

TIME-DEPENDENT PROPAGATION TIMES AND ENERGY LOSSES OF PROTONS IN THE HELIOSPHERE:

A SOLAR MODULATION MODELLING IN LIGHT OF NEW COSMIC-RAY DATA FROM OBSERVATIONS

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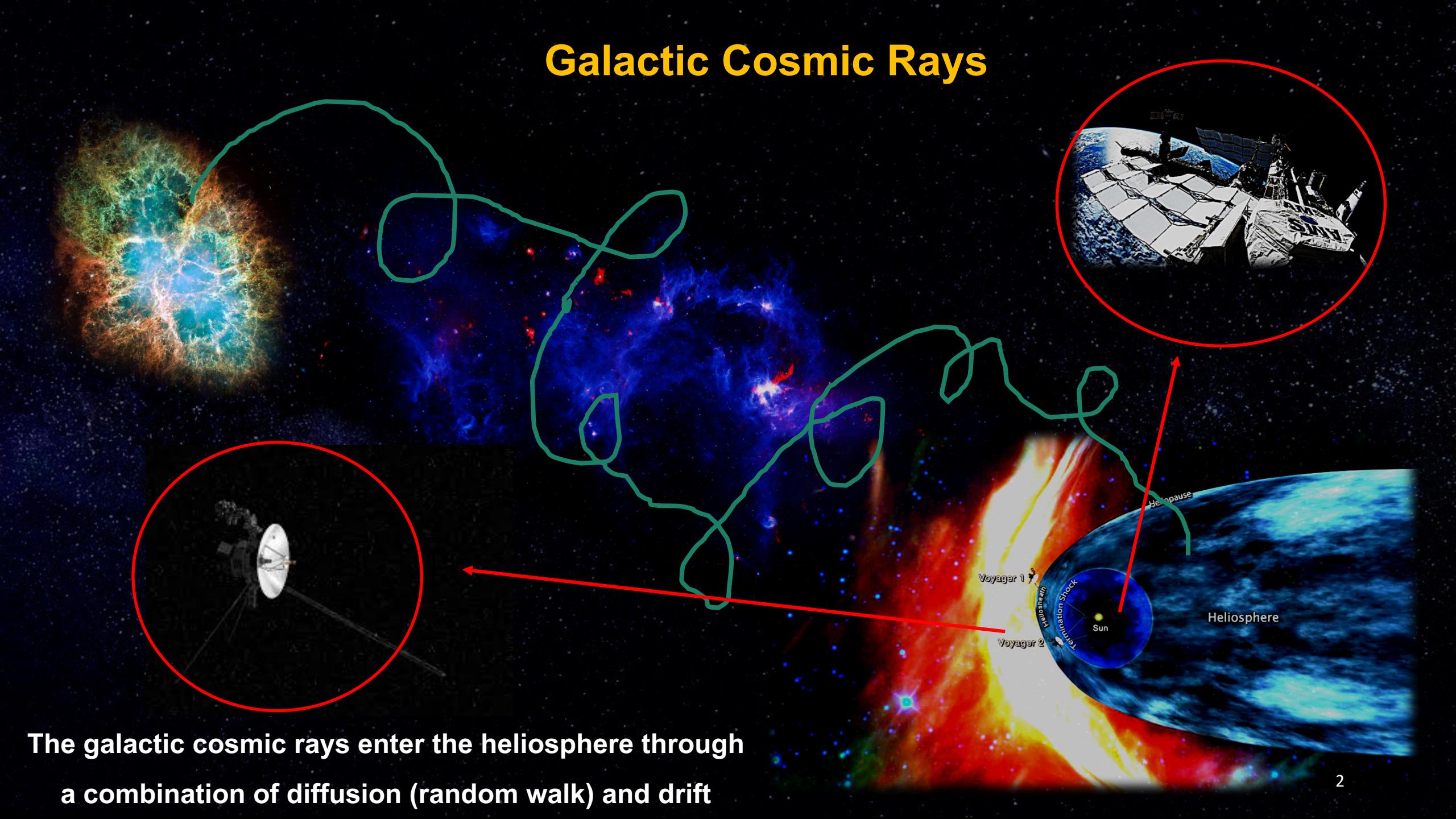
ASI - SSDC and INFN – Roma Tor Vergata, Italy

University of Perugia & INFN; ASI-UniPG 2019-2-HH.0

25 July 2022

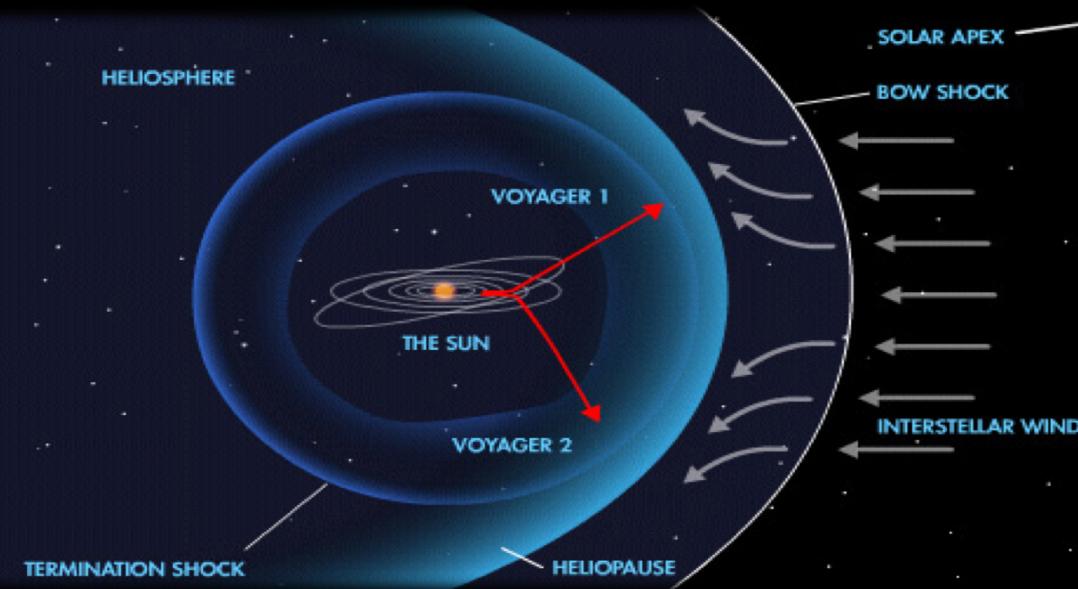


Galactic Cosmic Rays



The galactic cosmic rays enter the heliosphere through
a combination of diffusion (random walk) and drift

Solar Modulation



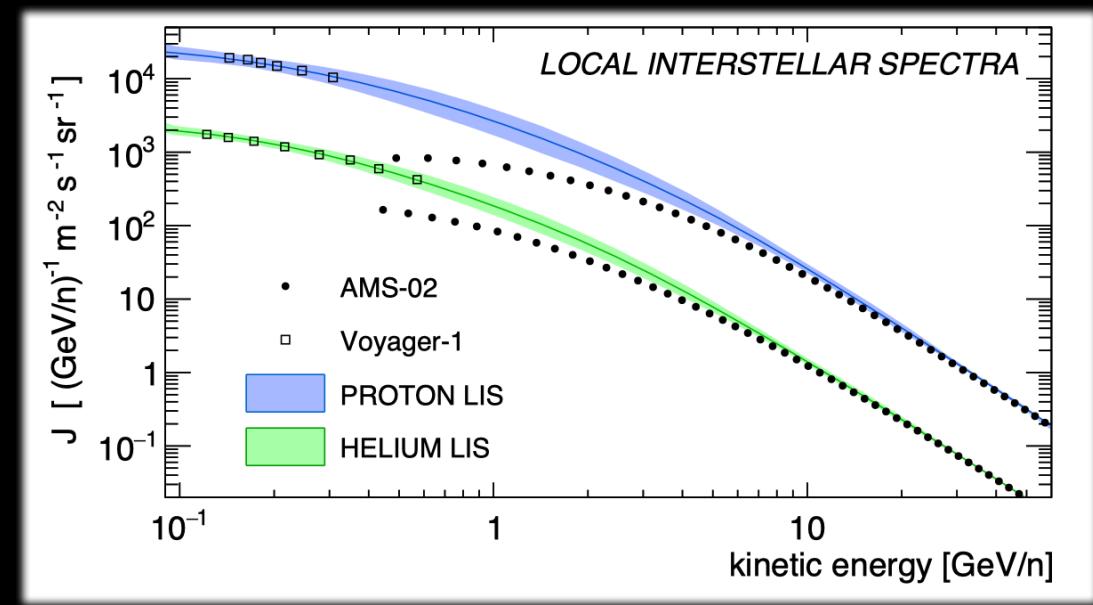
- This leads to significant global and temporal variations in their intensity and in their energy as a function of position inside the heliosphere.
- This process is identified as the **solar modulation** of cosmic rays .

