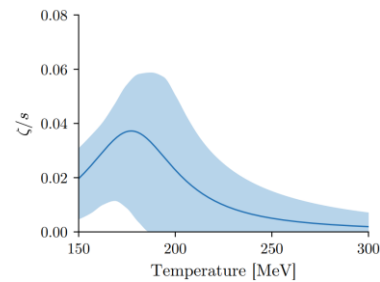
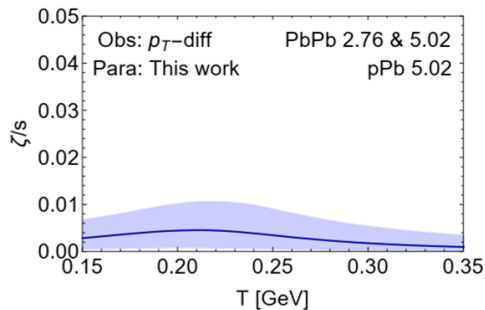
J. Bernhard, S. Moreland and S. Bass,
Nature Physics (2019)

JETSCAPE (2020)

Jacopo Ghiglieri, Guy Moore and Derek Teaney
QCD Shear Viscosity at (almost) NLO (2018)

Posterior distributions – bulk viscosity:

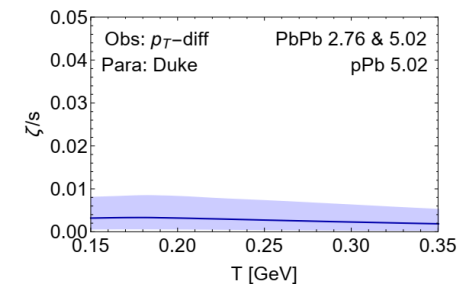
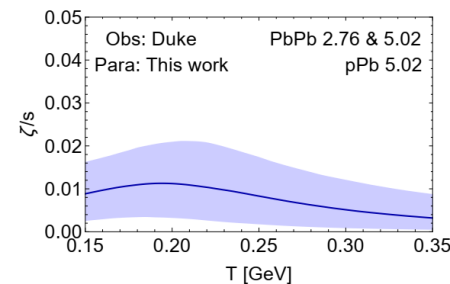
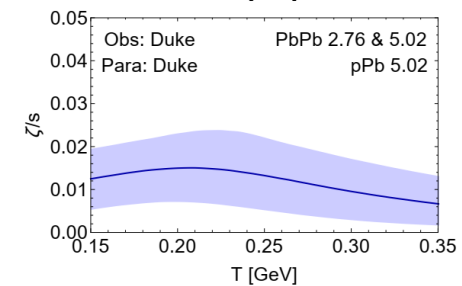
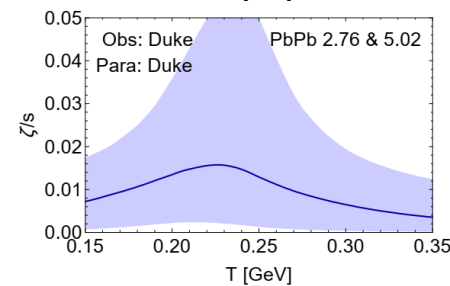
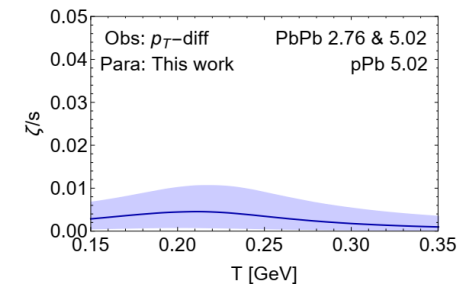
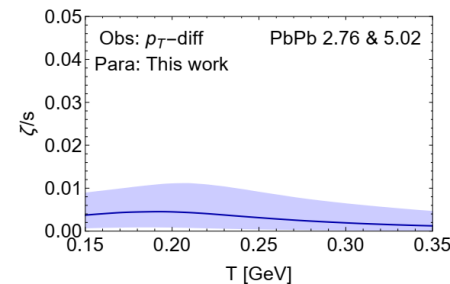
Much smaller, even consistent with zero



J. Bernhard, S. Moreland and S. Bass,
Nature Physics (2019)

Bulk viscosity, varied several aspects:

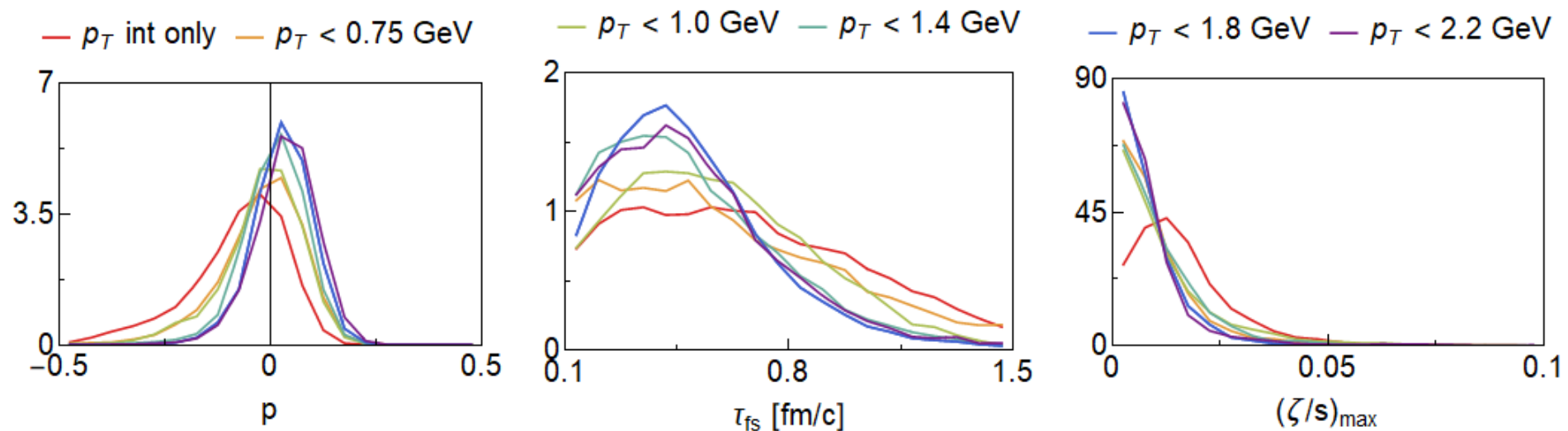
- More limited parameter set
 - All versus only 'Duke'
- Include or not include p-Pb collisions
- Include p_T -differential observables



Sensitivity to p_T -differential observables

Vary maximum p_T of observables to identify their constraining power

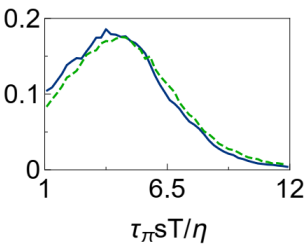
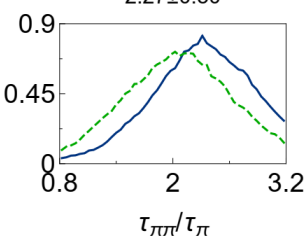
- First two bins till 1 GeV give strongest constraints (if at all, selection shown)
 - Also due to tougher statistics at higher p_T : emulator error
- Viscous corrections at freeze-out very uncertain at large p_T
 - Encouraging that results are quite insensitive to high p_T bins



Posterior distributions – 2nd order transport

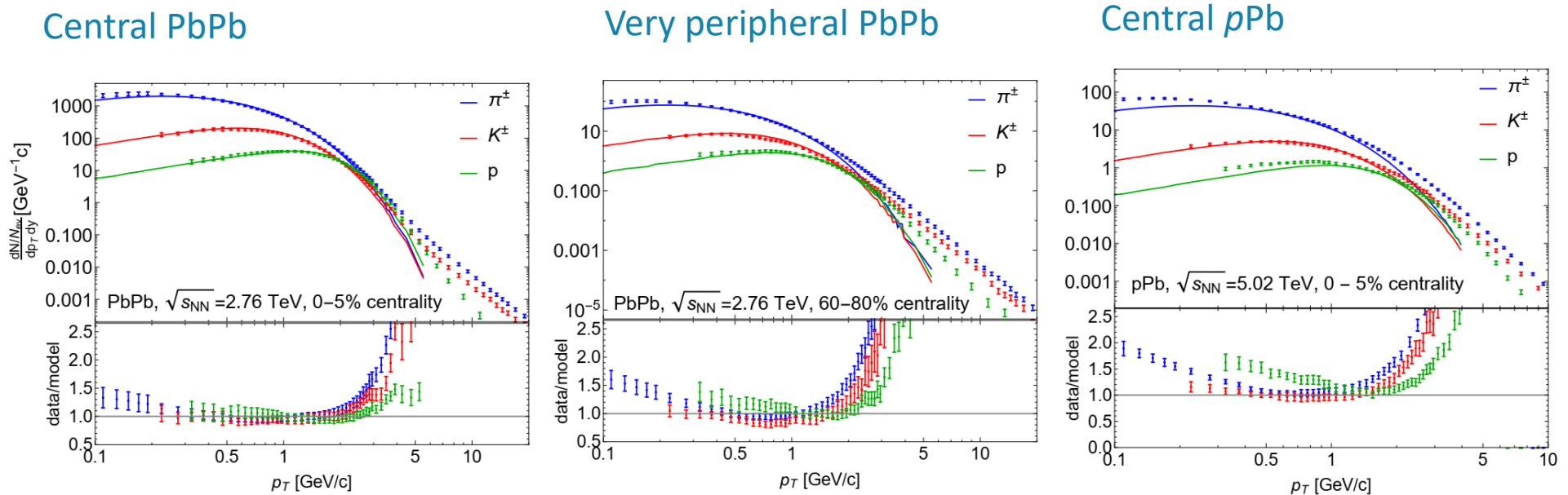
τ_π and $\tau_{\pi\pi}$ can be compared to strong and weak coupling values

