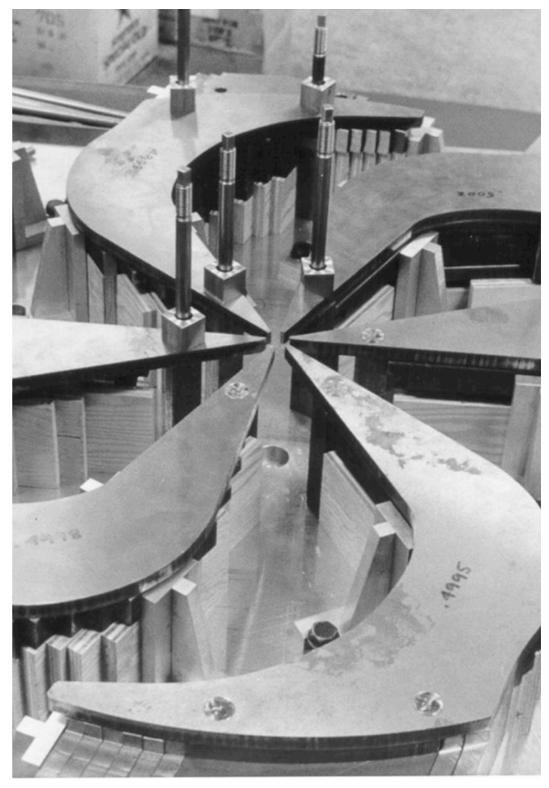
%TRIUMF

Electroweak baryogenesis, experimental status, progress and extensions

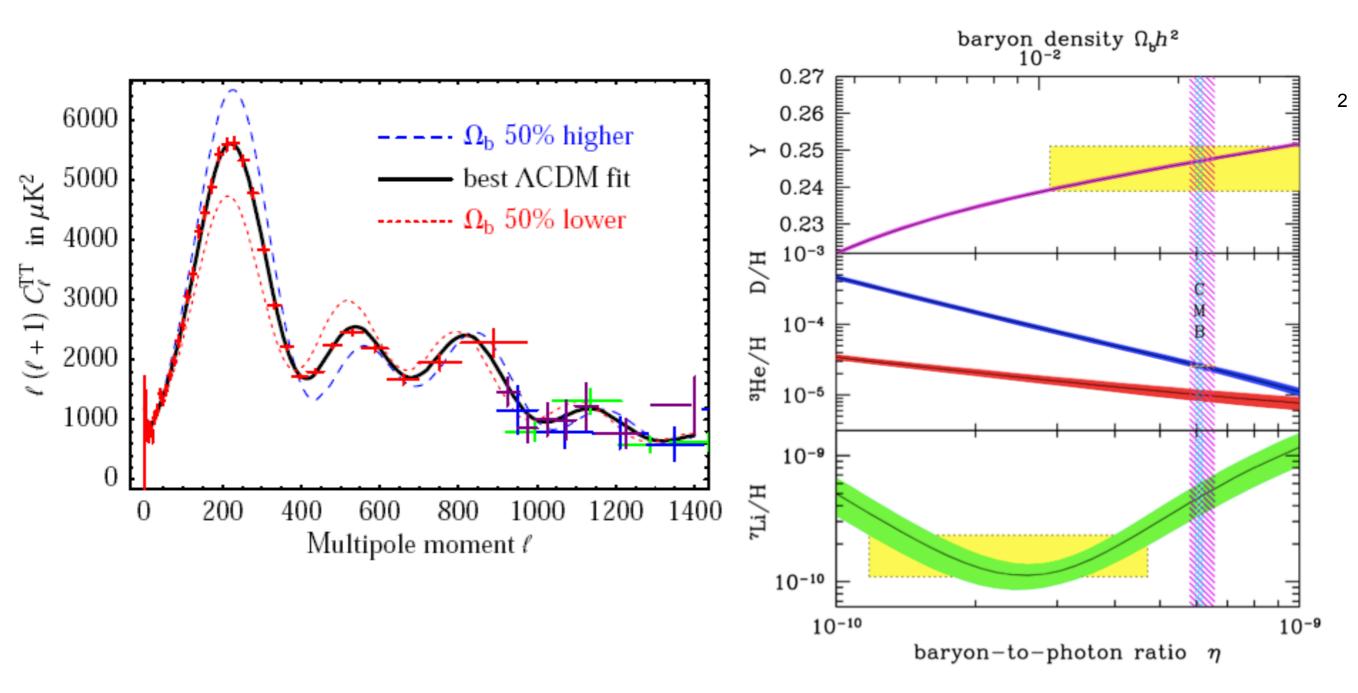
with Csaba Balazs, Michael Ramsey Musolf, David Morrissey, Seyda Ipek, Sebastian Ellis, Joydeep Roy, Jorinde Van de Vis, Jordy De Vries, Marieke Postma, Sujeet Akula, Liam Duncan, Peter Athron, Michael Bardsley, Peter Winslow, Djuna Croon, Toumas Tenkanen and Oli Gould



Discovery, accelerate

1

Baryon asymmetry



Ng. K. Francis (thesis)

Particle data group

Washout and Sakharov conditions

Inflation washes out initial baryon asymmetry (almost)

