Evolution of sum rules and positivity constraints

QCD Evolution-2016, NIKHEF, Amsterdam, June 2, 2016

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### Main Topics

#### Kinetic form of evolution

<u>X.Artru</u>, <u>M.Elchikh</u>, <u>J.-M.Richard</u>, <u>J.Soffer</u>, <u>O.V.Teryaev</u>.
 **Phys.Rept. 470 (2009) 1-92**; <u>arXiv:0802.0164</u> [hep-ph]

- SR and positivity for DGLAP
- Energy-momentum tensor and twist-3
   Sum Rules
- TMDs and positivity

#### Kinetic form of evolution

#### Gain-loss equation (Collins, Qiu'89)

$$\frac{dq(x)}{dt} = \frac{\alpha_s}{2\pi} \int_x^1 dy \, \frac{q(y)}{y} P(x/y) \qquad P_+(z) = P(z) - \delta(1-z) \int_0^1 P(y) \, dy$$
$$\frac{dq(x)}{dt} = \frac{\alpha_s}{2\pi} \left[ \int_x^1 dy \, \frac{q(y)}{y} P(x/y) - q(x) \int_0^1 P(z) dz \right]$$

• Extra change of variables z = y/x

$$\frac{dq(x)}{dt} = \frac{\alpha_s}{2\pi} \left[ \int_x^1 dy \, \frac{q(y)}{y} P(x/y) - \int_0^x dy \, \frac{q(x)}{x} P(y/x) \right]$$
Master equation

$$\frac{\mathrm{d}q(x,t)}{\mathrm{d}t} = \int \mathrm{d}y \left[ w(y \to x)q(y,t) - w(x \to y)q(x,t) \right]$$

$$w(y \to x) = \frac{\alpha_s}{2\pi} P(x/y) \frac{\theta(y-x)}{y}$$

### Positivity preservation

- at LO as P(z) > 0 (z<1)</p>
- Positivity may be violated only by loss term, but it is decreasing when function is approaching zero
- Decreasing exponents never turns to zero; to approach shore with zero velocity boat needs infinite time (V.I. Arnold)
- At NLO scheme dependent may be used to select the scheme

#### Preservation of convexity

Kramers Moyal expansion

 $\frac{\partial q(x)}{\partial t} = -\frac{1}{x} \exp\left[\frac{\partial}{\partial \ln(1/x)} \frac{\mathrm{d}}{\mathrm{d}n}\right] x q(x) \gamma(n) \Big|_{n=1}$ 

- Commutes with log derivative curvatture preserved by evolution
- Keeping lowest terms diffusion+drift approximation  $\frac{\partial q(x)}{\partial t} = \frac{1}{x} \left( v \frac{\partial (xq(x))}{\partial \ln(1/x)} + D \frac{\partial^2 (xq(x))}{[\partial \ln(1/x)]^2} \right)$ with  $v = \frac{\alpha_s}{2\pi} (5/4 - \pi^{2/3})$ , and  $D = \frac{\alpha_s}{4\pi} (-9/8 + 2\zeta(3))$  (NSLO)

### Spin dependent case

- Use original form of DGLAP for (positive) distributions with definite helicity  $\frac{dq_{\pm}(x)}{dt} = \frac{\alpha_s}{2\pi} [P_{+\pm}(x/y) \otimes q_{+}(y) + P_{+\mp}(x/y) \otimes q_{-}(y)]$ .  $P_{+\pm}(z) = (P(z) \pm \Delta P(z))/2$
- Required :positivity of helicity kernels

 $\left|\Delta P(z)\right| \leq P(z), \quad z<1$ 

Loss term in the diagonal kernel P<sub>++</sub> only!