- Both consistent, AdS/CFT slightly favoured for $\tau_{\pi\pi}$

	this work	AdS/CFT	kinetic theory
$\frac{\tau_\pi sT}{\eta}$	<p>4.5±2.1</p> 	$4 - \log(4) \approx 2.61$	5
$\frac{\tau_{\pi\pi}}{\tau_\pi}$	<p>2.27±0.50</p> 	$\frac{88}{35(2-\log 2)} \approx 1.92$	$\frac{10}{7} \approx 1.43$

MAP: maximum a posteriori: spectra

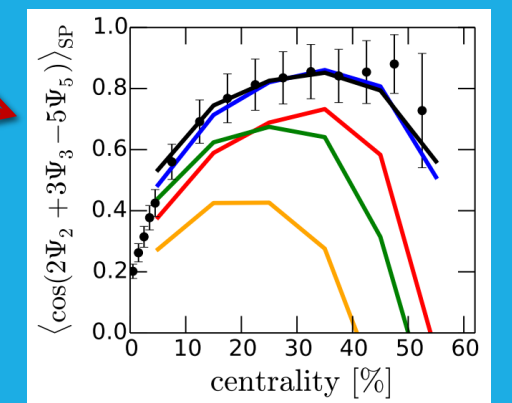
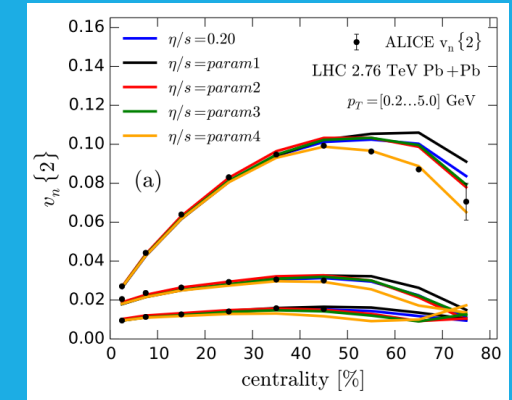
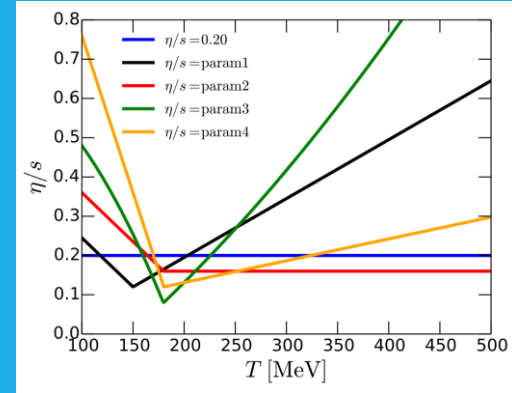
1. High statistics run at (almost) optimal parameters, compared with ALICE data
 - 400k events (PbPb) and 4M events (pPb)



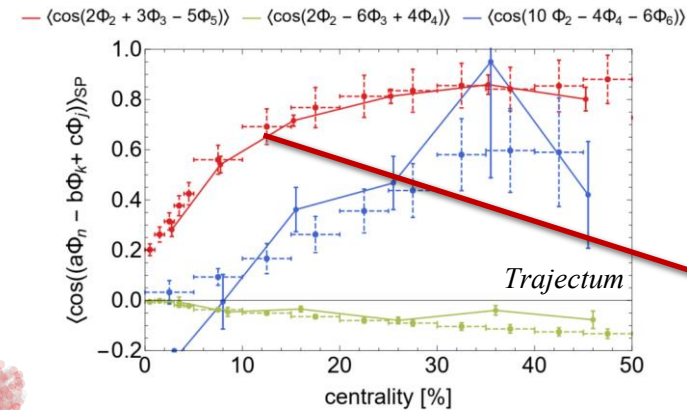
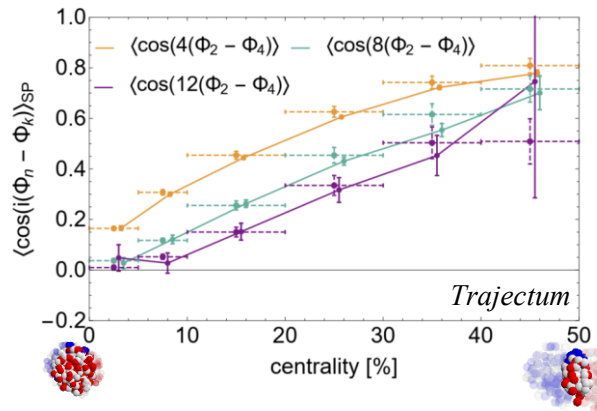
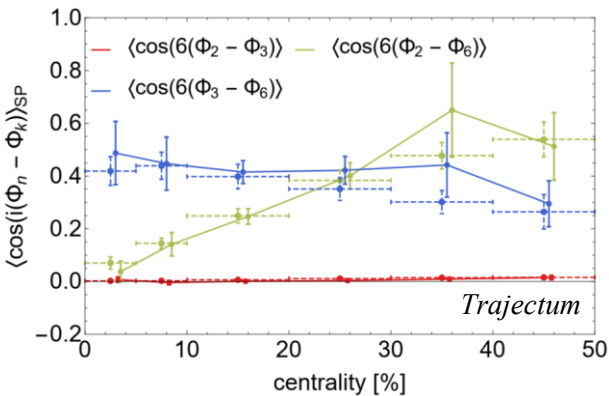
$$\frac{d\bar{N}}{d\varphi} = \frac{\bar{N}}{2\pi} \left(1 + 2 \sum_{n=1}^{\infty} \bar{v}_n \cos(n(\varphi - \bar{\Psi}_n)) \right)$$

MAP: PbPb event plane angles

1. A non-trivial check: the event-plane angle correlations (not used in Bayesian)
2. Non-trivial agreement: needs a very specific $\eta/s(T)$ (even when fitting v_n)

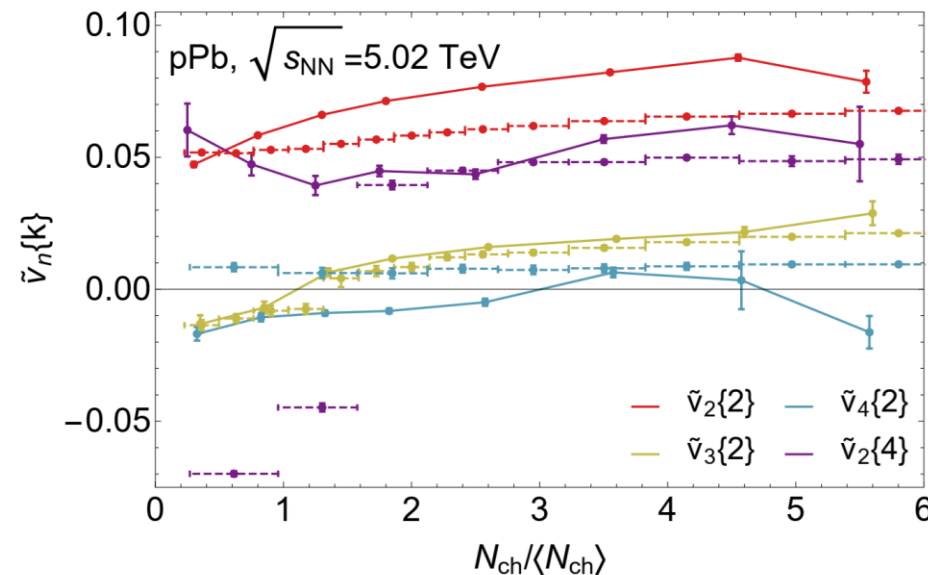


Data by **ATLAS**



MAP: p Pb anisotropic flow

1. Emulator and MCMC are less precise for p Pb: uncertainty is statistical only
2. Shows potential to obtain imaginary $v_n\{2\}$ (= negative $\tilde{v}_n\{k\}$), in agreement with ATLAS low multiplicity result
3. Sheds new light on discussion of hydro versus sign of $\langle\langle 2 \rangle\rangle_n$



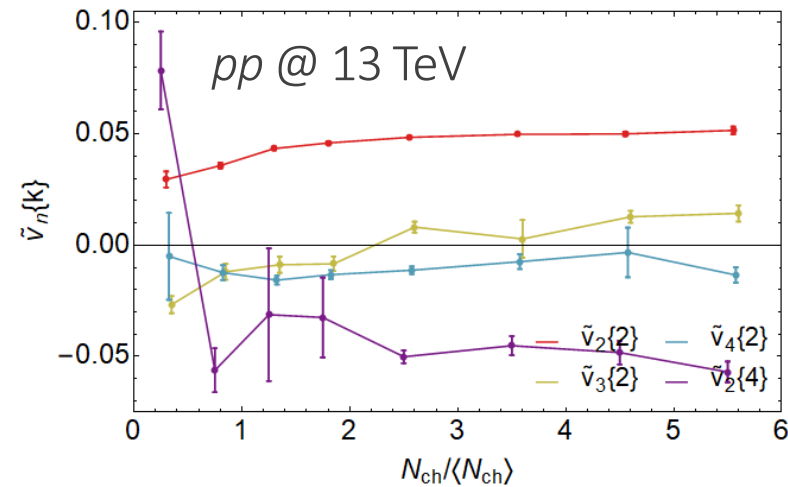
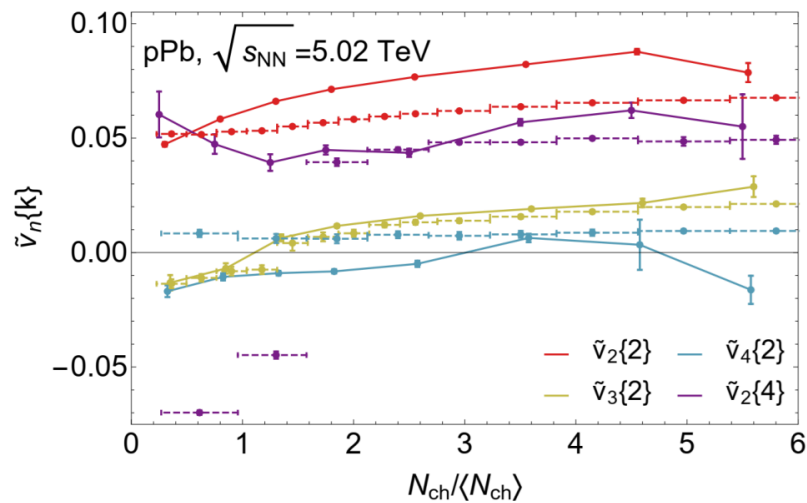
$$Q_n = \sum_{i=1}^M e^{in\phi_i}$$