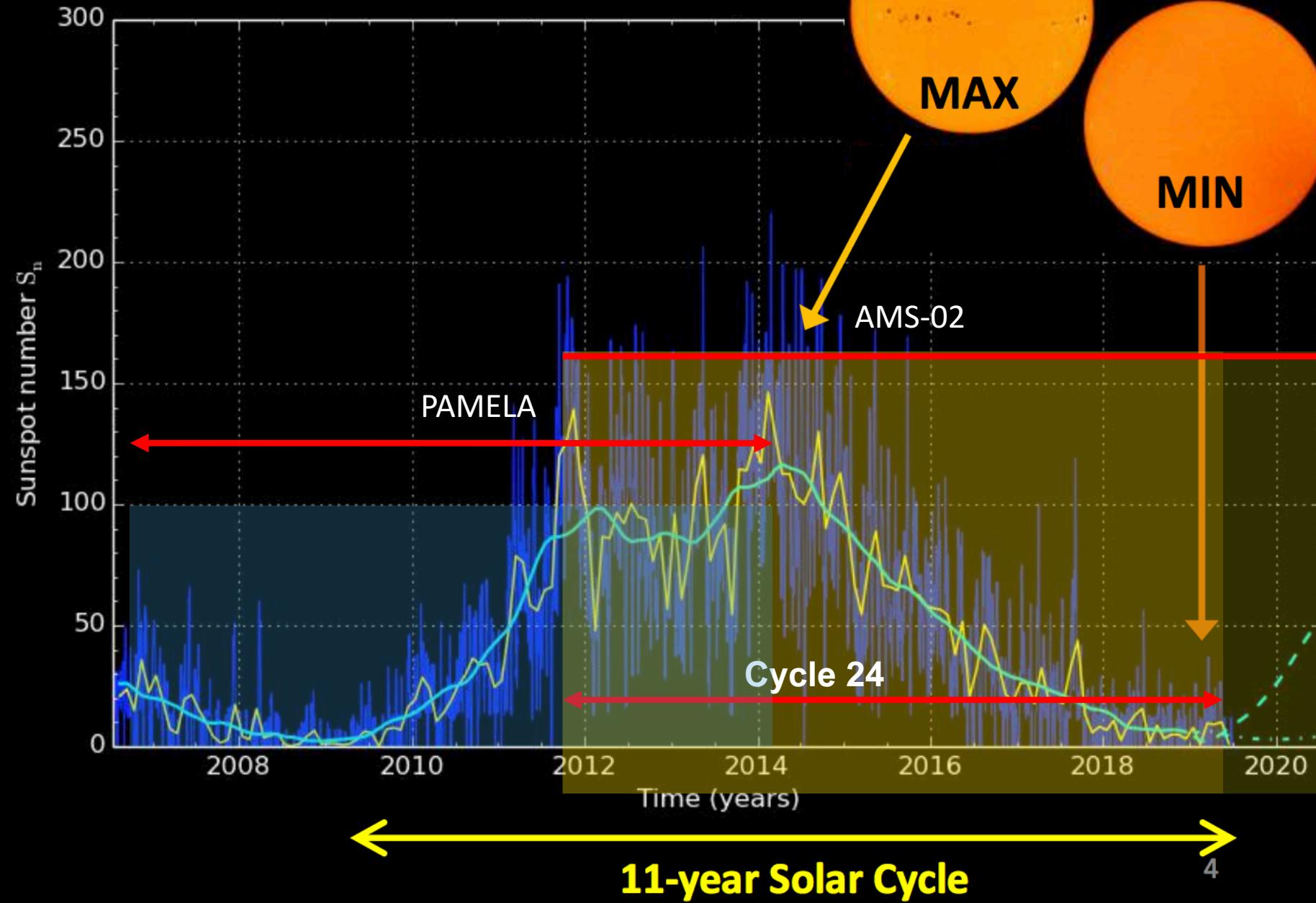
The Golden Age:

Data from Interstellar medium taken by Voyager 1 & 2

Main attributes:

- Out-flowing Solar wind
- Frozen-in turbulent Magnetic Field
- Cyclic time-dependent activity





Propagation in the heliosphere

The Parker evolution equation for the cosmic ray density

Diffusion

Small-scale Magnetic field irregularities

Drifts

Large-Scale B field, (gradients & curvature)

Sources

Galactic CRs from interstellar space

$$\frac{\partial f}{\partial t} = \nabla \cdot [K \cdot \nabla f] - V \cdot \nabla f - \langle v_D \rangle \cdot \nabla f + \frac{1}{3} (\nabla \cdot V) \frac{\partial f}{\partial \ln p} + Q(r, p, t)$$

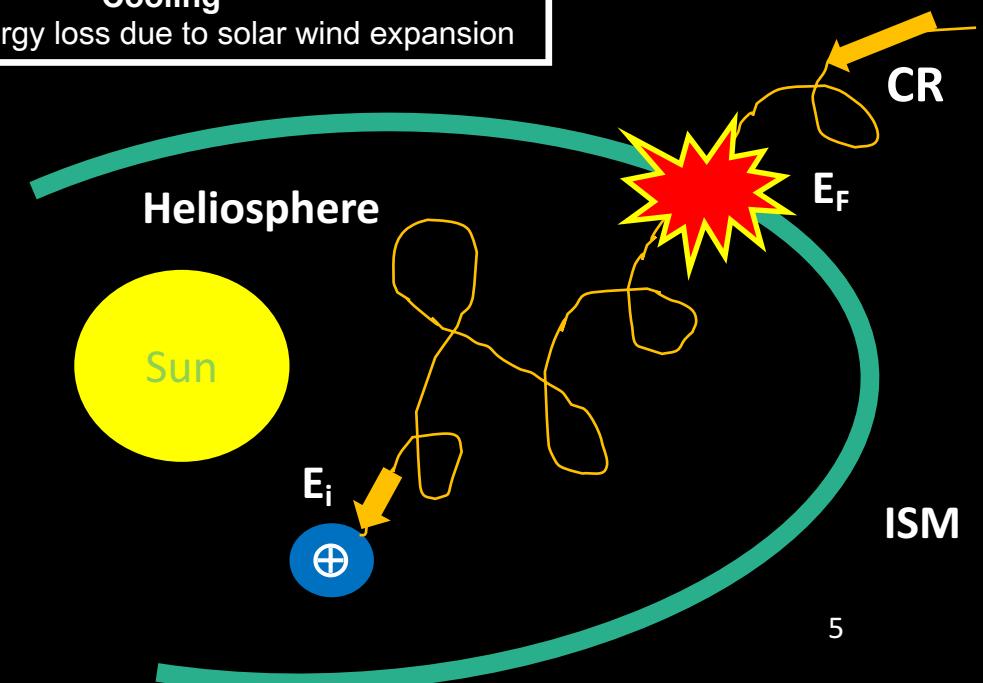
Convection

Solar wind moving out from the Sun

Cooling

Adiabatic energy loss due to solar wind expansion

- Numerical integration with $\partial/\partial t=0$ solves the Stochastic Differential Equation approach with a backward time integration based on SOLARPROP. [Kappl, Comp. Phys. Comm. 207 (2016) 386–399]



Propagation Model

The calculations

- Parker equation for the cosmic ray (CR) density in the heliosphere
- Stochastic method: Monte-Carlo simulation of CR trajectories
- Grid scan over free parameters and statistical data analysis

Key data

- Voyager-1: the CR proton flux outside the heliosphere (LIS)
- AMS-02: monthly variations of the CR proton fluxes since 2011
- PAMELA: monthly variations of the CR proton fluxes since 2006