#### Singlet case

#### Coupled kinetic equations ('reaction diffusion system")

$$\begin{aligned} \frac{\mathrm{d}q_{\pm}(x)}{\mathrm{d}t} &= \frac{\alpha_s}{2\pi} \left[ P_{+\pm}^{qq}(x/y) \otimes q_+(y) + P_{+\mp}^{qq}(x/y) \otimes q_-(y) \right] \\ &+ P_{+\pm}^{qG}(x/y) \otimes G_+(y) + P_{+\mp}^{qG}(x/y) \otimes G_-(y) , \\ \frac{\mathrm{d}G_{\pm}(x)}{\mathrm{d}t} &= \frac{\alpha_s}{2\pi} \left[ P_{+\pm}^{Gq}(x/y) \otimes q_+(y) + P_{+\mp}^{Gq}(x/y) \otimes q_-(y) \\ &+ P_{+\pm}^{GG}(x/y) \otimes G_+(y) + P_{+\mp}^{GG}(x/y) \otimes G_-(y) \right] \end{aligned}$$

- **Positivity**  $|\Delta P^{ij}(z)| \le P^{ij}(z), z < 1; i, j = q, G$
- Loss terms only in the diagonal kernels

Transversity: Soffer bound stability

Positive quantities (of mixed chirality)

 $Q_{\pm}(x) = q_{\pm}(x) \pm h_1(x)$ 

Kinetic equations

$$\frac{\mathrm{d}Q_{\pm}(x)}{\mathrm{d}t} = \frac{\alpha_s}{2\pi} (P_{+\pm}^Q(x/y) \otimes Q_+(y) + P_{+\mp}^Q(x/y) \otimes Q_-(y))$$

$$P_{++}^Q(z) \equiv \frac{P_{qq}^{(0)}(z) + P_h^{(0)}(z)}{2} = \frac{C_F}{2} \left[ \frac{(1+z)^2}{(1-z)_+} + 3\delta(1-z) \right]$$

$$P_{+-}^Q(z) \equiv \frac{P_{qq}^{(0)}(z) - P_h^{(0)}(z)}{2} = \frac{C_F}{2} (1-z).$$

 Kernels: positive (LO) for z < 1, loss terms diagonal

### "Scale arrow"

- Evolution is kinetic when going to UV and "antikinetic" in the IR
- Positivity preserved when going to UV
- Evolving backwards: small deviations in the UV explode in the IR
- "Turbulence" in the confinement region
- Soffer bound (for d) may be saturated at low scale only: otherwise backward evolution would violate it (M.Radici talk)



#### Current (~"particles" number) conservation

$$\int_0^1 dx \, \frac{dq(x)}{dt} = \int_0^1 \int_0^1 dx dy \, [w(y \to x)q(y) - w(x \to y)q(x)] = 0$$

 Singlet case – EMT conservation: 2 sorts of particles (numbers ~ xq(x), xG(x)) mutually transforming one to another ("Mass conservation in reaction-diffusion system") Energy-momentum tensor related sum rules

- PDF's momentum SR, first indication for gluons
- GPDs Ji's SRs
- Follow from momentum and angular momentum conservation
- Evolution of angular momenta the same as momenta may be obtained in the kinetic way (OT'98) if  $\int_0^1 dxx \Delta P_{Gq}(x) = \frac{1}{2} \int_0^1 dxx P_{Gq}(x)$

## Gluonic poles and Energy-Momentum tensor

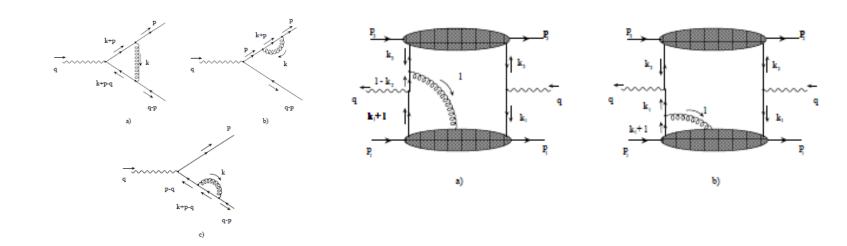
- Consider twist 3 (= relevant moment of Sivers function being infinite tower of twists) gluonic pole
- EMT forward matrix element  $\langle P | T_{\mu\nu} | P \rangle = 2P_{\mu}P_{\nu}$ , no spindependent structure  $P_{\mu}\mathcal{E}_{vSPn}$  (similar to B~E with q -> n) :

$$\sum \int \int dx_1 dx_2 \frac{T(x_1, x_2)}{x_1 - x_2} = 0$$

- Naively: Valid identically due to symmetry properties implied by T-invariance)
- However: such pole in physical processes should get imaginary part due to EMGI and related contour gauge(DY process: Anikin, OT, PLB2010,2015; EPJ2015) – analog of the choice of Wilson line.