$$\langle 2 \rangle_n = \frac{|Q_n|^2 - M}{M(M-1)}$$

$$v_n\{2\} = \sqrt{\langle\langle 2 \rangle\rangle_n}$$

MAP: pp anisotropic flow

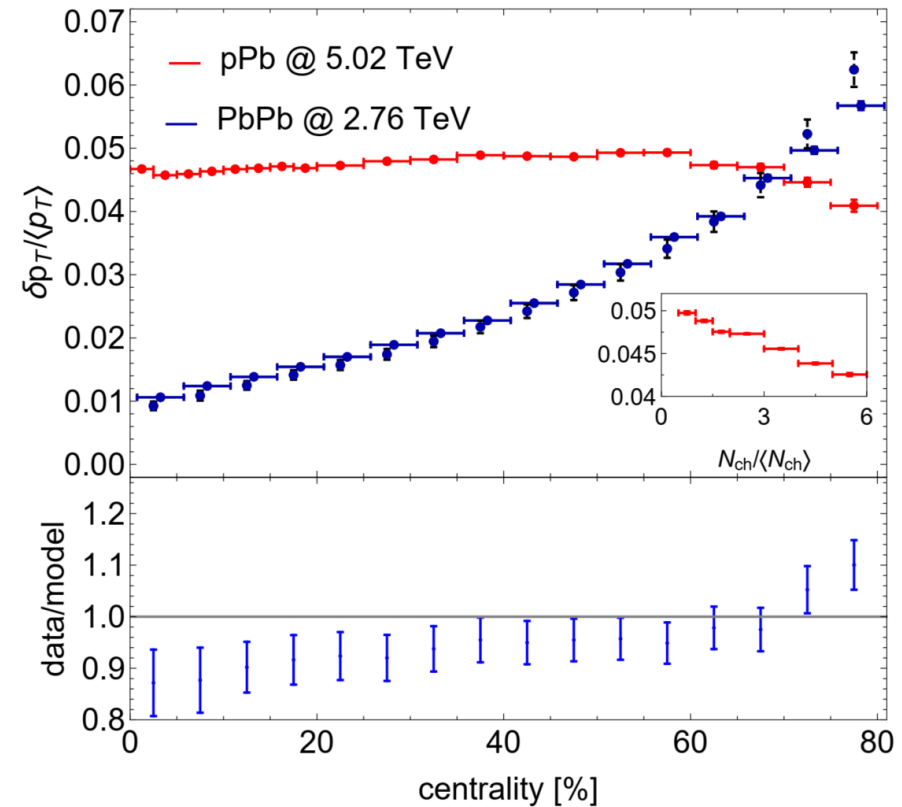
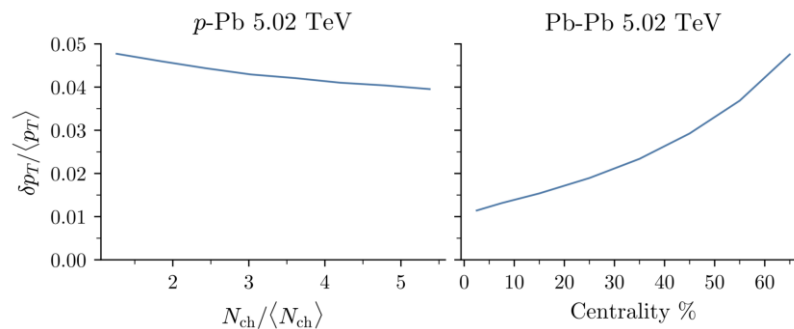
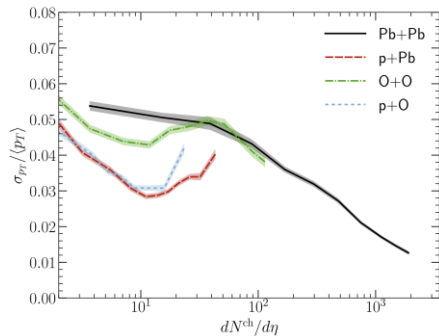
1. Preliminary results for pp ; different sign for $v_2\{4\}^2$ (?)



$$\delta p_T^2 = \left\langle \left\langle (p_{T,i} - \langle p_T \rangle)(p_{T,j} - \langle p_T \rangle) \right\rangle \right\rangle$$

A prediction for p Pb: momentum fluctuations

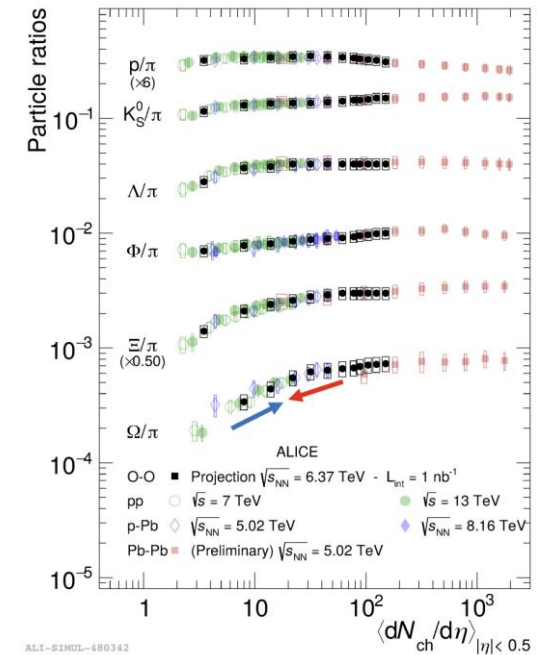
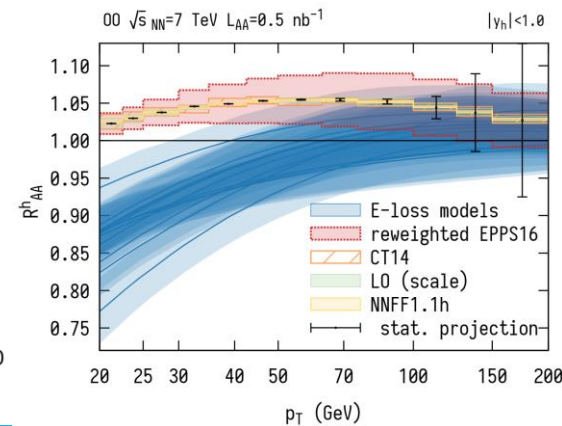
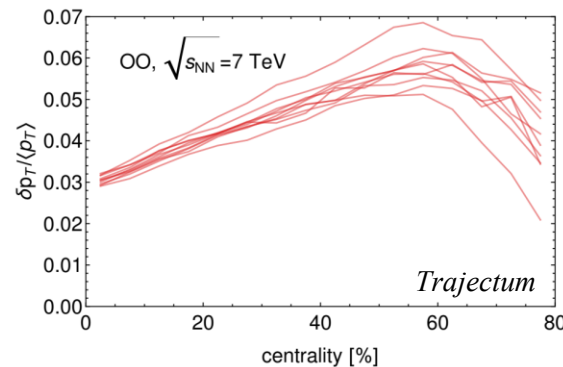
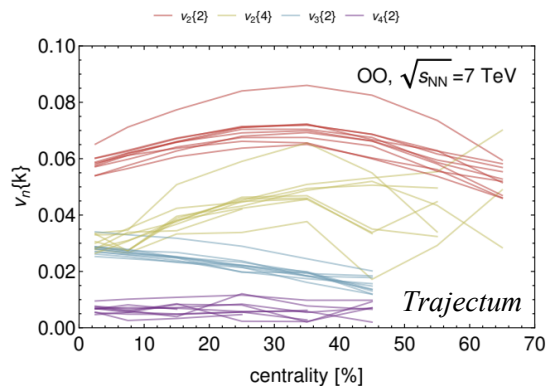
1. Rather hard to find observables that have not been measured
2. Typically hard to match to data
3. Roughly comparable to results by Moreland et al, higher than Shen et al:



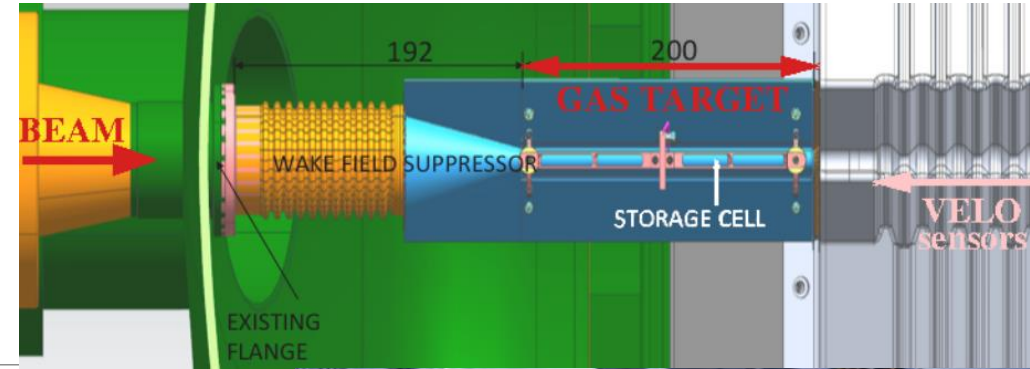


Oxygen & small systems: LHC as a light ion collider

1. Will help resolving the 'flow in small systems' puzzle
2. A well-controlled environment to discover *jet quenching* in small systems
3. Bridges the gap between pPb and PbPb for strangeness enhancement
4. pO: of crucial interest in modelling cosmic ray showers (LHCf)