Sakharov conditions

- C and CP violation
- Violation of BNC
- Departure from Equilibrium*

How it fulfils the Sakharov conditions

Inflation washes out initial baryon asymmetry (almost)

Sakharov conditions

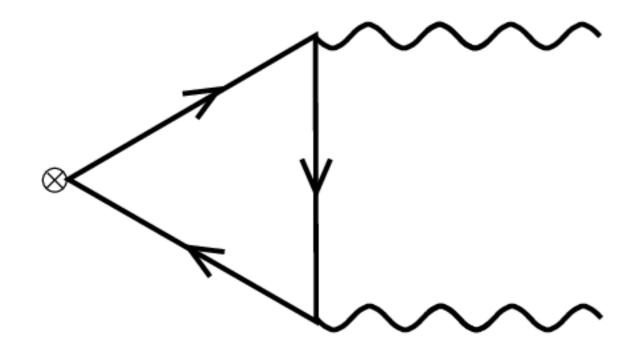
- C and CP violation
- Violation of BNC
- Departure from Equilibrium*
- -- CPV interactions with Higgs
- <—— Non-perturbative processes (sphalerons)</p>

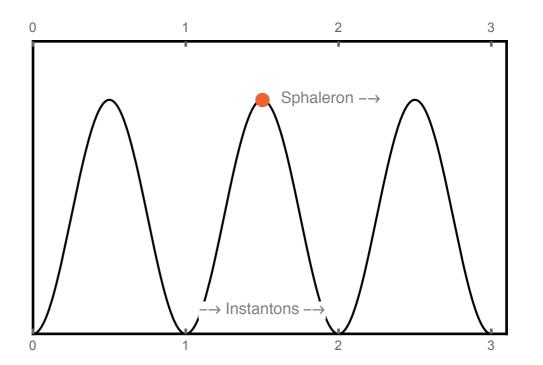
4

-- Electroweak phase transition

BNC violation: sphalerons

 $\partial_{\mu}J^{\mu}_{B+L} \sim W\tilde{W}$ $\partial_{\mu}J^{\mu}_{5} \sim G\tilde{G}$

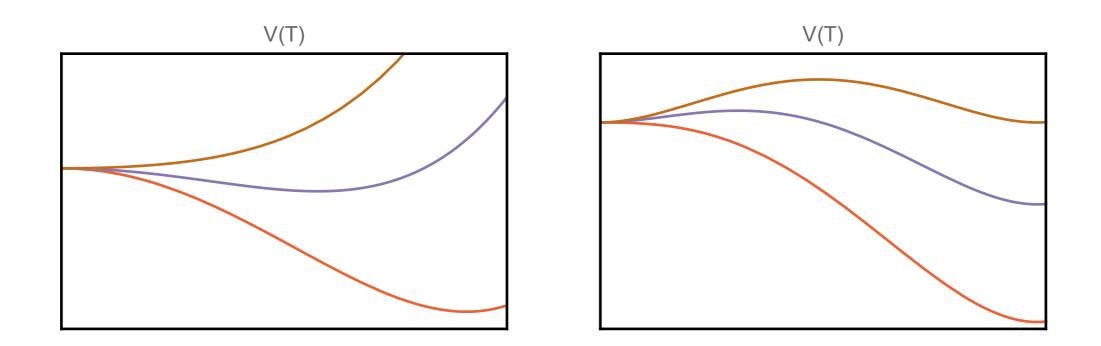




5

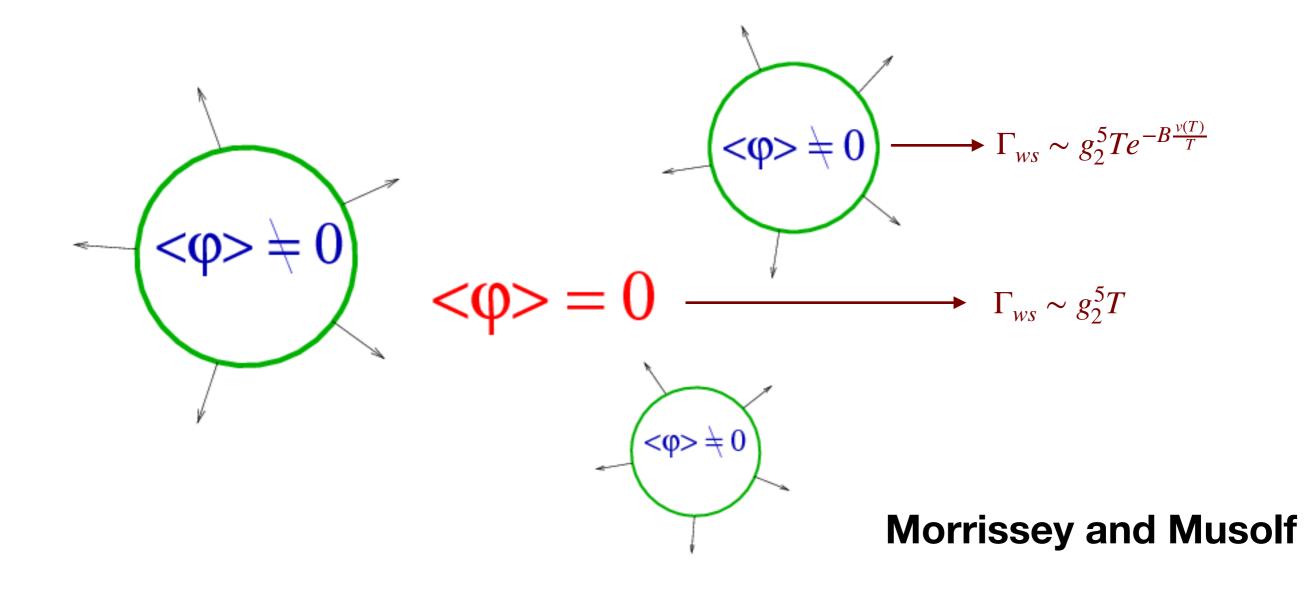
Want more details? read my book!