Electromagnetric Gauge Invariance in DY process

- Extra diagram factor 2 in transverse (TM integrated) asymmetry
- Follows also from EM GI
- May be studied at COMPASS, NICA 1408.3959



# Pole prescription and Burkardt SR (OT'14)

Pole prescription provides ("T-odd") symmetric part!  $\sum \int \int dx_1 dx_2 \frac{T(x_1, x_2)}{x_1 - x_2 + i\varepsilon} = 0$ 

SR: 
$$\sum \int dx T(x,x) = 0$$

- Burkardt SR+Boer-Mulders-Pijlman relation
- Pole prescription way to account for dynamics
- Cf with analogous SR (Schafer,OT'01) for Collins functions where only TM conservation was necessary

# Validity for separate parton species

- Can it be valid separately for each quark flavour (and gluons) : nodes (Boer, Prokudin)?
- Valid if structures P<sub>µ</sub>E<sub>vSPn</sub> forbidden for TOTAL conserved EMT do not appear for each flavour
- Structure contains (besides S) gauge vector n: GI (and/or rotational invariant) separation of EMTs – forbidden: SR valid separately!

Direct test of Twist-3 sum rule evolution (J. Zhou'15 - talk)

Multiplicative evolution:

$$\frac{\partial \sum_{q,\bar{q},g} [T_F^q(\mu^2) + T_F^{\bar{q}}(\mu^2) + T_G^{(+)}(\mu^2)]}{\partial \ln \mu^2} = -\frac{\alpha_s}{2\pi} C_A \sum_{q,\bar{q},g} [T_F^q(\mu^2) + T_F^{\bar{q}}(\mu^2) + T_G^{(+)}(\mu^2)]$$

#### Satisfied IF valid at some scale

 R.h.s. from "extra terms" (now confirmed) of Braun, Manashov and Pirnay Comparing Burkardt and Burkhardt-Cottingham SRs

- Twist 3 contribution to BC SR also evolve multiplicatively in large N<sub>C</sub> limit (Ali, Braun, Hiller; Braun, Korchemsky, Manashov)
- BC is related to rotational invariance
- Momentum SR -> BSR: transition from longitudinal to transverse
- BSR ~ JiSR (OT'06) where ANGULAR momentum conservation is required

"Spontaneous" conservation of Burkardt SR

- Pure non-Abelian rotational properties due to gluon self-interaction?
- Straightforward generalization: If valid separately for each flavour and gluons at some points also remains stable!

### Positivity and TMD

- Low-x BFKL evolution
- Master-type form for UGDF f(x<sub>B</sub>,k<sub>T</sub>) with longitudinal time t=ln(1/x<sub>B</sub>) and transverse coordinate x = ln k<sub>T</sub>

$$\frac{\mathrm{d}f(x,t)}{\mathrm{d}t} = \int \mathrm{d}y \left[ w_+(y \to x)f(y,t) - w_-(x \to y)f(x,t) \right]$$

- Contains exponential growth besides diffusion and drift
- It is possible to separate these effects

#### **BFKL** as a master equation

#### Redefined function and kernel

 $f_{\sigma}(x,t) = f(x,t)\sigma(x), \quad \bar{f}_{\sigma}(x,t) = f_{\sigma}(x,t)/\langle f_{\sigma}(t) \rangle \quad \langle f_{\sigma}(t) \rangle = \int dx f_{\sigma}(x,t).$ 