Light ions and SMOG2 LHCb at fixed target



1. Interesting idea: 'contaminate' beam with gas (only at LHCb)
2. Fixed target (gas is at rest); options: H, He, N, O, Ar, Ne, Kr, Xe, ...
3. Lower energy ($\sqrt{s_{NN}} \simeq 110 \text{ GeV}$): complementary to colliding set-up
4. Possible with p , Pb and O in the beam
5. Data taking simultaneous: sizeable integrated lumi: 100 pb^{-1}

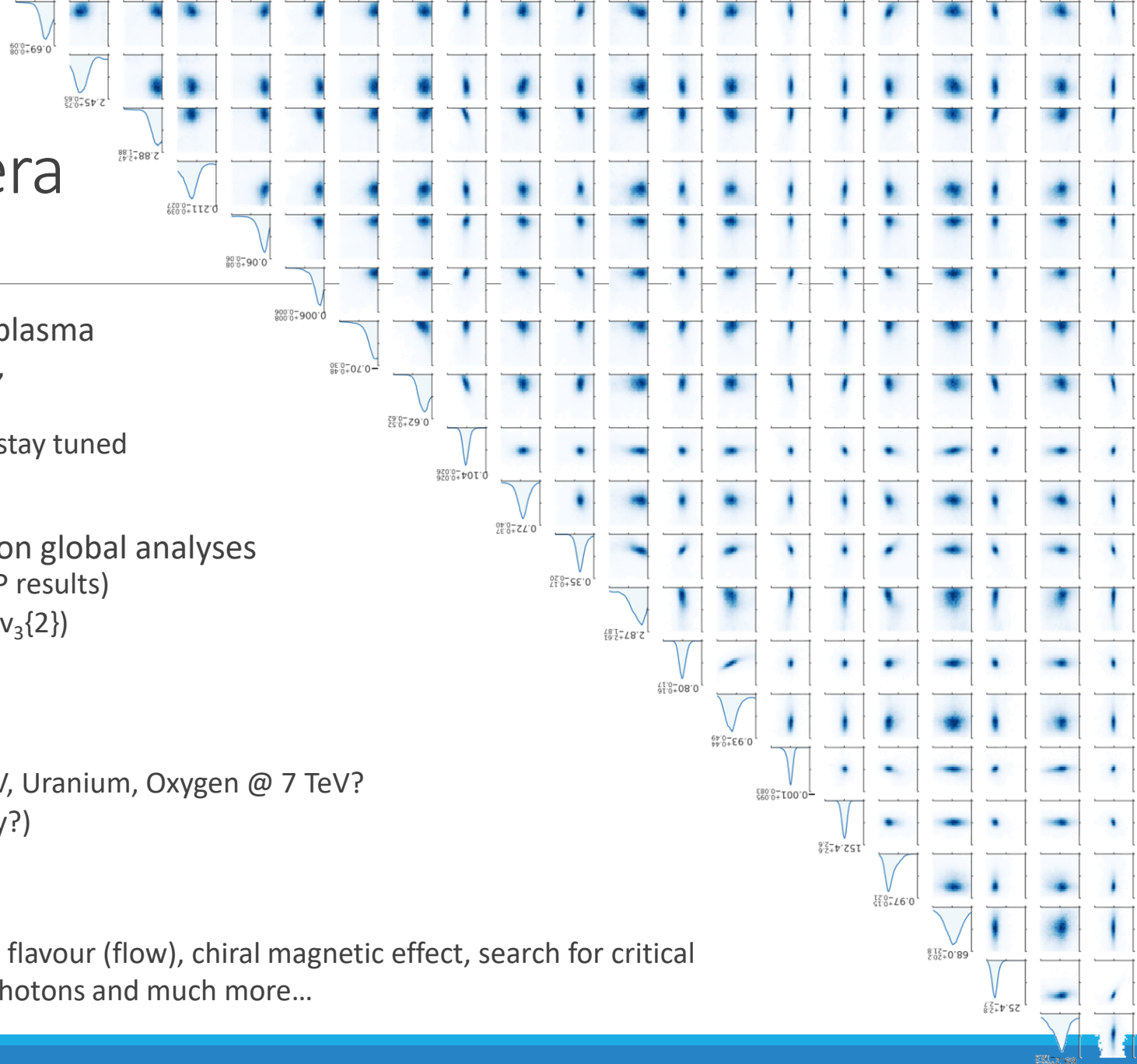


Global analysis perspective: need for a wide variety of colliding systems and energies



Theory of HIC in the LHC era

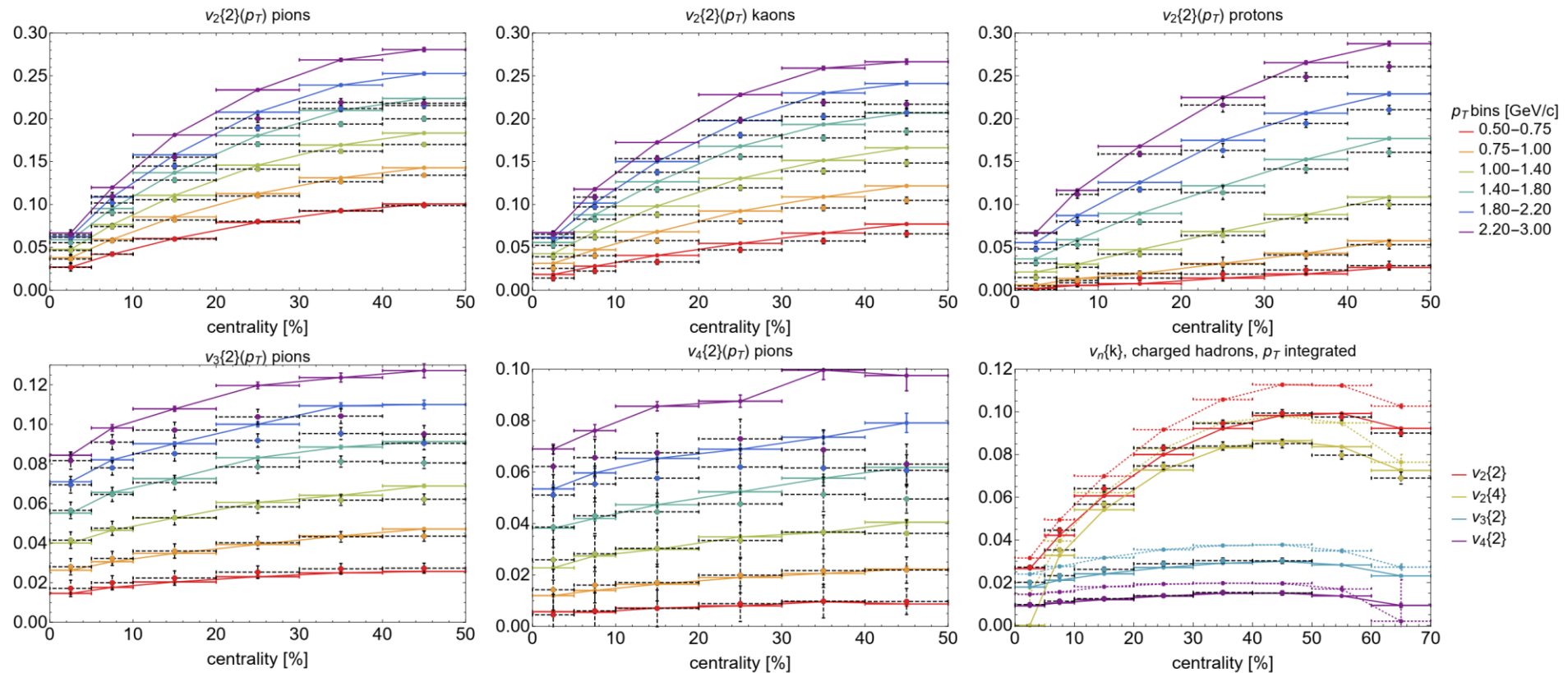
1. A road to precision analysis of the quark-gluon plasma
 - Measuring transport and initial stage 'beyond η/s ,
 - **Revisited bulk viscosity: surprisingly small**
 - Hint towards second order transport coefficients: stay tuned
2. p_T -differential anisotropic flow sheds new light on global analyses
 - An excellent fit, including event-plane angles (MAP results)
 - Interesting MAP results on flow in pPb (imaginary $v_3\{2\}$)
 - Initial stage and degrees of freedom: stay tuned
3. Only our first study
 - Include more systems and energies: Au @ 200 GeV, Uranium, Oxygen @ 7 TeV?
 - Finer observables (but sometimes statistics hungry?)
4. Did I skip anything?
 - progress in jet (substructure) modifications, heavy flavour (flow), chiral magnetic effect, search for critical point, rapidity dependence, bound quark states, photons and much more...