Electroweak symmetry at high temperature

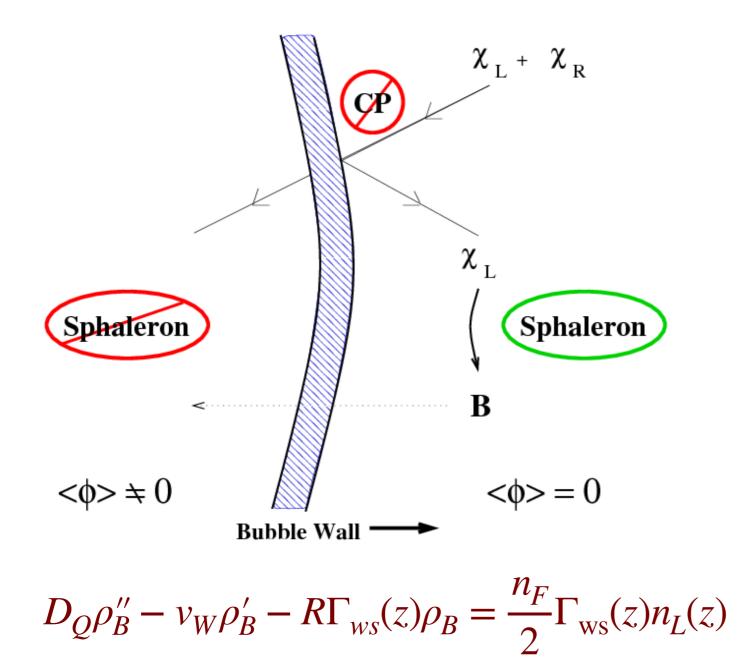


7

Electroweak symmetry at high temperature



Electroweak symmetry at high temperature



Morrissey and Musolf

8

Experimental signatures:

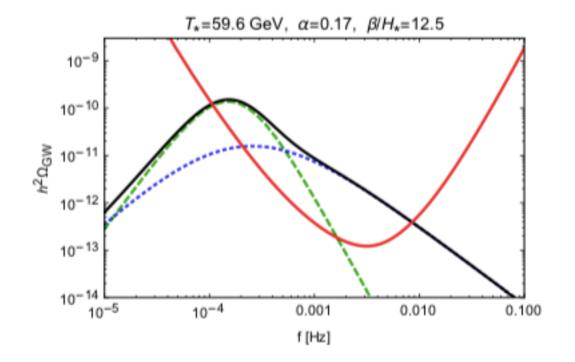
Gravitational waves Collider signatures EDMs

Experimental signatures of PT: GWs 1

1st order transitions have 3 contributions:

- Acoustic contribution
- Turbulent contribution
- Bubble collision

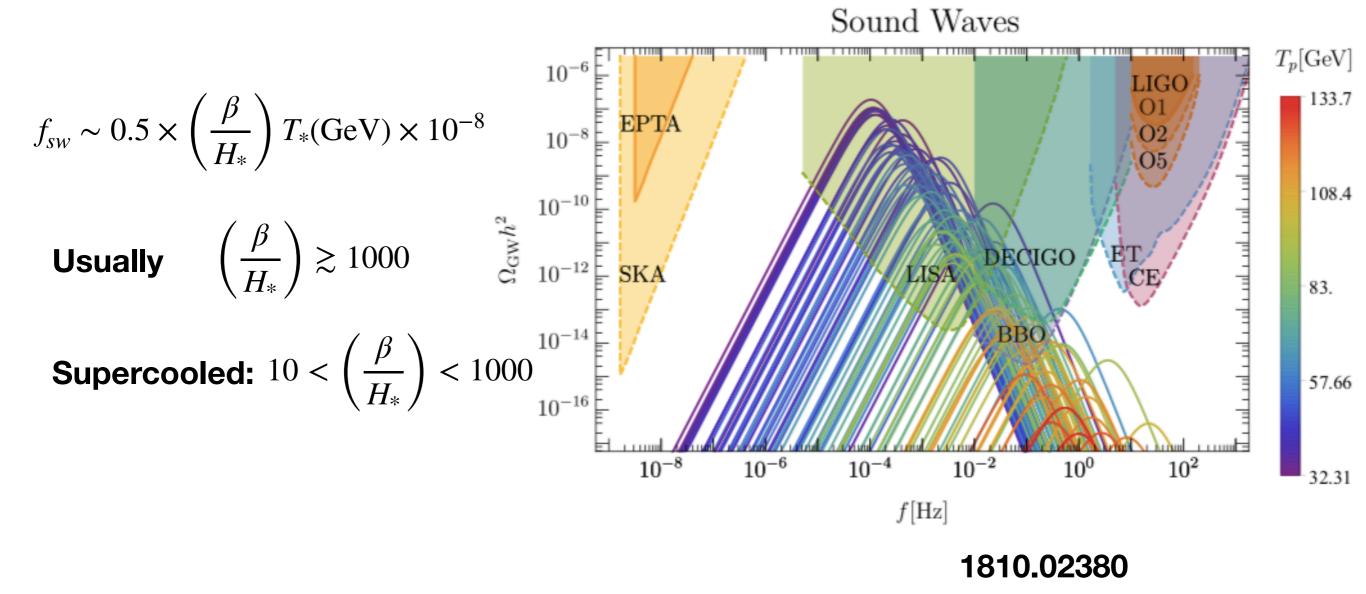
All obey a broken power law



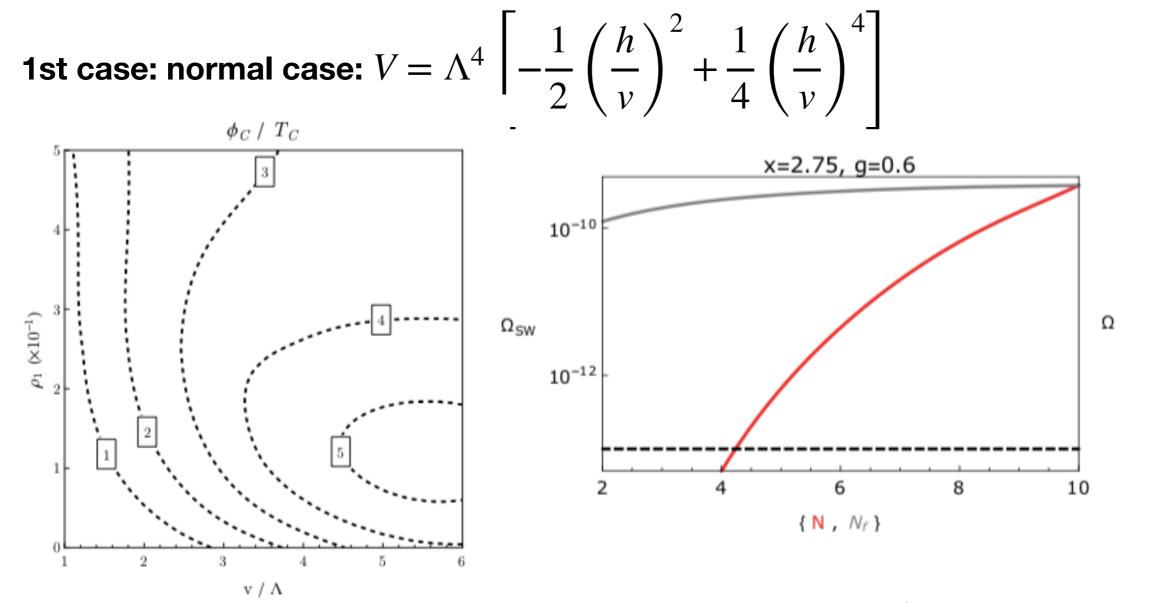
arxiv:1801.04268

11

Experimental signatures of PT: GWs 3



What drives the strength of the GW?



SM has N=3, Nf=1, v/ Λ =1.6 - not even SFO!

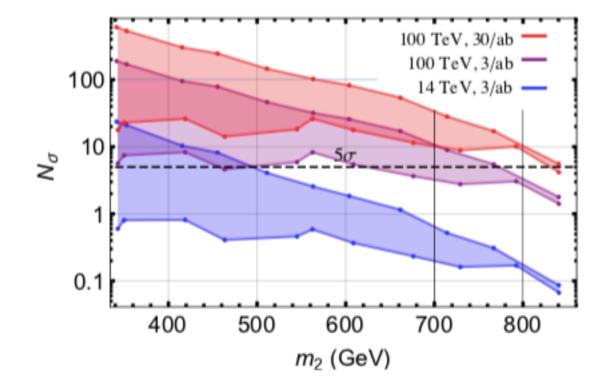
1812.02747,1806.02332

What drives the strength of the GW?