 $w_{+}(x, y) \rightarrow w_{\sigma}(x, y) = w_{+}(x, y)\sigma(x)/\sigma(y)$ 

• Master equation for  $\overline{f_{a}(\mathbf{x},t)}$ 

$$\frac{\mathrm{d}\bar{f}_{\sigma}(\mathbf{x},t)}{\mathrm{d}t} = \int \mathrm{d}\mathbf{y} \left[ w_{\sigma}(\mathbf{y} \to \mathbf{x}) f_{\sigma}(\mathbf{y}) - w_{\sigma}(\mathbf{x} \to \mathbf{y}) \bar{f}_{\sigma}(\mathbf{x}) \right]$$

•  $\sigma(x)$ -eigenfunction (BFKL:  $(k_T)^a$ , 0 < a < 1)
Of  $\int dx w_+ (y \to x) \sigma(x) = \left[\lambda_\sigma + \int dx w_- (y \to x)\right] \sigma(y)$ • Growth power (  $\lambda_0 = \lambda_\sigma - \frac{v_\sigma^2}{4D_\sigma} = 4 \ln 2$ )  $\lambda_\sigma = \int dx [w_\sigma (y \to x) - w_- (y \to x)]$ 

# Positivity for BFKL and its extensions

- Master form preserved positivity
- Scale arrow directed towards low x
- Unitarization: Nonlinear local terms do not affect positivity
- Coordinate space: BK preserve positivity
- But: Fourier transform of positive functions is only positive definite: for any real x and complex z  $\sum f(x_i-x_j)z_iz_j^* > 0$

# Positivity and TMD factorization

- W term must turn negative to have zero moment – violation of positivity (in physical momentum space!) signals of inapplicability of approximation
- Two scale arrows for CS and scale evolution
- Positivity preserved for
- Possible qualitative description: reversal of transverse scale arrow (change of log sign) at Q<sub>T</sub>~Q and its restoration by change W -> Y?

## **Conclusions/Outlook**

- Kinetic interpretation of evolution naturally describes positivity and sum rules preservation
- Scale arrows directed towards large p<sub>T</sub> and small x (combined by angle arrow?)
- Burkardt SR in twist 3 approach is controlled by energy-momentum conservation + dynamics (pole prescription)
- Spontaneous conservation of Burkardt SR and its generalization for seprate flavours
- TMD evolution: scale arrow reversal at  $Q_T \sim Q$ ?

Can "standard" conservation of BSR be imposed?

- Twist 3 perfectly survives Abelian limit
- "Extra terms" are pure non-Abelian
- Small x (IR) effect: could it be the room for extra subtraction?
- Recall axial anomaly (Carlitz,Collins, Mueller)
   correct IR limits crucial
- Pairs of anomalies: V vs A, EMT conservation vs Trace
- Could the subtraction be related to trace anomaly?!

## 1-st moments - EM, 2-nd -Gravitational Formfactors

 $\langle p'|T^{\mu\nu}_{q,g}|p\rangle = \bar{u}(p') \Big[ A_{q,g}(\Delta^2) \gamma^{(\mu} p^{\nu)} + B_{q,g}(\Delta^2) P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}/2M ] u(p)$ 

Conservation laws - zero Anomalous Gravitomagnetic Moment :  $\mu_G = J$  (g=2)

$$\begin{split} P_{q,g} &= A_{q,g}(0) & A_q(0) + A_g(0) = 1 \\ J_{q,g} &= \frac{1}{2} \left[ A_{q,g}(0) + B_{q,g}(0) \right] & A_q(0) + B_q(0) + A_g(0) + B_g(0) = 1 \end{split}$$

- May be extracted from high-energy experiments/NPQCD calculations
- Describe the partition of angular momentum between quarks and gluons
- Describe ainteraction with both classical and TeV gravity

### Electromagnetism vs Gravity

#### Interaction – field vs metric deviation

- $M = \langle P' | J^{\mu}_{q} | P \rangle A_{\mu}(q) \qquad \qquad M = \frac{1}{2} \sum_{q,G} \langle P' | T^{\mu\nu}_{q,G} | P \rangle h_{\mu\nu}(q)$
- Static limit