Back up

MAP: PbPb anisotropic flow

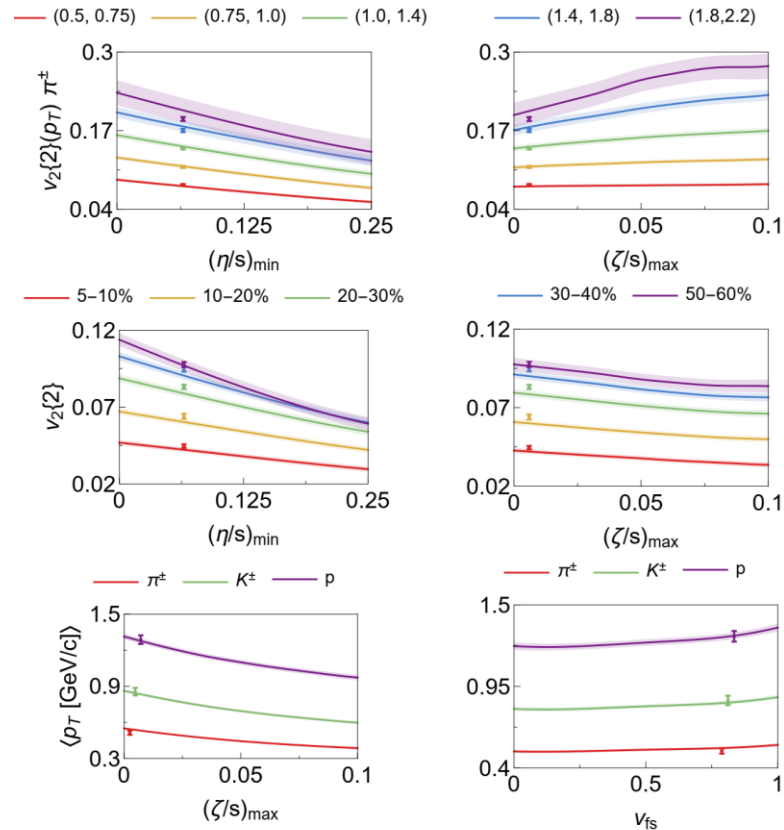
1. Anisotropic flow matches well, except for a few high p_T bins



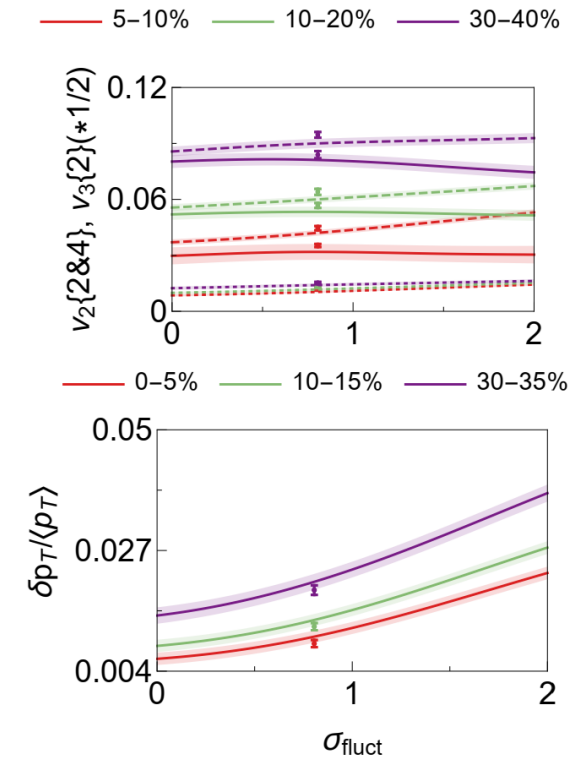
The emulator: Viscosities and fluctuations

also note: emulator uncertainty (50-60%, or $v_2\{4\}$)

v_2 increases in every p_T bin, but decreases on average for ζ .

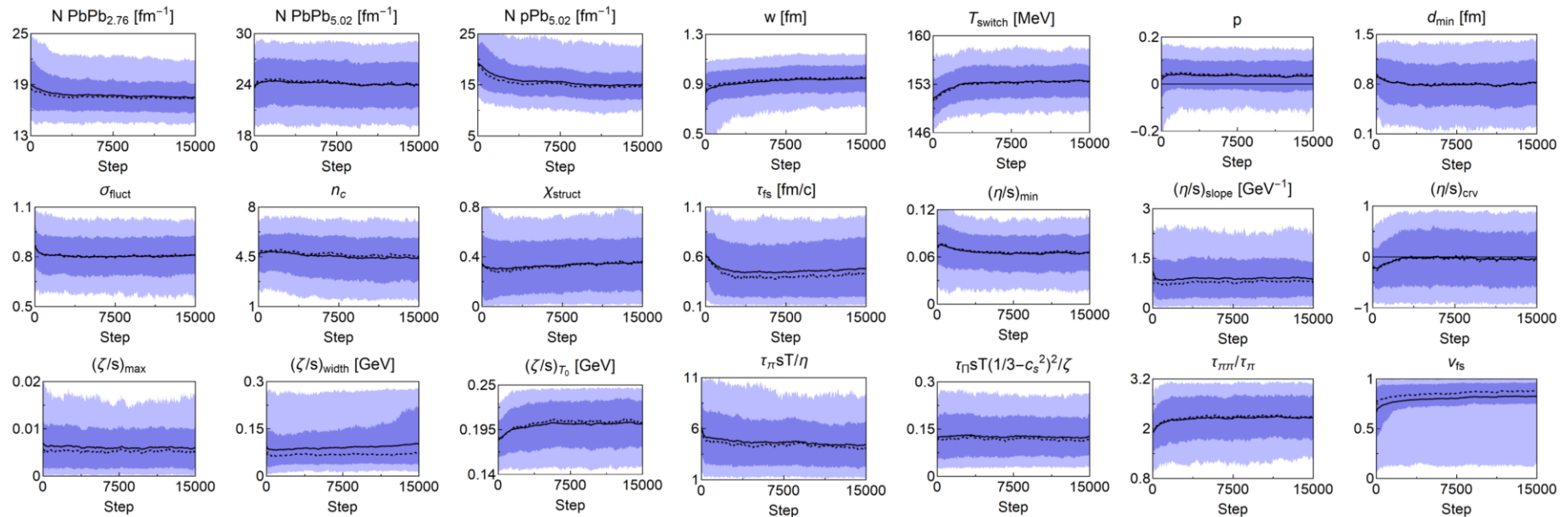


$v_2\{2\} - v_2\{4\}$ increases when increasing fluctuations



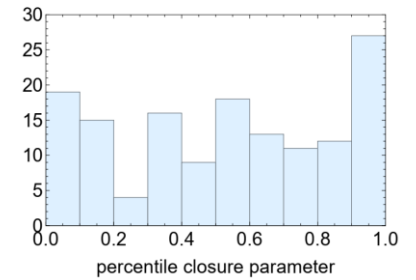
Making the chain: (pt)emcee

1. Constructing a chain of posteriors (MCMC); we use emcee
2. Necessary to verify convergence:

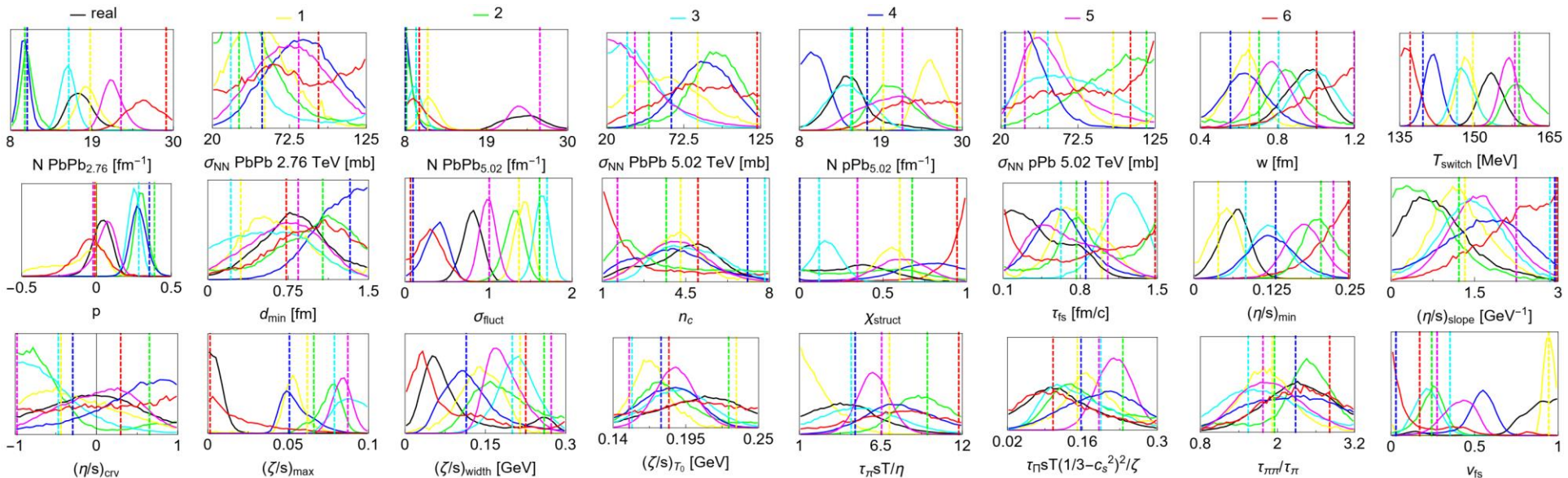


3. Still quite expensive to emulate:
600 walkers * 15000 steps * 10% acceptance with 2000 points = 3 weeks

Selected results: closure test

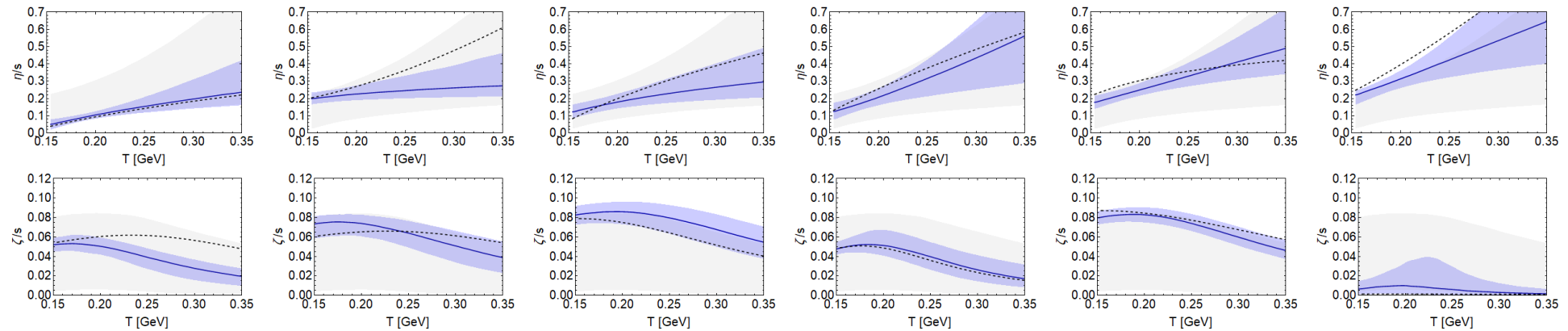


1. We chose six random parameter points (sometimes at edge of prior)
 - Try to extract parameters from model-generated 'experimental' data
2. Verifies model + shows sensitivity data on parameter
 - Output indeed consistent with input
 - Sensitive to viscosities, less so for second order