Supercooled case:

$$\phi = ah + bs, \quad V \sim \Lambda(T)^4 \left[\frac{(3 - 4\alpha(T))}{2} \left(\frac{\phi}{v(T)} \right)^2 - \left(\frac{\phi}{v(T)} \right)^3 + \alpha(T) \left(\frac{\phi}{v(T)} \right)^4 \right]$$
$$V = \Lambda(T)^4 \left[(2 - 3\alpha(T)) \left(\frac{h}{v(T)} \right)^2 - \left(\frac{h}{v(T)} \right)^4 + \alpha(T) \left(\frac{h}{v(T)} \right)^6 \right]$$

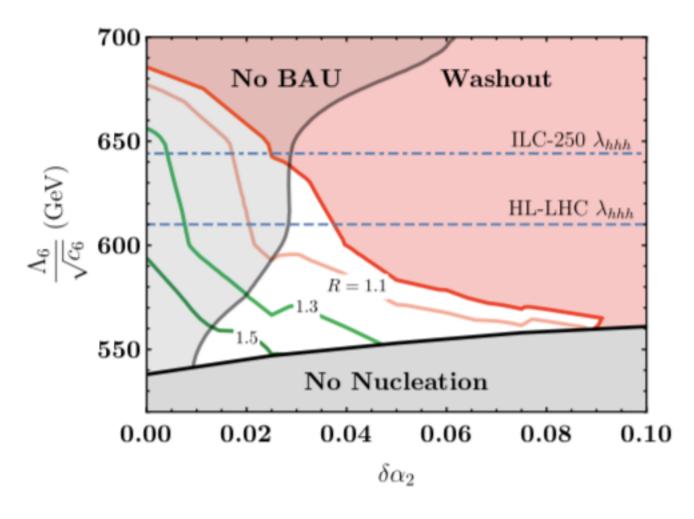
Experimental signatures of PT: Collider 1



 $b\bar{b}\gamma\gamma$ and 4τ final states

1605.06123

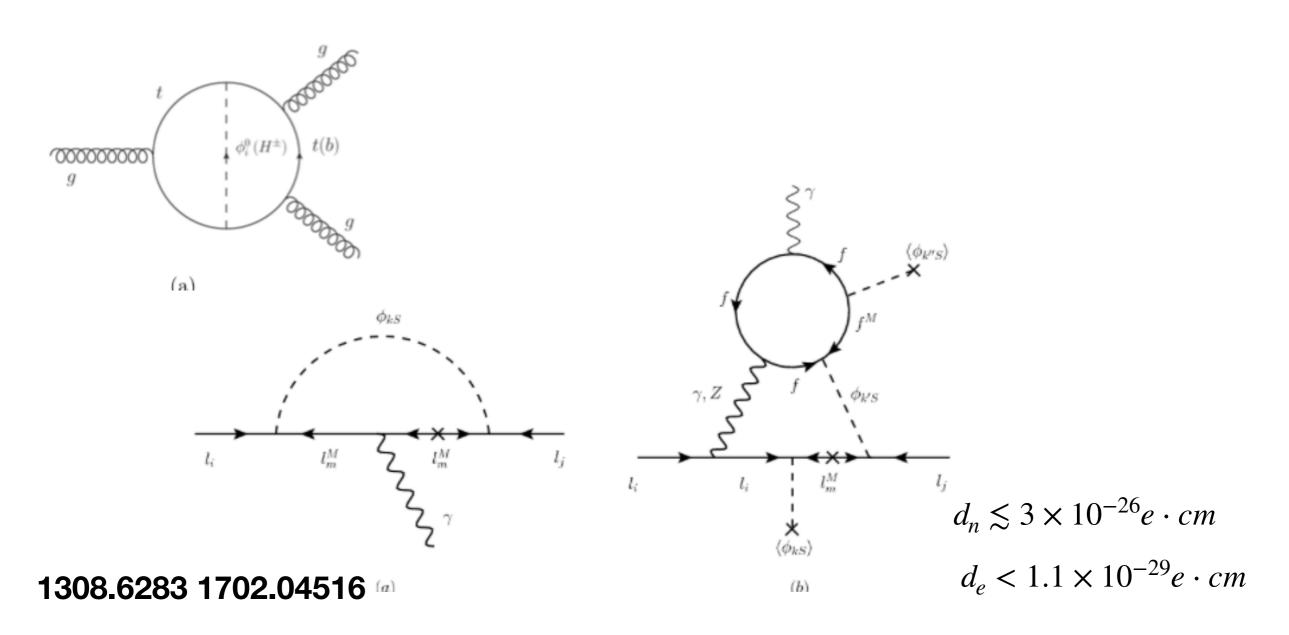
Experimental signatures of PT: Collider 2



1905.11994

16

Experimental signatures of CPV: EDMs 2



Calculation issues:

CP violating source

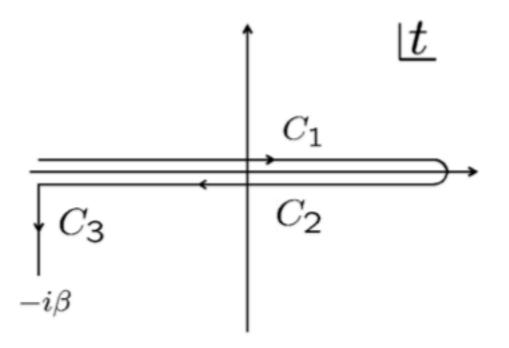
Electroweak phase transition

Calculation of CP source

Cartoon of thermal field theory

Four types of 2 point correlators

 G^{++}, G^{+-}, G^{-+} and G^{--}



Calculation of CP source

Cartoon of calculation

 $G = G_0 + \Delta G$

 $\mathcal{E}(x)G_{ij}^{+-}(x,y) - G_{ij}^{+-}(x,y)\mathcal{E}(y) = \mathcal{E}(x)\Delta G_{ij}^{+-}(x,y) - \Delta G_{ij}^{+-}(x,y)\mathcal{E}(y)$