 $\langle P|J^{\mu}_{q}|P\rangle = 2e_{q}P^{\mu}$ 

$$\sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle = 2P^{\mu}P^{\nu}$$
$$h_{00} = 2\phi(x)$$

$$M_0 = \langle P | J^{\mu}_q | P \rangle A_{\mu} = 2e_q M \phi(q) \qquad M_0 = \frac{1}{2} \sum_{q,G} \langle P | T^{\mu\nu}_i | P \rangle h_{\mu\nu} = 2M \cdot M \phi(q)$$

Mass as charge – equivalence principle

### Equivalence principle

- Newtonian "Falling elevator" well known and checked
- Post-Newtonian gravity action on SPIN known since 1962 (Kobzarev and Okun') – not checked on purpose but in fact checked in atomic spins experiments at % level (Silenko,OT'07)
- Anomalous gravitomagnetic moment iz ZERO or
- Classical and QUANTUM rotators behave in the SAME way

#### Gravitomagnetism

Gravitomagnetic field – action on spin – ½ from  $M = \frac{1}{2} \sum_{q,G} \langle P' | T^{\mu\nu}_{q,G} | P \rangle h_{\mu\nu}(q)$ 

$$\vec{H}_J = \frac{1}{2} rot \vec{g}; \ \vec{g}_i \equiv g_{0i}$$
 spin dragging twice  
smaller than EM

- Lorentz force similar to EM case: factor  $\frac{1}{2}$ cancelled with 2 from  $h_{00} = 2\phi(x)$ Larmor frequency same as EM  $\vec{H}_L = rot\vec{g}$
- Orbital and Spin momenta dragging the same Equivalence principle  $\omega_J = \frac{\mu_G}{J}H_J = \frac{H_L}{2} = \omega_L$

# Equivalence principle for moving particles

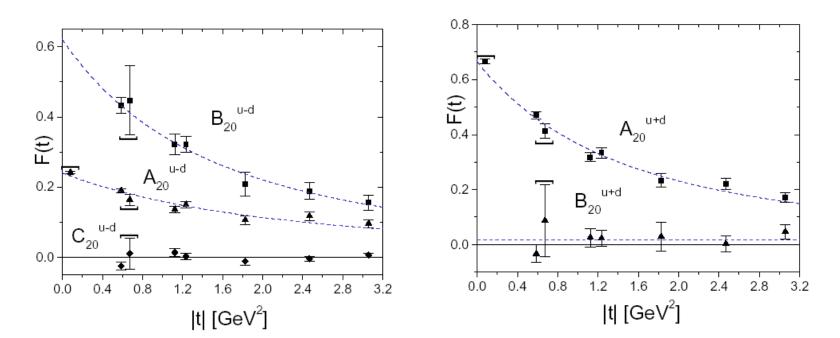
- Compare gravity and acceleration: gravity provides EXTRA space components of metrics h<sub>zz</sub> = h<sub>xx</sub> = h<sub>yy</sub> = h<sub>00</sub>
- Matrix elements DIFFER

 $\mathcal{M}_g = (\epsilon^2 + p^2) h_{00}(q), \qquad \mathcal{M}_a = \epsilon^2 h_{00}(q)$ 

Ratio of accelerations:  $R = \frac{\epsilon^2 + p^2}{\epsilon^2}$  - confirmed by explicit solution of Dirac equation (Silenko, O.T.)

# Generalization of Equivalence principle

Various arguments: AGM ≈ 0 separately for quarks and gluons – most clear from the lattice (LHPC/SESAM)



Extended Equivalence Principle=Exact EquiPartition

- In pQCD violated
- Reason in the case of EEP- no smooth transition for zero fermion mass limit (Milton, 73)
- Conjecture (O.T., 2001 prior to lattice data) – valid in NP QCD – zero quark mass limit is safe due to chiral symmetry breaking
- Supported by smallness of E (isoscalar AMM)
- Polyakov Vanderhaeghen: dual model with E=0

## EEP and AdS/QCD

- Recent development calculation of Rho formfactors in Holographic QCD (Grigoryan, Radyushkin)
- Provides g=2 identically!
- Experimental test at time –like region possible