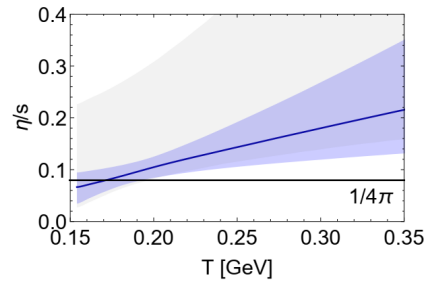


Closure test: viscosities

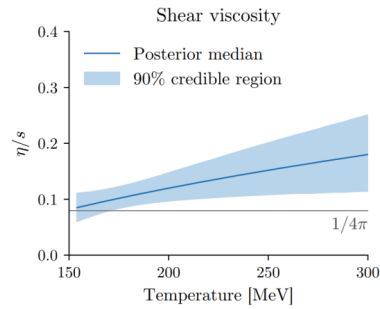
1. Closure test works well for both viscosities
2. Most sensitive to low-T region



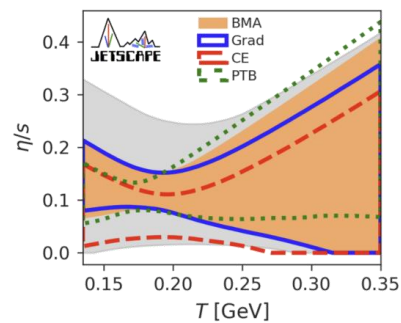
Posterior distributions – shear viscosity



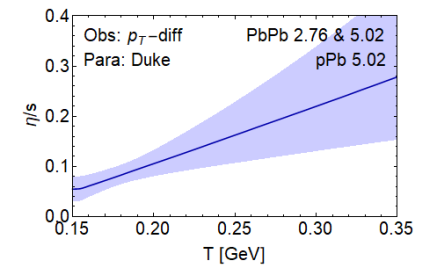
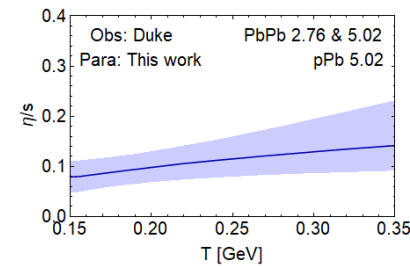
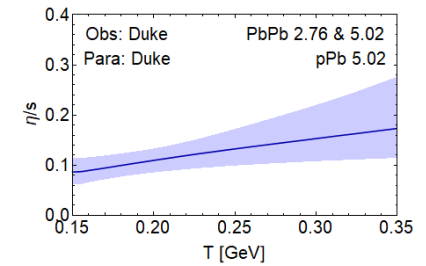
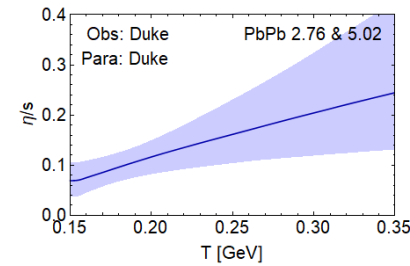
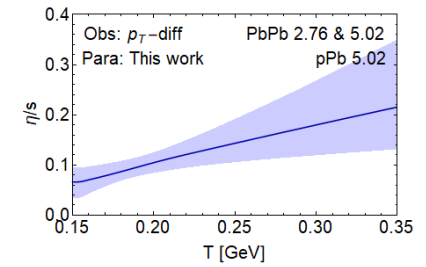
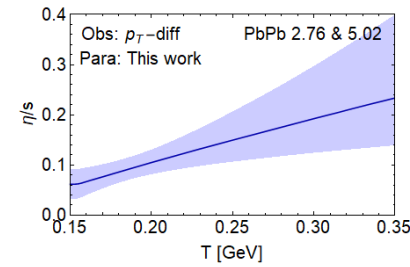
Current work (2020)



J. Bernhard, S. Moreland and S. Bass, Nature Physics (2019)



JETSCAPE (2020)



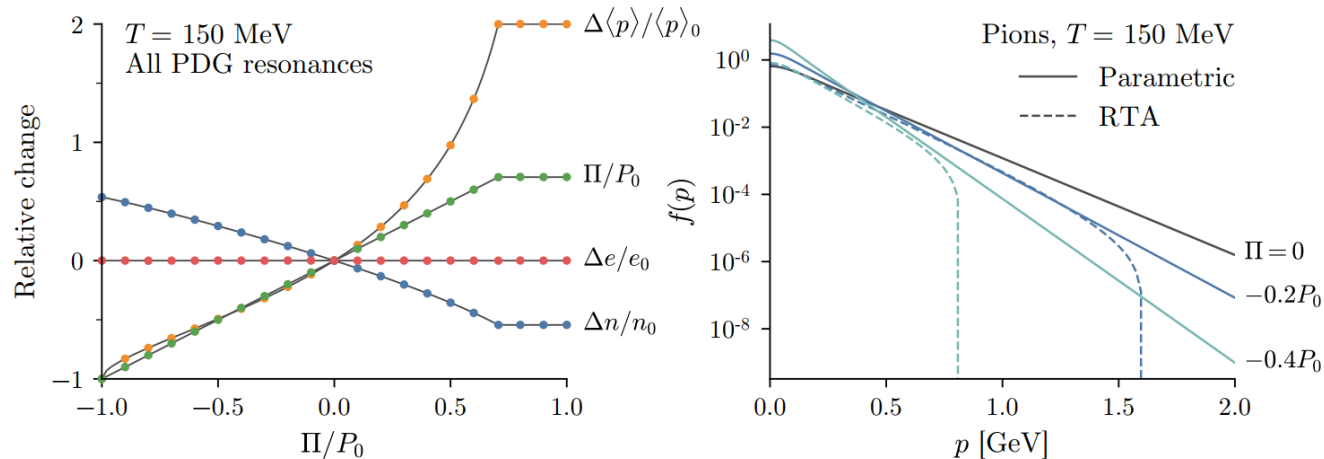
Particlization: viscous corrections

$$T^{\mu\nu} = \sum_{\text{sp}} g \int \frac{d^3p}{(2\pi)^3} \frac{p^\mu p^\nu}{E} f(p),$$

1. Particles in fluid restframe cannot be in thermal equilibrium
2. Several methods that (only/mostly) agree for small deviations

$f(p) \rightarrow z_{\text{bulk}} f(p + \lambda_{\text{bulk}} p)$ parametric, rescale p : fix z and λ such that e and P match

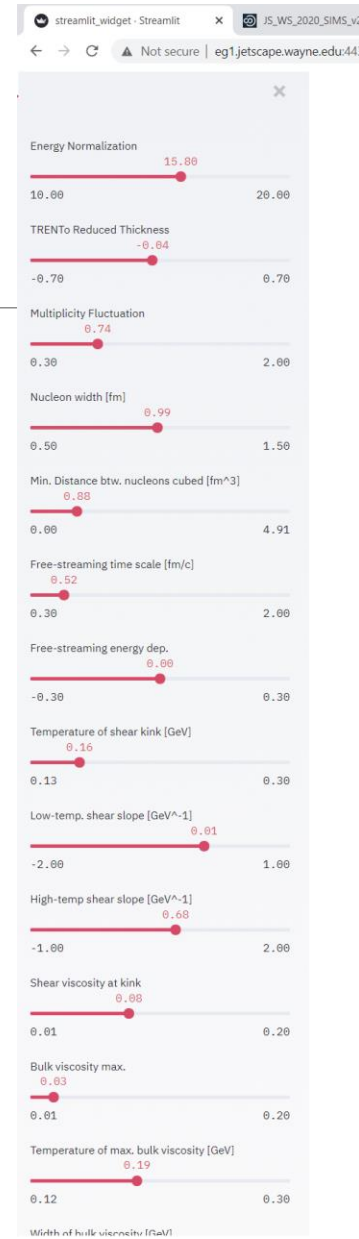
$\delta f = f_0(1 \pm f_0) \frac{\tau}{ET} \left[\frac{1}{2\eta} p^i p^j \pi_{ij} + \frac{1}{\zeta} \left(\frac{p^2}{3} - c_s^2 E^2 \right) \Pi \right]$ change $f(p)$ directly, motivated by RTA



'Parametric' clearly better at high p_T , but somewhat ad-hoc and species independent

Comparison with JETSCAPE

Results seem to be in relatively good agreement. Data is quite consistent without a sizeable bulk viscosity.



The experimentally measured observables by the [ALICE collaboration](#) are shown as black dots.

The last row displays the temperature dependence of the specific shear and bulk viscosities (red lines), as determined by different parameters on the left sidebar.