"Semi Classical force" <-> LHS, Self-energies <-> RHS

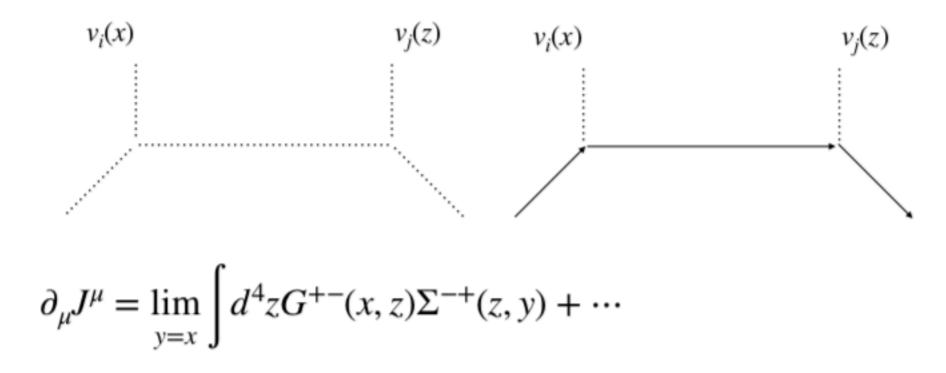
VEV insertion approach:

- 1. ignore off diagonals
- 2. Expand around x=y
- 3. Use electroweak symmetric basis

Space time varying basis!

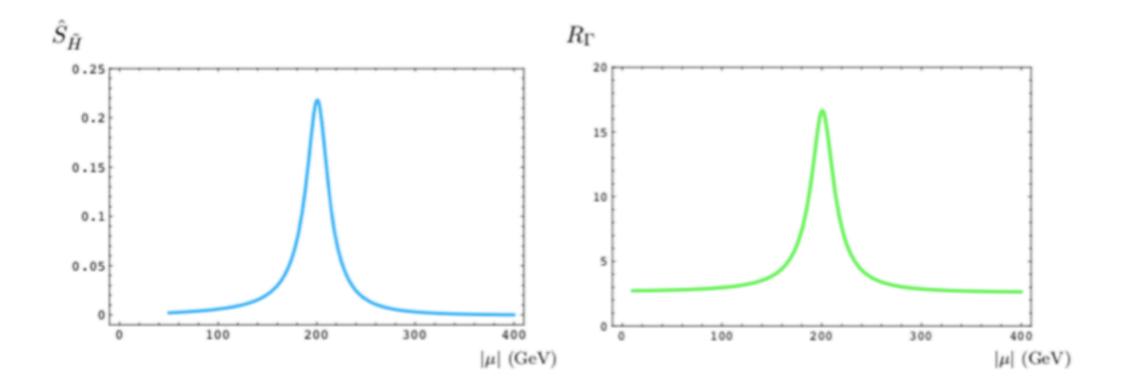
Calculation of CP source

Vev insertion sources



Calculation of CP source

Resonant behaviour of sources



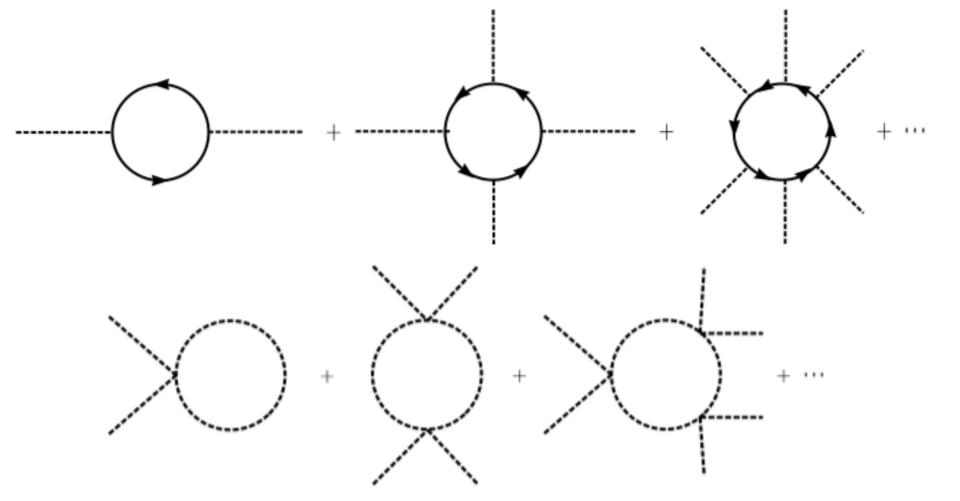
Lee, hep-ph/0412354

Calculation of CP source

Proper handling of CP sources still an open problem!