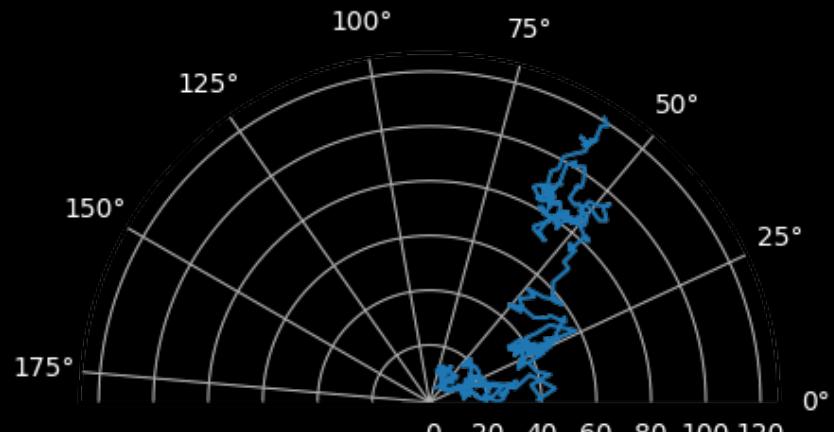
Key ingredients

1. Rigidity- and time-dependent diffusion: $K(R, t)$

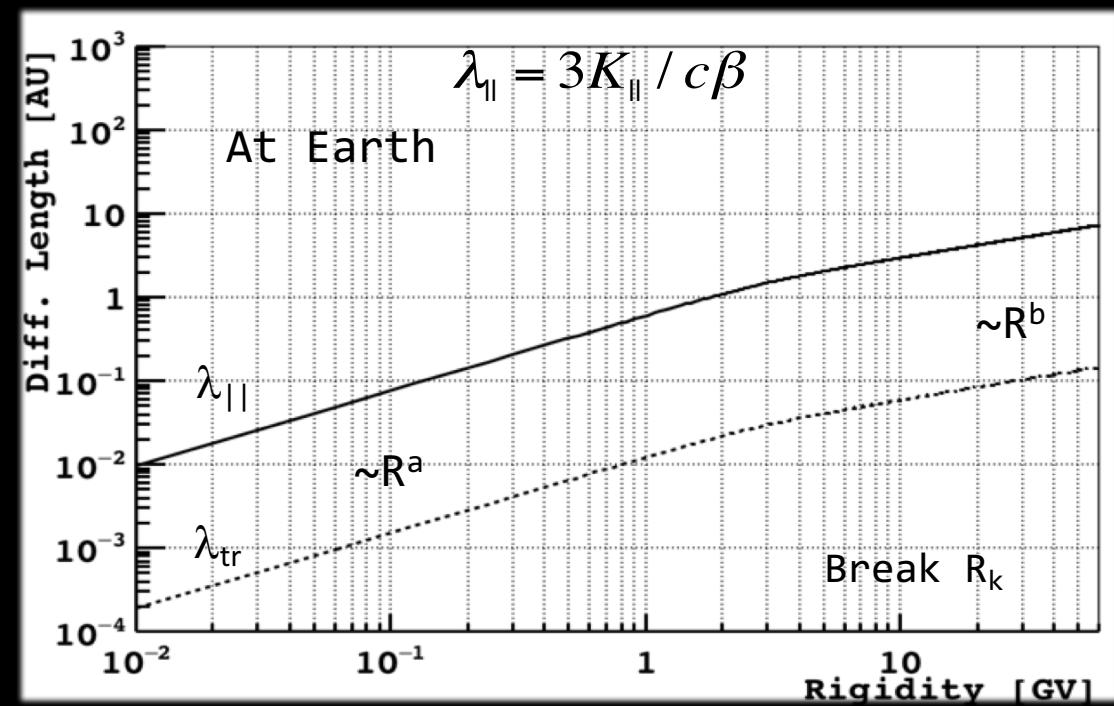
$$K(R, t) = \begin{cases} K_0(t) R_k^{a(t)} & (R_k < 3\text{GV}) \\ K_0(t) R_k^{b(t)} & (R_k > 3\text{GV}) \end{cases}$$

2. Improved 2D description of the heliosphere

3. Time-dependent B-field, current sheet, solar wind



Example: 1 GeV proton trajectory



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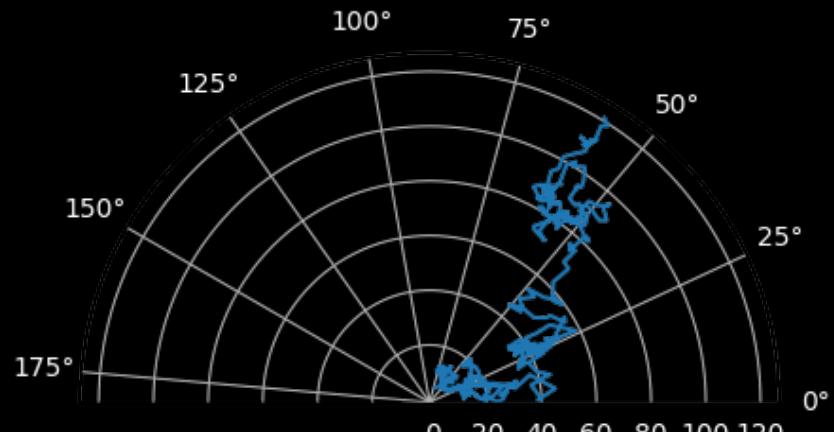
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Example: 1 GeV proton trajectory

The model has 6 parameters:

1. Tilt Angle $\alpha_T(t)$,
2. HMF at Earth $B_0(t)$,
3. Polarity of the HMF $A(t)$,

} Fix the HS status at a given epoch

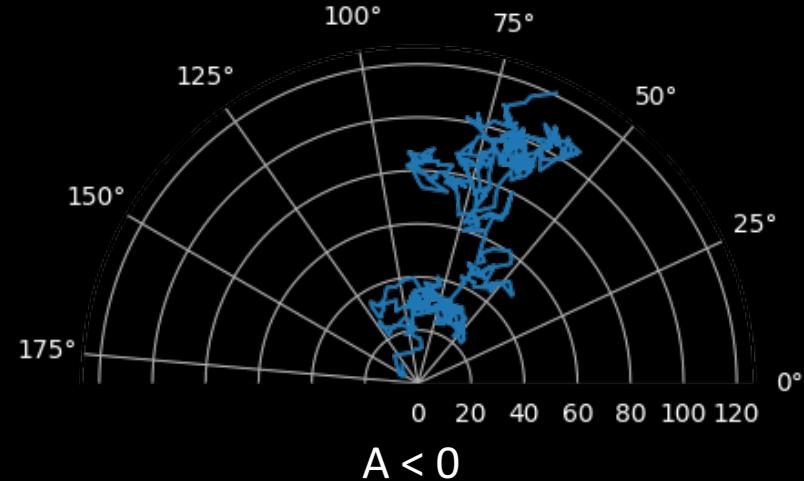
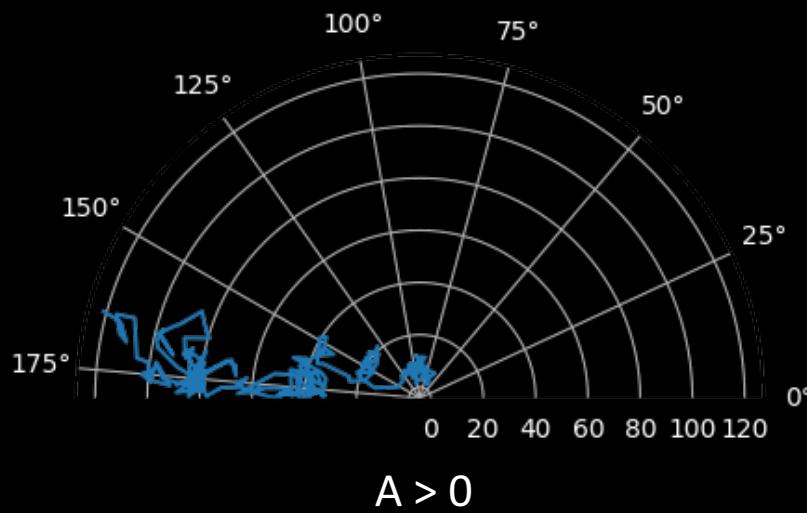
4. Diffusion coefficient normalization $K_o(t)$,
5. Diffusion spectral indices, $a(t)$ and $b(t)$, and R_k respectively
6. The break was fixed at $R_k = 3$ GV

} Fit the data

Propagation Model

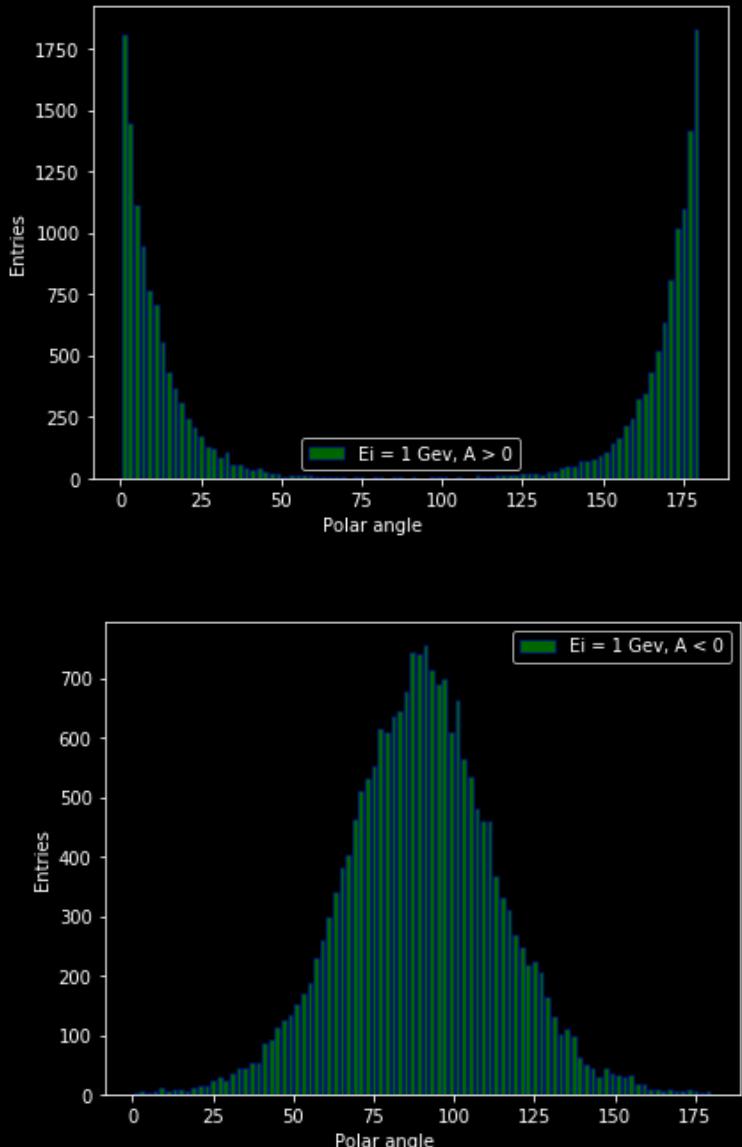
Key results

- ✓ *Data-driven determination of the key diffusion parameters as function of time*
- ✓ *Relations between the CR diffusion parameters and solar activity indices*
- ✓ *Inter-relations among the CR parameters*