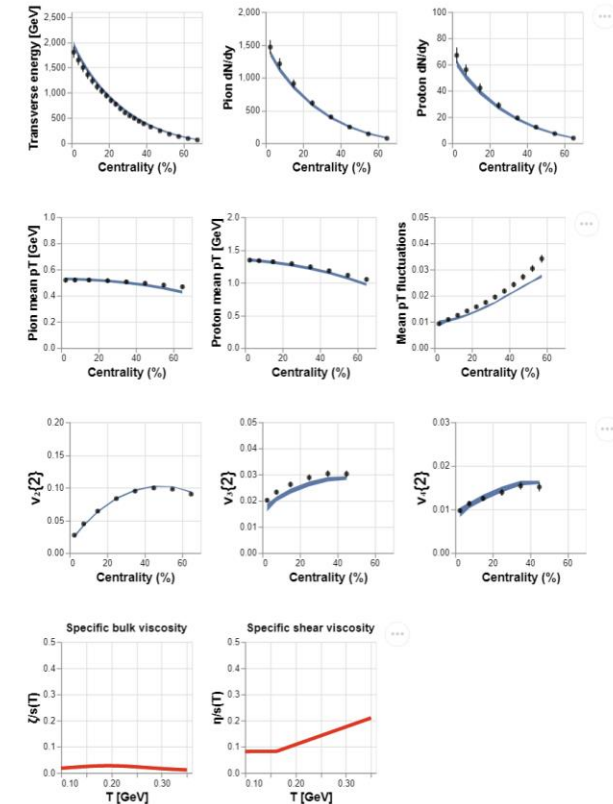
By default, these parameters are assigned the values that fit the experimental data *best* (maximize the likelihood).

An important modelling ingredient is the particization model used to convert hydrodynamic fields into individual hadrons. Three different viscous correction models can be selected by clicking the "Particization model" button below.

Particization model

Pratt-Torrieri-Bernhard

[Reset](#)

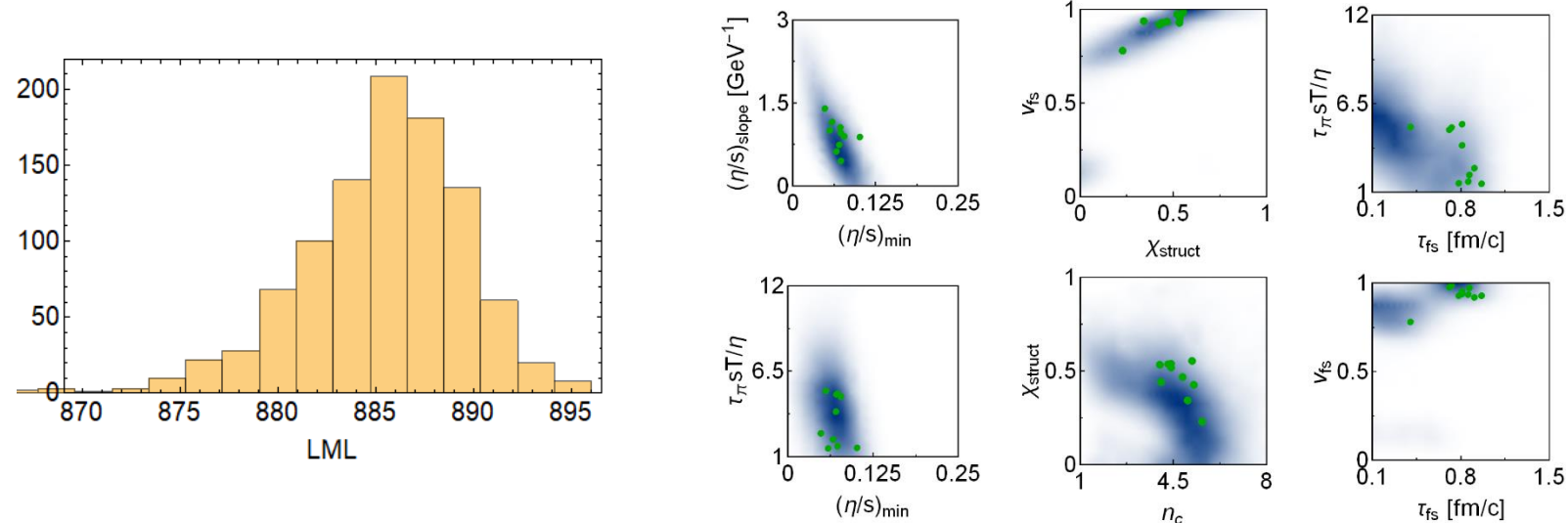


MAP: maximum a posteriori

$$\mathcal{P}(\mathbf{x}|\mathbf{y}_{\text{exp}}) = \frac{e^{-\Delta^2/2}}{\sqrt{(2\pi)^n \det(\Sigma(\mathbf{x}))}} \mathcal{P}(\mathbf{x})$$

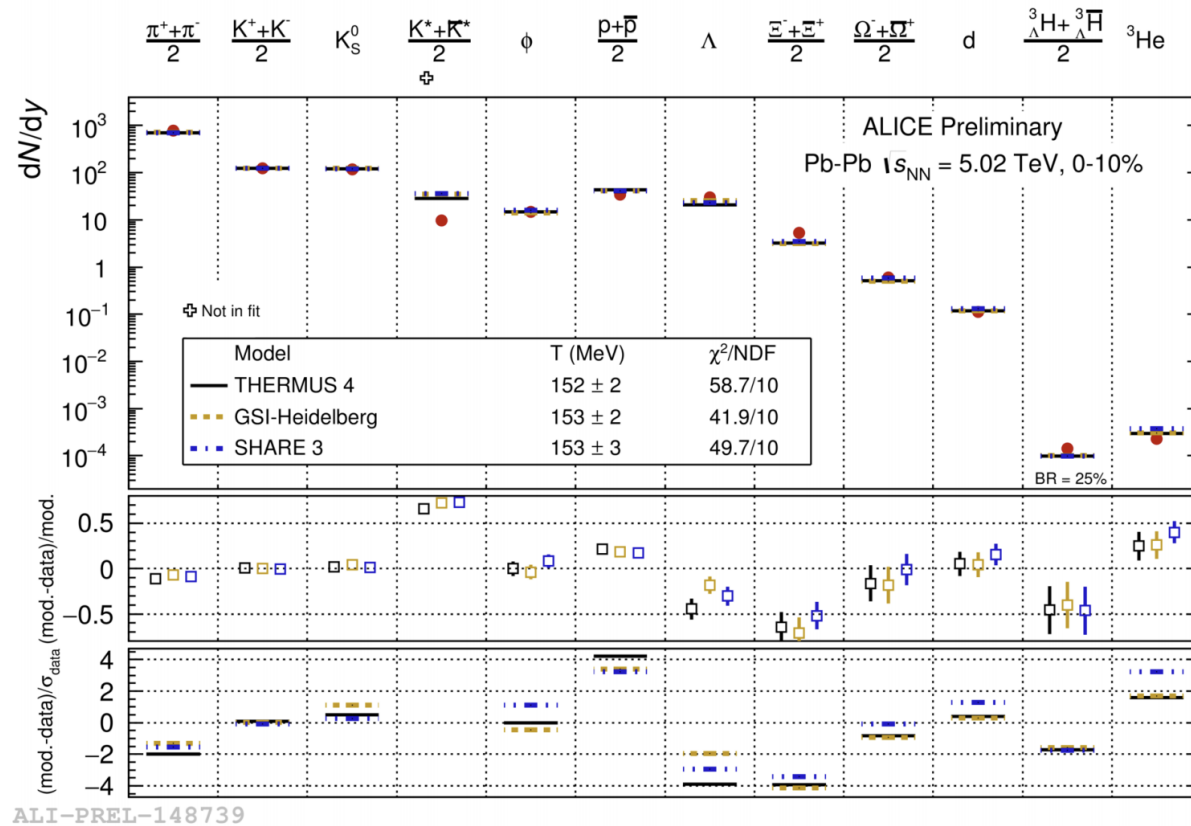
with $\Delta^2 = (\mathbf{y}(\mathbf{x}) - \mathbf{y}_{\text{exp}}) \cdot \Sigma(\mathbf{x})^{-1} \cdot (\mathbf{y}(\mathbf{x}) - \mathbf{y}_{\text{exp}})$,

1. Subtle to find the 'true' maximum of the LML:
2. take 3000 points from the posterior chain, plot highest 10 LMLs (LML > 893):



3. We decided to use expectation value of each param as 'MAP' (LML = 884)

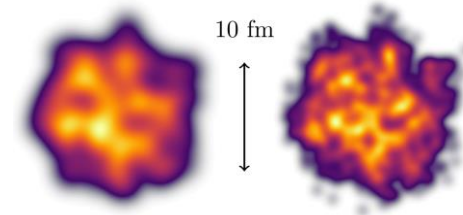
Strangeness and thermal fits: close but not perfect



Initial geometry: two (three?) uncertainties

1. The structure of nucleons

- n_c constituents of Gaussian subwidth v within a nucleon of width w
- Nucleons placed according to MC Glauber



2. How do colliding (sub)nucleons deposit their energy? $\mathcal{T} = \left(\frac{\mathcal{T}_A^p + \mathcal{T}_B^p}{2} \right)^{1/p}$

- For $p = 0$ we get $\mathcal{T} = \sqrt{\mathcal{T}_A \mathcal{T}_B}$: close to EKRT or Holography ($\mathcal{T} \approx (\mathcal{T}_A \mathcal{T}_B)^{4/9}$)
- Does not quite allow binary scaling?