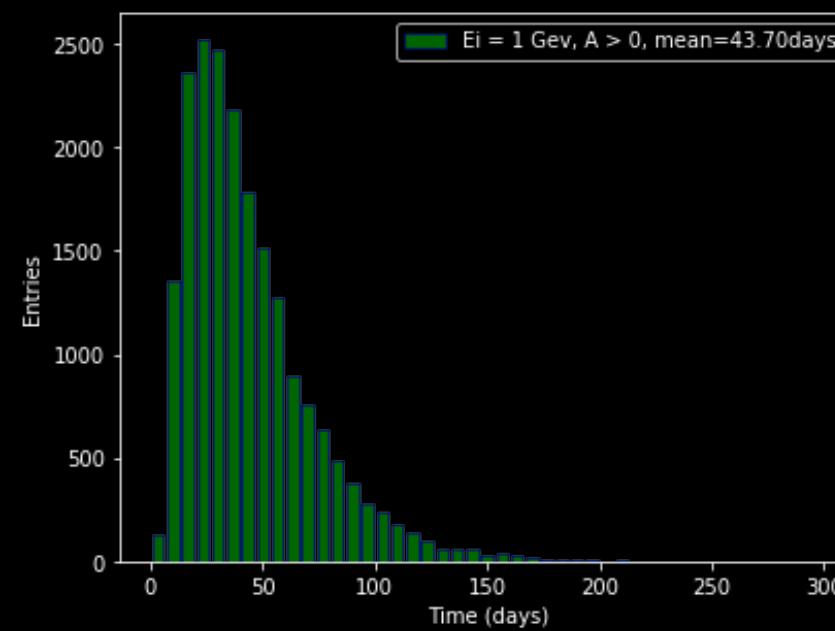
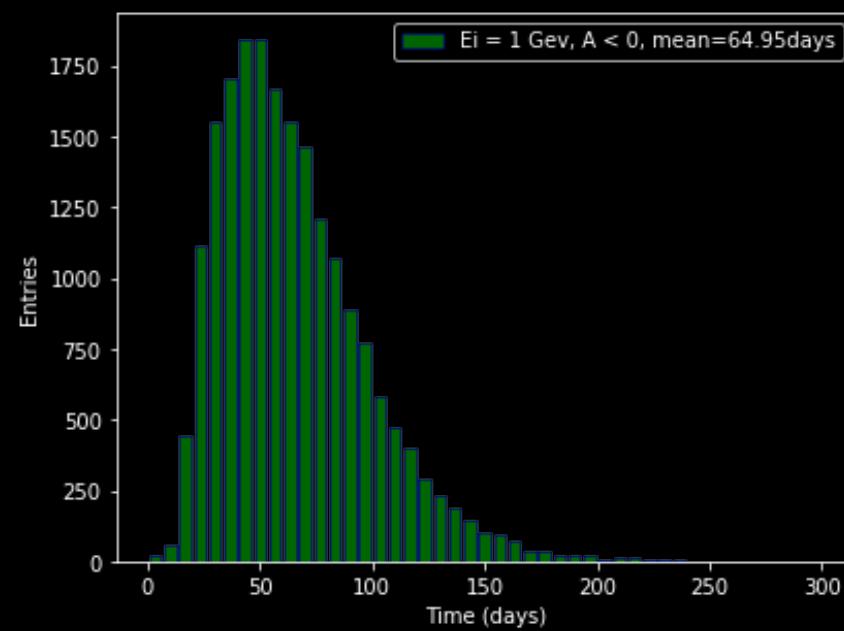


Proton trajectories in different polarities

In the $A > 0$ drift cycle, GCR protons generally drift from the polar regions to reach Earth, whereas, in the $A < 0$ cycle, they mainly drift along the Heliospheric Current Sheet

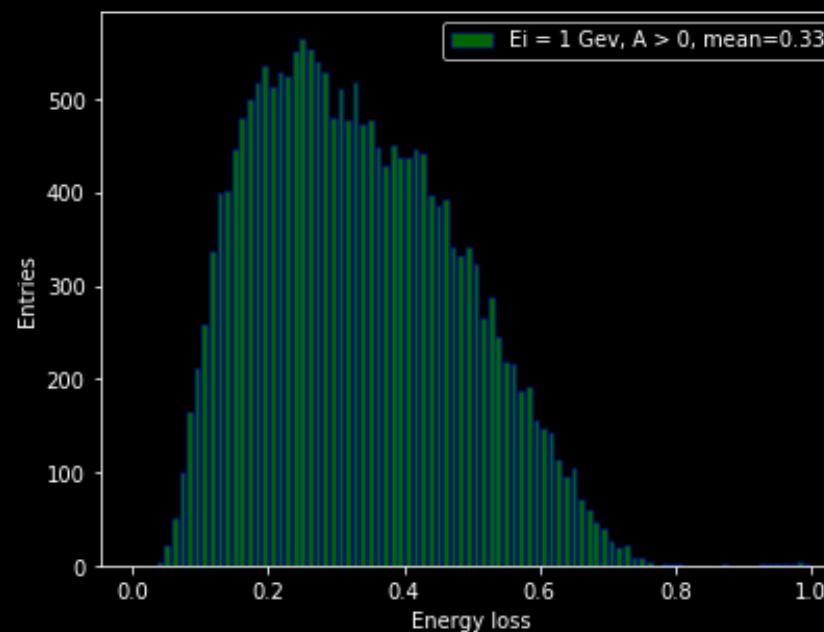
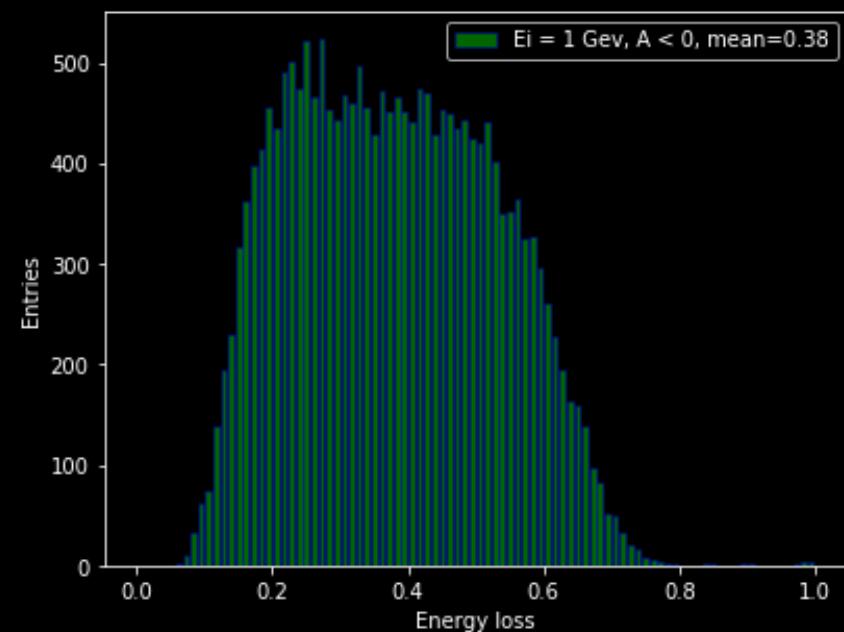


Propagation Model



Binned Propagation time

This is consistent with general drift considerations: For the $A < 0$ cycle, the protons that reach Earth have to drift inward along the heliospheric current sheet (HCS), taking a much longer time to reach Earth than protons which simply drift toward Earth from the polar regions in the $A > 0$ cycle.



Binned Energy Loss

E_F = At the boundary

E_i = Near the Earth

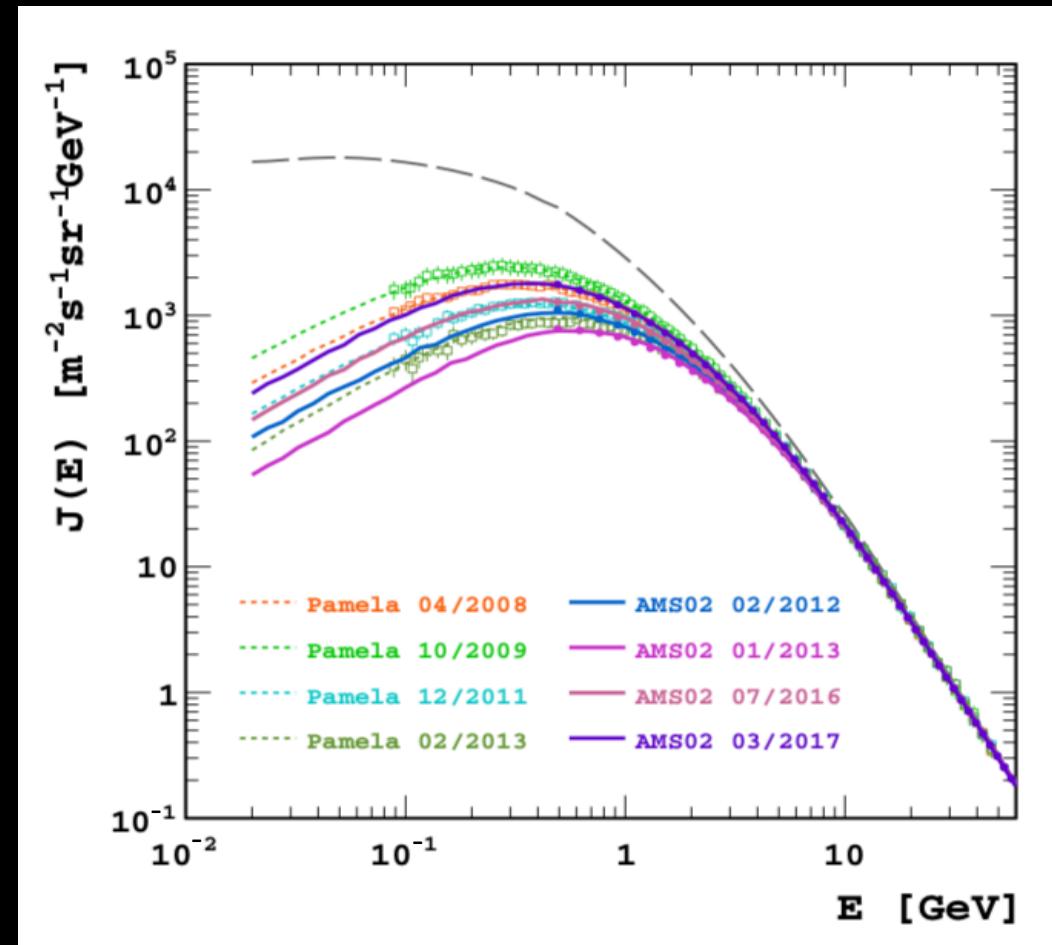
$$\text{Energy Loss} = \frac{E_F - E_i}{E_F}$$

Results: Application to AMS and PAMELA

The new precise data from AMS02 and PAMELA experiments offer a unique possibility to study the solar modulation over a long period of time

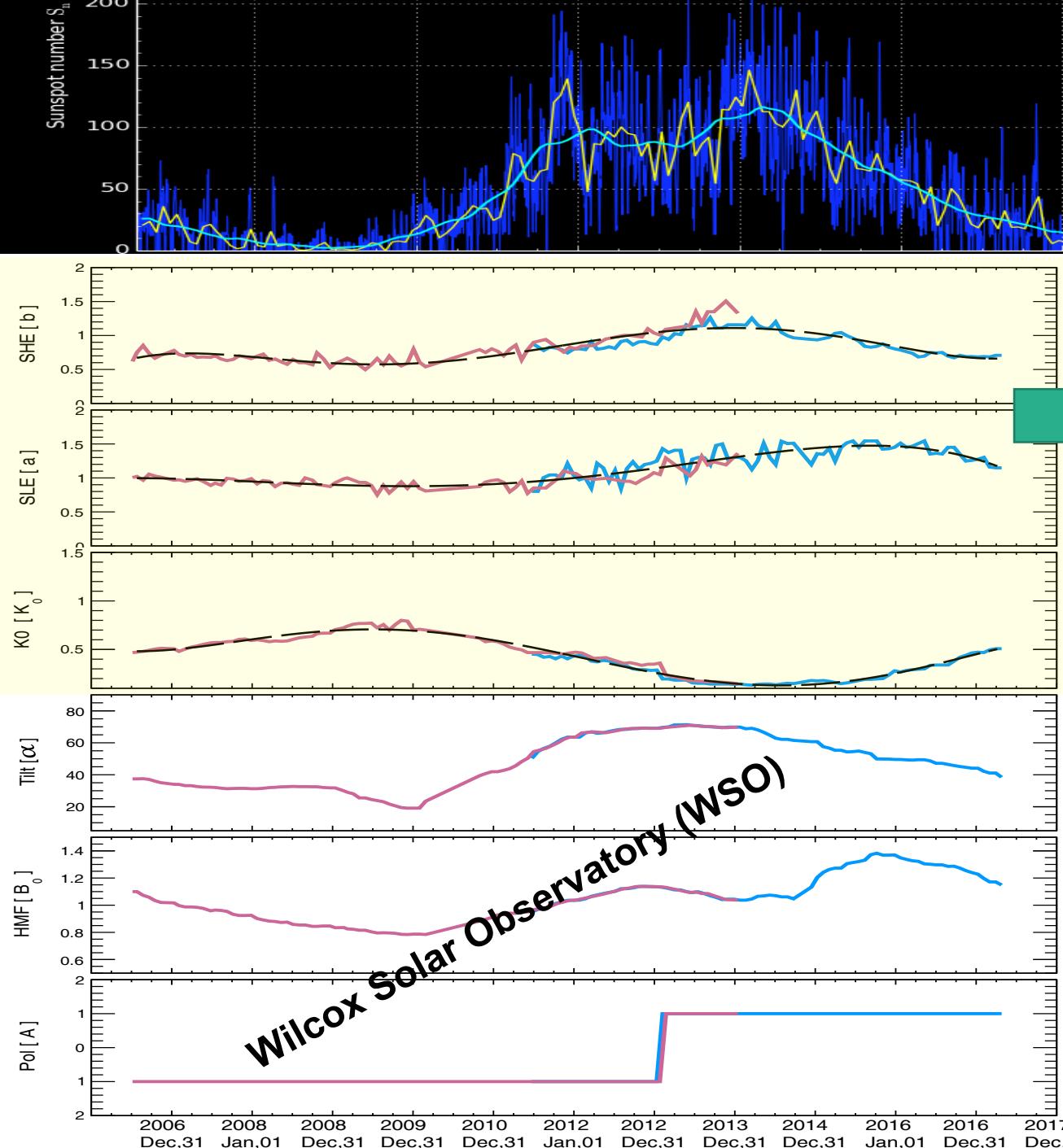
The model was applied to:

- Bartel-rotation averaged Pamela data sets from 2006/6 to 1/2014 [ApJ 742 (2011) 10, ApJ 765 (2013) 91]
- Bartel-rotation averaged AMS02 data sets from 2011/7 to 2017/6 [PRL 121 (2018) 0511]



Fiandrini et al. 2021
Phys.Rev. D. 104, 023012

Results: Application to AMS and PAMELA



A set of modulated proton fluxes were produced over the points of a discrete 6-D grid of model parameters:

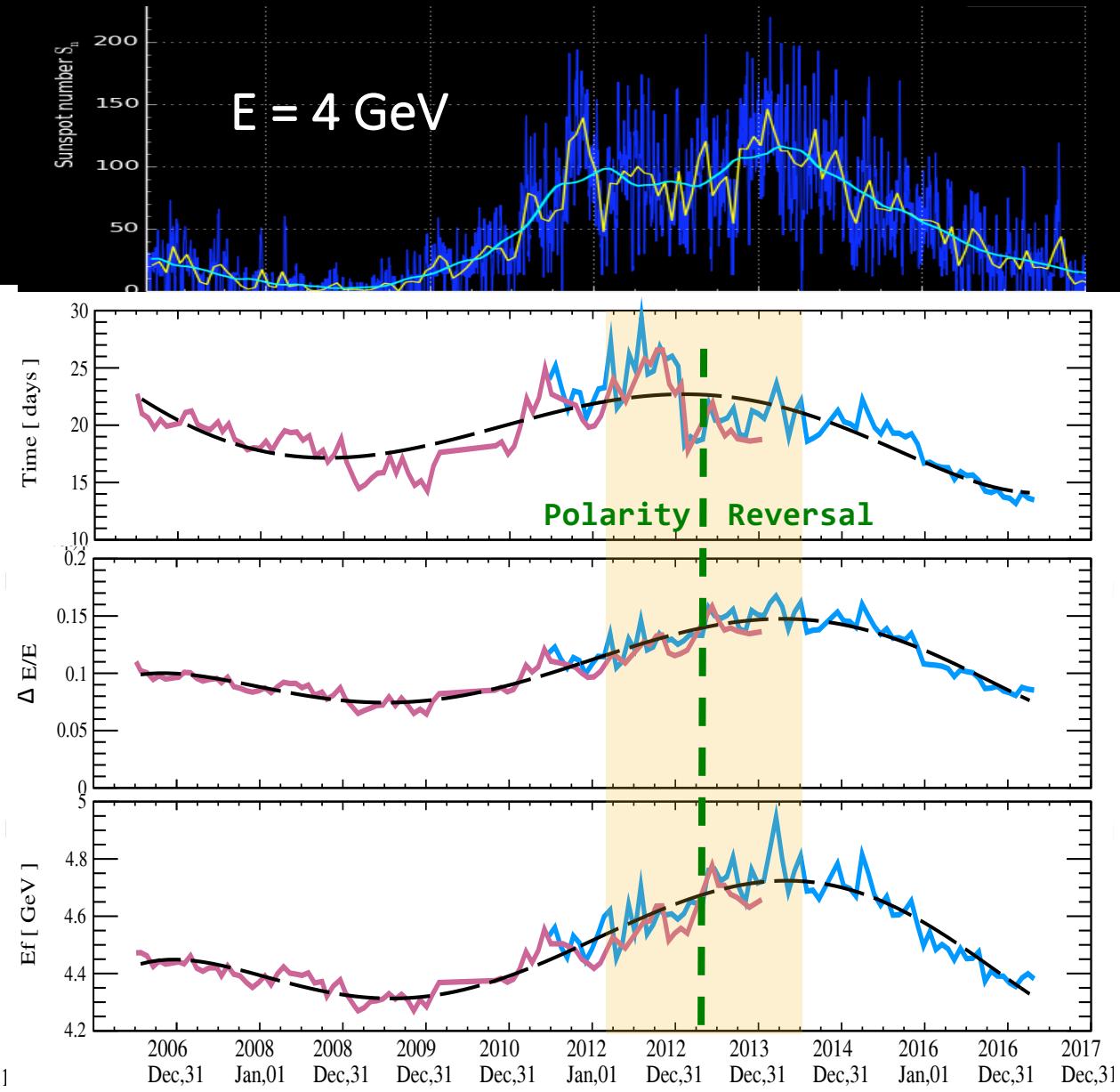
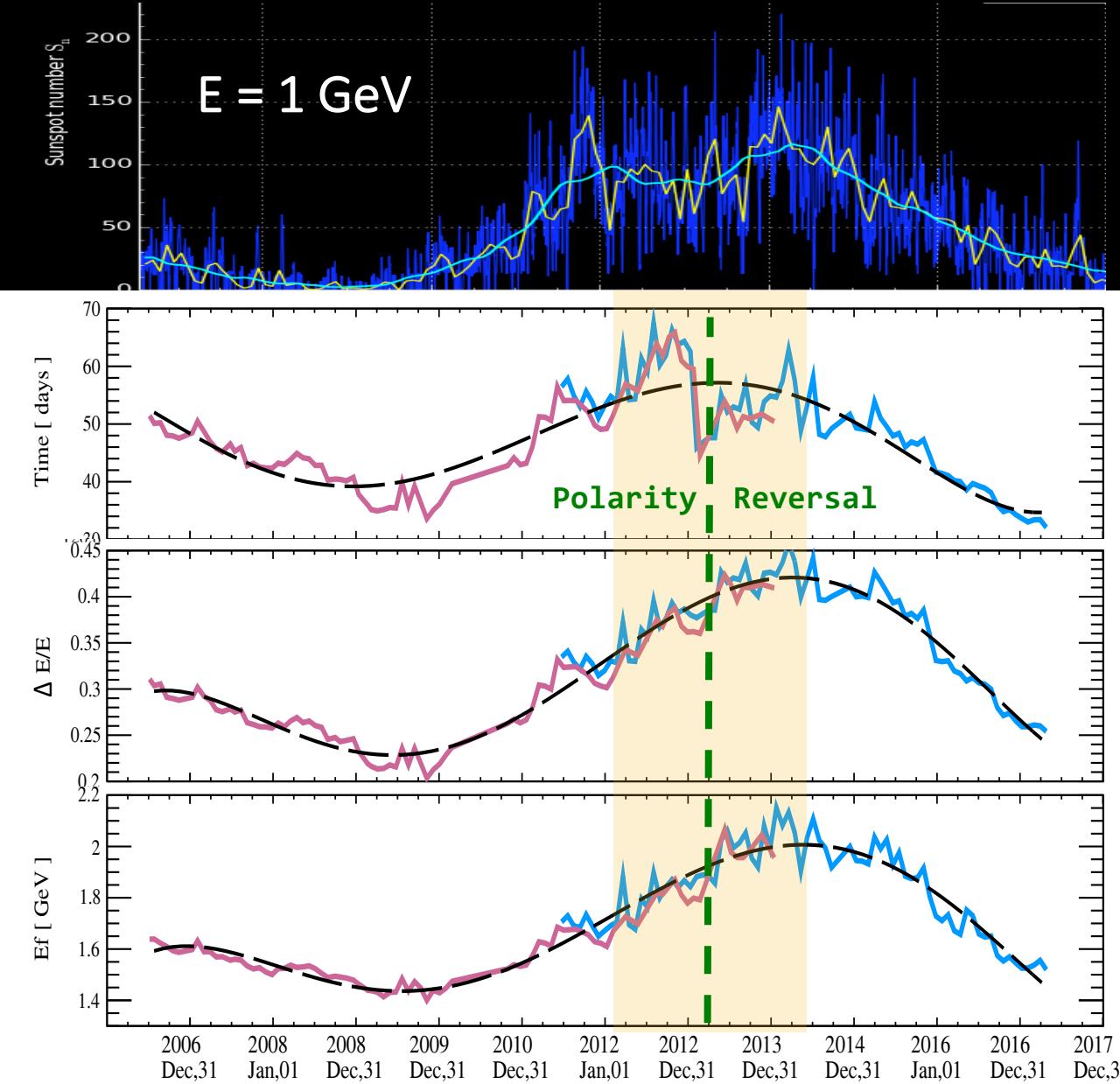
- α ($5^\circ, 75^\circ$) with step 5°
- B_0 (0.7, 1.5) in units of 5 nT with step 0.1
- Polarity A = 1 and -1
- K_0 (0.1, 0.9) with step 0.1 in units of $1.678 \times 10^{23} \text{ cm}^2\text{s}^{-1}$
- a (0.45, 1.4) with step 0.05
- b (0.45, 1.65) with step 0.05

Total of 1.215×10^6 models fluxes $J_m(E, q)$ with 2×10^3 pseudo-particles for each of the 60 energy bins from 20 MeV to 200 GeV with log step, one for each of the parameters vector q (a, B_0, A, K_0, a, b)