3. (Quantum) fluctuations in the above: Gamma-distribution:

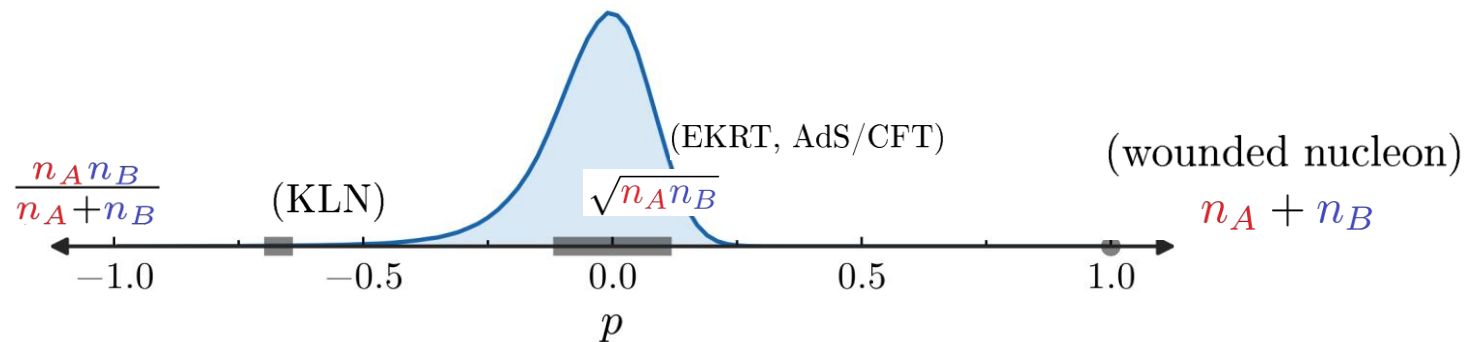
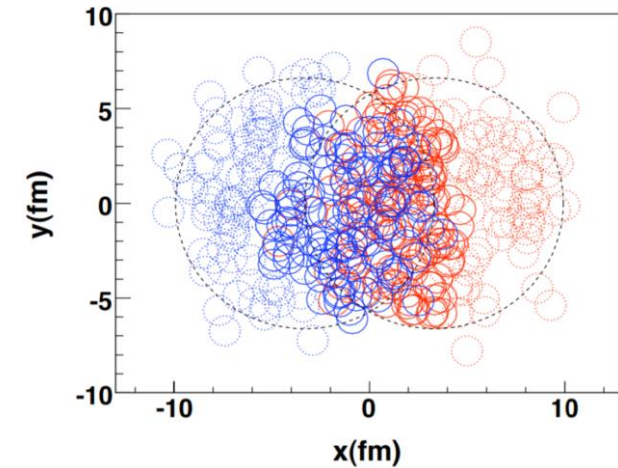
- Goes beyond MC Glauber fluctuations

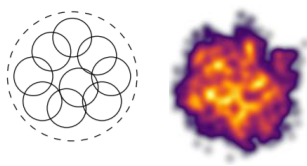
$$p(T) = \frac{1}{\Gamma(1/\sigma)\sigma^{1/\sigma}} T^{(1-\sigma)/\sigma} e^{-T/\sigma}$$

Constraints on initial conditions

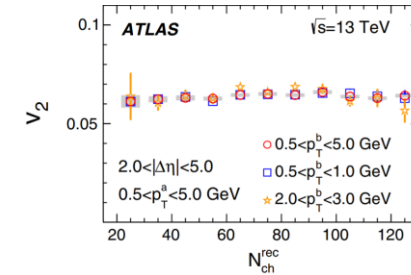
Trento parametrization allows for many models:

- Distinguishes KLN, EKRT or AdS/CFT, wounded nucleons
- Data clearly rules out KLN and wounded nucleons

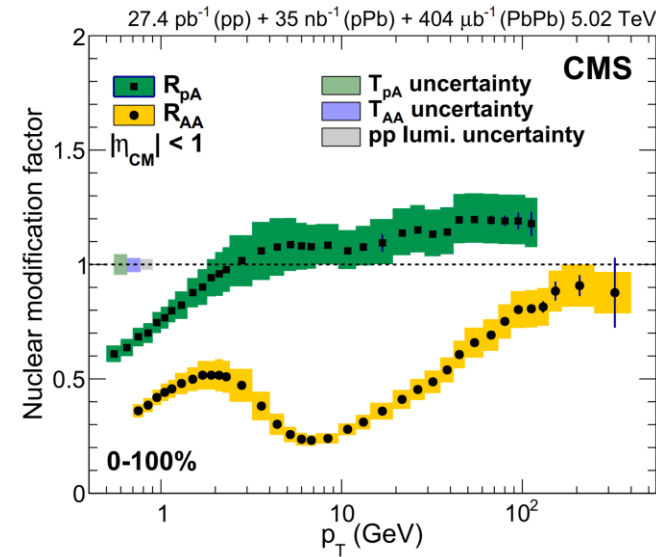
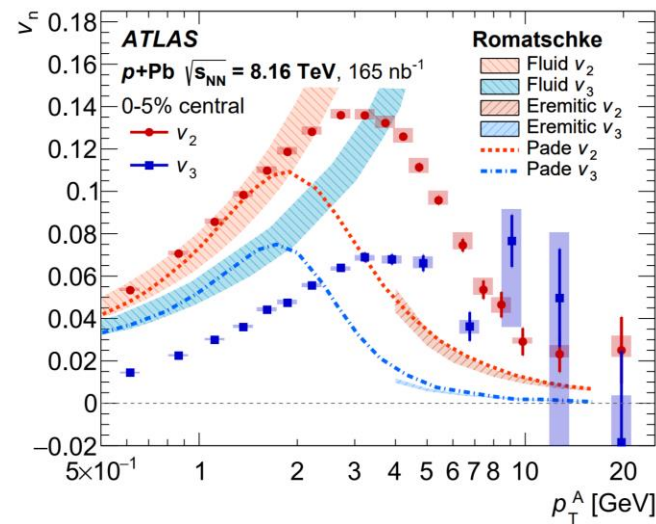




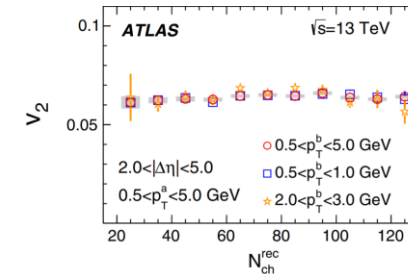
A puzzle: flow in pPb or pp collisions?



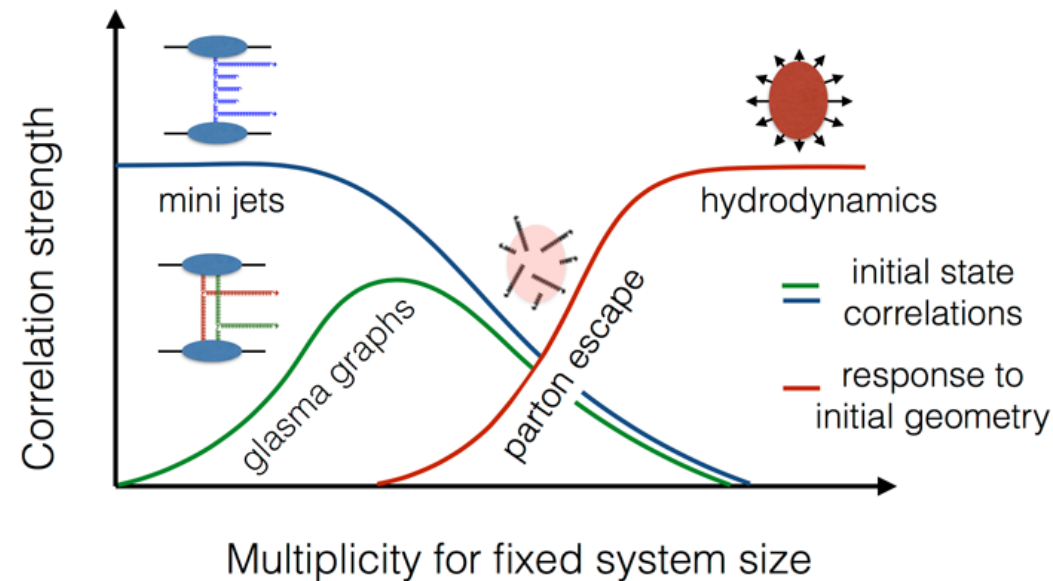
1. There seems to be flow
 - Quite some modeling, but everything consistent with hydro (does not proof hydro!)
2. But: nuclear modification > 1: **no jet energy loss**, but nuclear effects dominate



A puzzle: flow in pPb or pp collisions?



1. It should be possible to `turn off' hydro (small system)
2. Tantalising option: Combination of mini-jets/glasma connecting to hydro
 - Challenge: hard to explain constant v_2 ; spectra not necessarily thermal
 - Quenching versus flow is challenging for any model



The future

1. Comprehensive (Bayesian) analysis, with more complete dataset
2. For small systems: runs with light or intermediate ions:
 - **Pb–Pb at $\sqrt{s_{NN}} = 5.5$ TeV**, $L_{\text{int}} = 13 \text{ nb}^{-1}$ (ALICE, ATLAS, CMS), 2 nb^{-1} (LHCb)
 - **pp at $\sqrt{s} = 5.5$ TeV**, $L_{\text{int}} = 600 \text{ pb}^{-1}$ (ATLAS, CMS), 6 pb^{-1} (ALICE), 50 pb^{-1} (LHCb)
 - **pp at $\sqrt{s} = 14$ TeV**, $L_{\text{int}} = 200 \text{ pb}^{-1}$ with low pileup (ALICE, ATLAS, CMS)
 - **p–Pb at $\sqrt{s_{NN}} = 8.8$ TeV**, $L_{\text{int}} = 1.2 \text{ pb}^{-1}$ (ATLAS, CMS), 0.6 pb^{-1} (ALICE, LHCb)
 - **pp at $\sqrt{s} = 8.8$ TeV**, $L_{\text{int}} = 200 \text{ pb}^{-1}$ (ATLAS, CMS, LHCb), 3 pb^{-1} (ALICE)
 - **O–O at $\sqrt{s_{NN}} = 7$ TeV**, $L_{\text{int}} = 500 \mu\text{b}^{-1}$ (ALICE, ATLAS, CMS, LHCb)
 - **p–O at $\sqrt{s_{NN}} = 9.9$ TeV**, $L_{\text{int}} = 200 \mu\text{b}^{-1}$ (ALICE, ATLAS, CMS, LHCb)
 - **Intermediate AA**, e.g. $L_{\text{int}}^{\text{Ar–Ar}} = 3\text{--}9 \text{ pb}^{-1}$ (about 3 months) gives NN luminosity equivalent to Pb–Pb with $L_{\text{int}} = 75\text{--}250 \text{ nb}^{-1}$

Similar in pPb or PbPb

1. pPb geometry intrinsically more spherical \rightarrow lower v_2