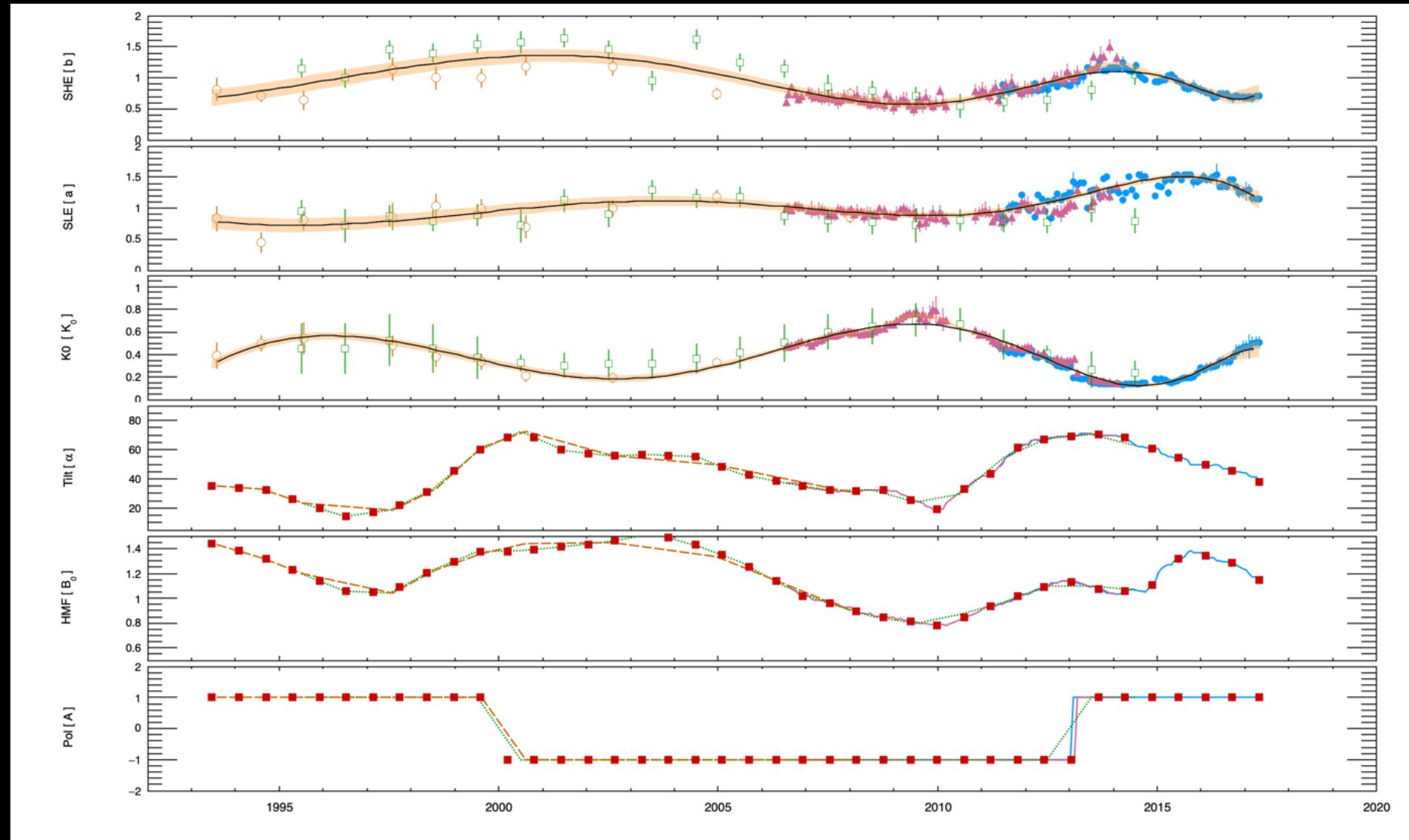


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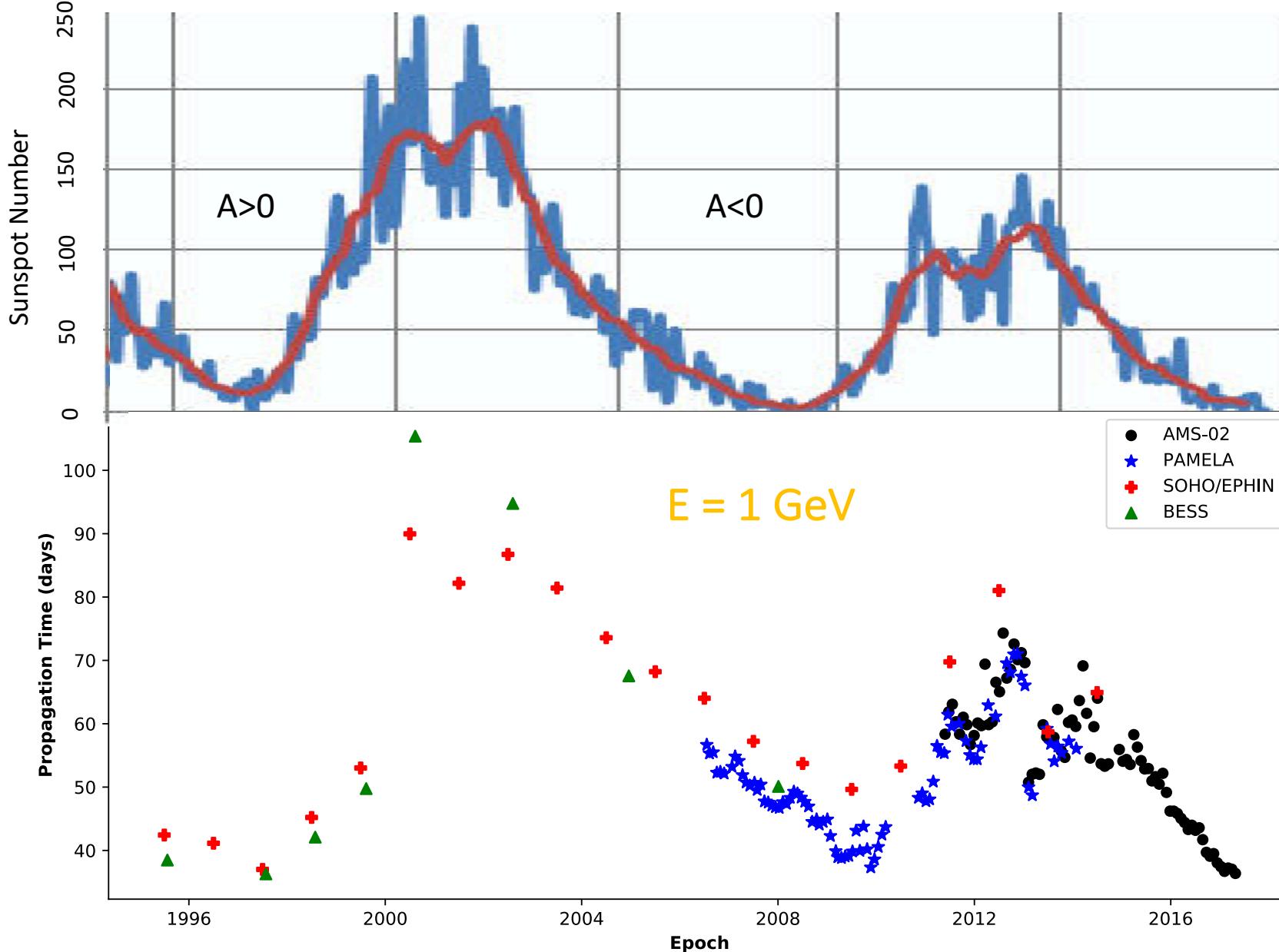


Applying the model for **Bess**, **SOHO/EPHIN**, **PAMELA**, and **AMS-02**

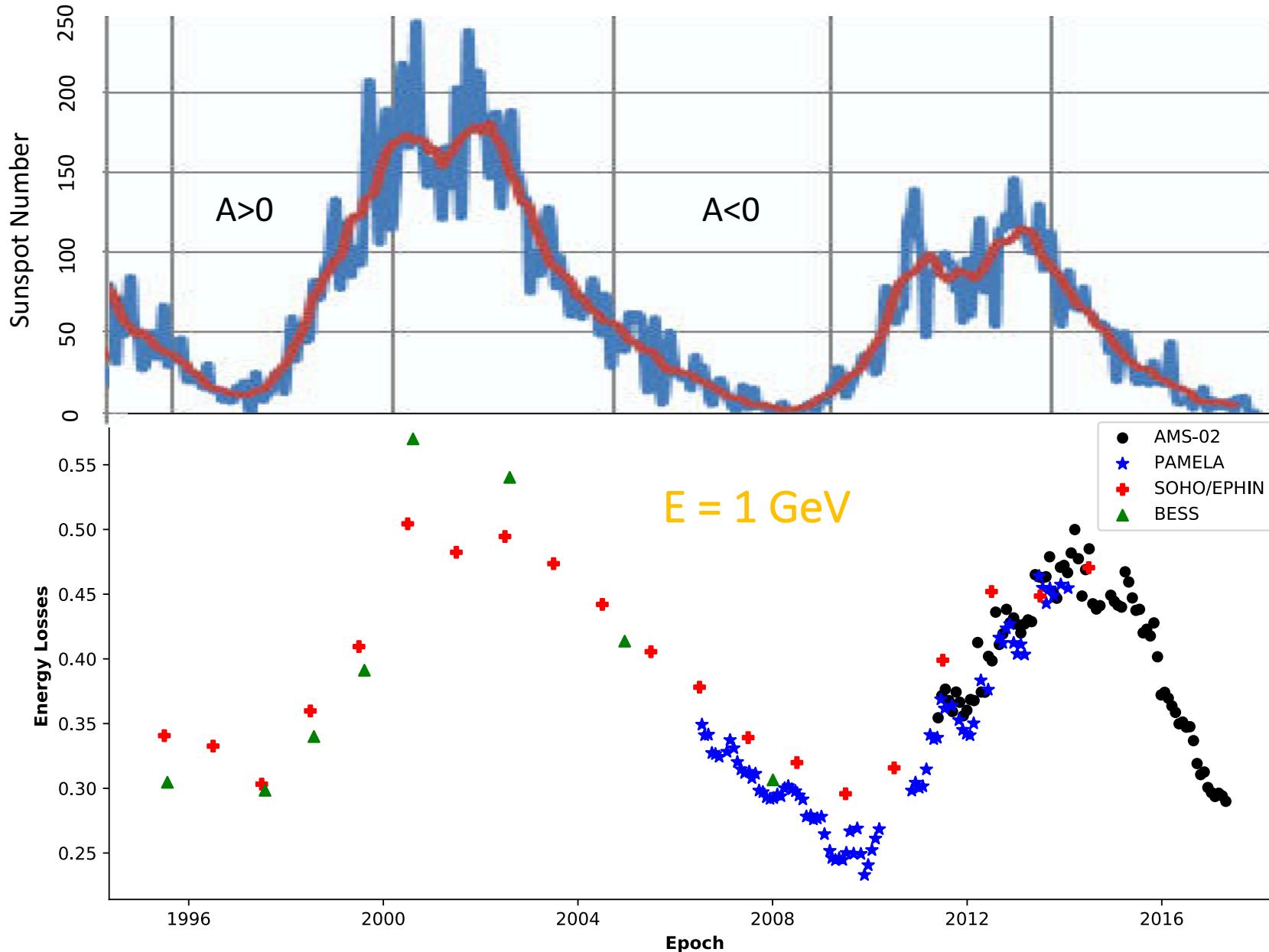


By N. Tomassetti

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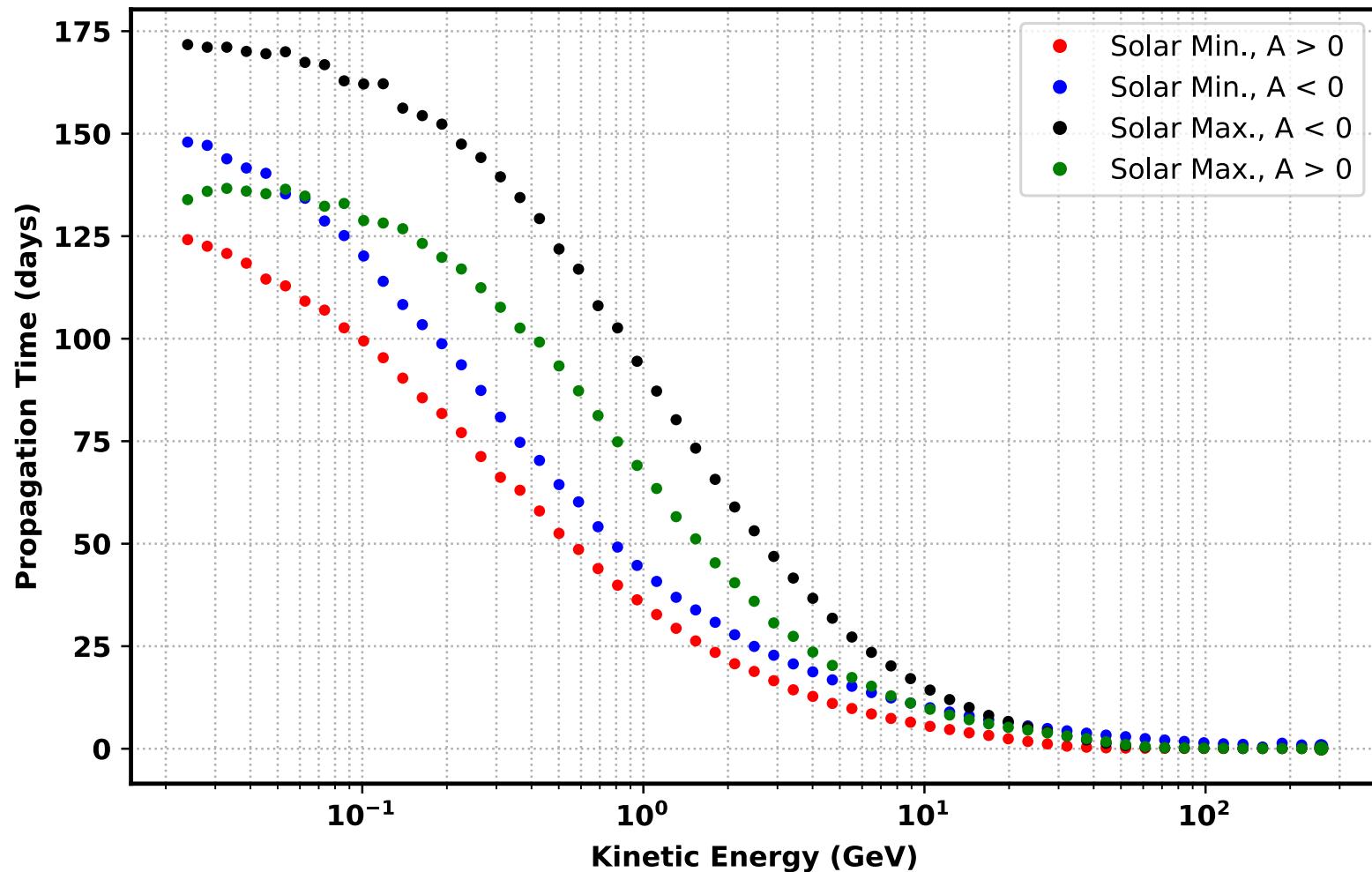
Results: Energy at Earth- Propagation Time

1997, BESS

2009, PAMELA

2002, BESS

2014, AMS-02



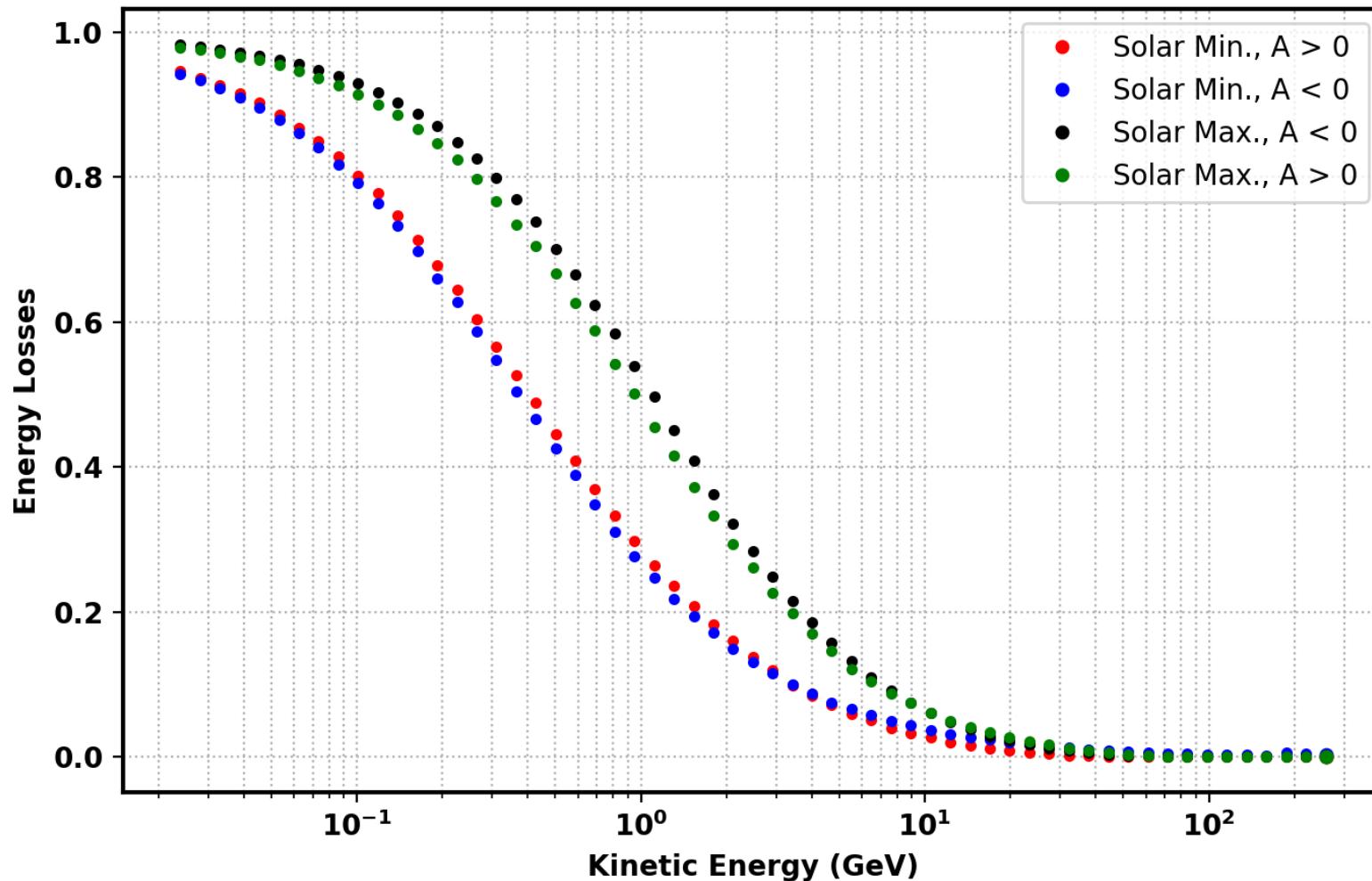
Results: Energy at Earth- Energy Losses

1997, BESS

2009, PAMELA

2002, BESS

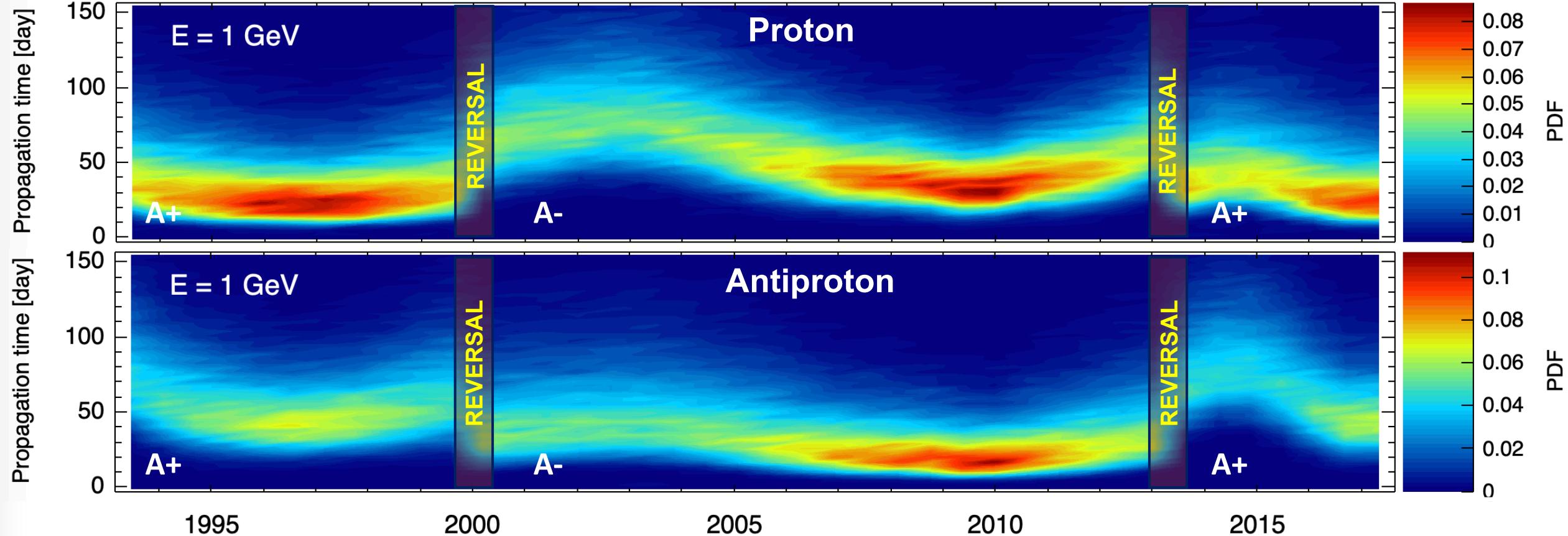
2014, AMS-02



Conclusions

- We use a stochastic model where we simulate particle trajectories.
- The model (the parameters of CR diffusion) is tuned to the observational time dependent CR fluxes: AMS-02, PAMELA, BESS, and SOHO/EPHIN, covering solar cycles
- Other physical inputs of the model are taken from direct observations of the Sun (WSO: tilt, B-field intensity, polarity)
- We are looking at particle trajectories, propagation times of CRs, and their energy losses, using our model that is global tuned over a large time interval
- The evolution of CR propagation time and their energy losses are correlated with solar activities

Back-up Slides ...



Energy loss fraction