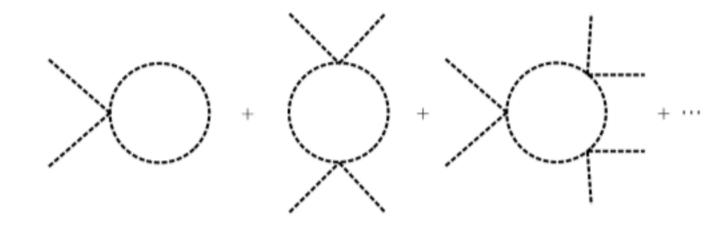
Calculation of phase transition (what to calculate)^{2³}

 $V(\phi) \to V_0(\phi) + V_1(\phi, T)$



Calculation of phase transition (what to calculate)²⁴

 $V(\phi) \rightarrow V_0(\phi) + V_1(\phi, T)$



$$\frac{dV_1}{dm^2} = \int \frac{d^4p}{(2\pi)^4} i\Delta^{++}(p,T)$$

 $\Delta^{++}(p,T)=\Delta^{++}(p,0)+\delta^{++}(p,T)$

 $V_1(\phi, T) = V_{CW}(\phi) + V_T(\phi, T)$

Calculation of phase transition (what to calculate)²⁵

$$\frac{\partial^2 \phi}{dr^2} + \frac{2}{r} \frac{d\phi}{dr} = \frac{dV}{d\phi}$$

$$S_E = 4\pi \int r^2 dr \left[\frac{1}{2} (\phi')^2 + V(\phi) \right]$$

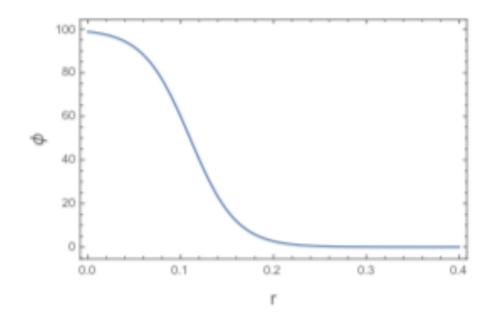
$$S_E / T_N \sim 140$$

$$\alpha = \frac{\Delta V - \frac{1}{4} T \Delta (dV/dT)}{\rho_R} \Big|_{T_N}$$

$$\beta = T \frac{d(S_E / T)}{\rho_R}$$

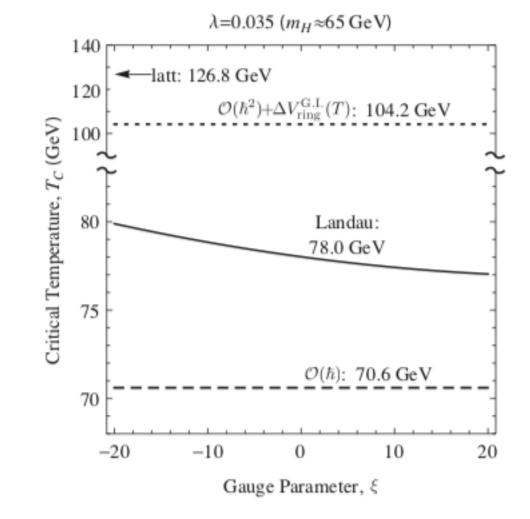
dT

 H_*



Calculation of phase transition Issue 1:

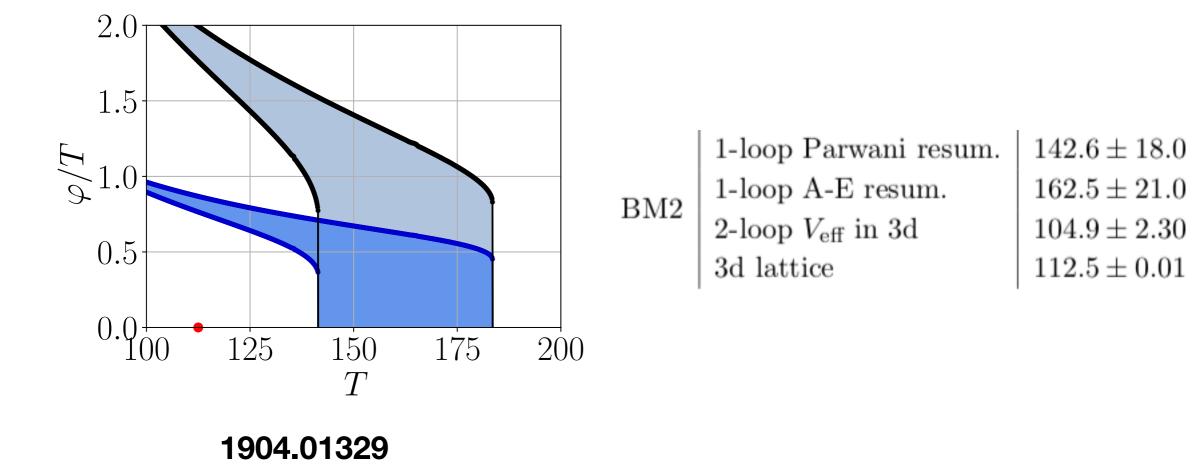
Gauge dependence



1101.4665

Calculation of phase transition Issue 2:

Renormalization dependence

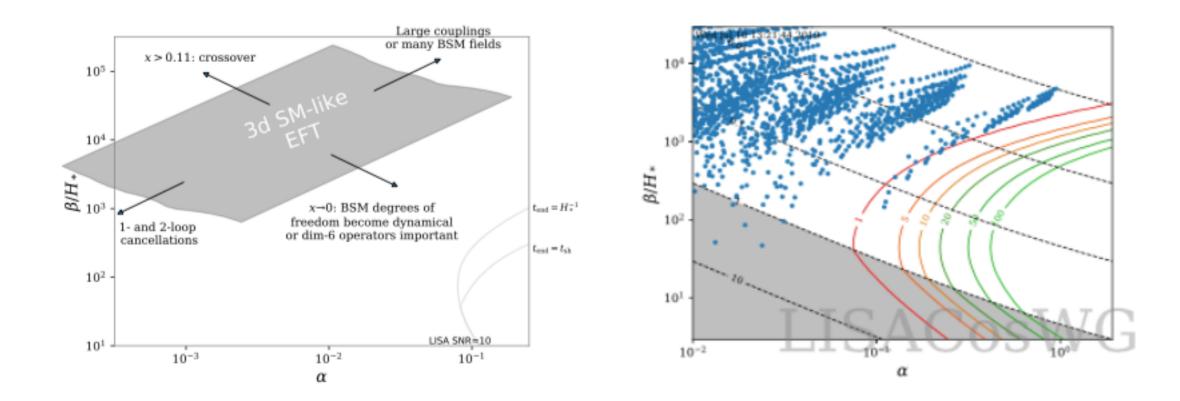


Calculation of phase transition Issue 2:

Renormalization dependence

- 3d (dimensional reduction)
 - Integrate out heavy matsubara modes
 - Manifestly gauge invariant
 - Super-renomalizable only need to go to 2 loop level
 - 2 loop calculations are easy in 3d
 - Includes all order resummations by construction
 - Lattice calculations more tractable
 - Only works in when HT expansion is valid

Calculation of phase transition Issue 2:



1903.11604

Croon, Gould, Tenkanen and White, upcoming

Summary of calculation techniques in EWBG:

Great deal of theoretical uncertainty on both the calculation of CP violating sources and phase transitions

Very much a work in progress!

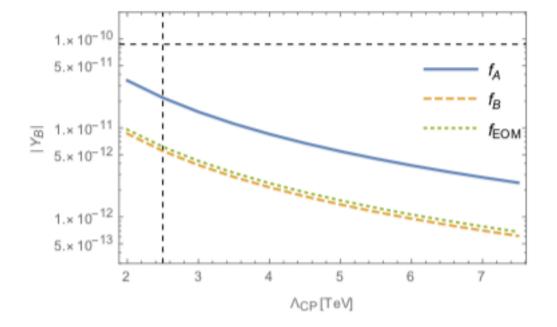
Status of some favourite EWBG models

Status of EFT baryogenesis

First an odd feature: Consider the source generated by 2 dim-6 operators related by eom

$$S_{\mathcal{O}_{DD}}^{\mathcal{O}\mathcal{P}} \sim \frac{1}{\Lambda^2} \left[v(x)\partial_t \left(\partial_\mu \partial^\mu v(x) \right) - \partial_t v(x) \left(\partial_\mu \partial^\mu v(x) \right) \right],$$
$$S_{\mathcal{O}_{\partial V/\partial H}}^{\mathcal{O}\mathcal{P}} \sim \frac{1}{\Lambda^2} \left[v(x)\partial_t \left(\left. \frac{\partial V_{\rm SM}}{\partial H} \right|_{v(x)} \right) - \partial_t v(x) \left(\left. \frac{\partial V_{\rm SM}}{\partial H} \right|_{v(x)} \right) \right] + \mathcal{O}\left(\frac{1}{\Lambda^4} \right).$$

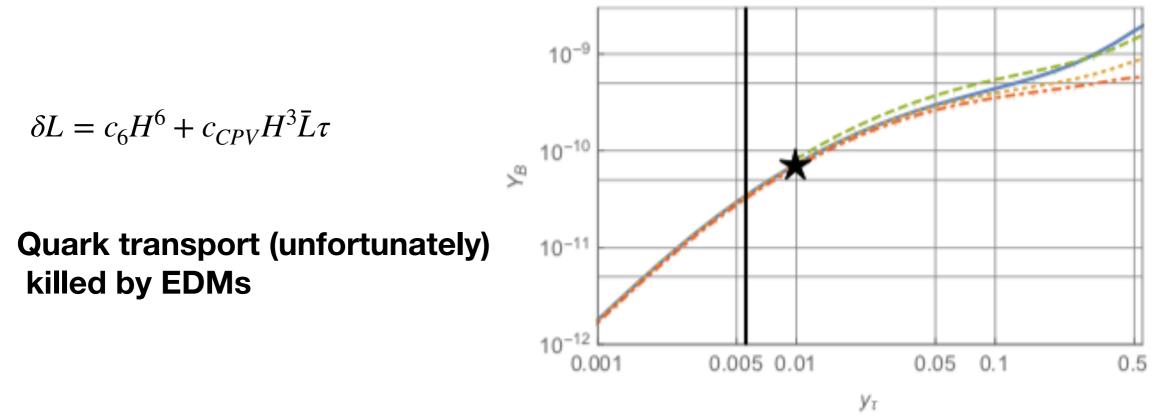
Degeneracy is broken unless dim 8 piece is included!



1710.04061

Status of EFT baryogenesis

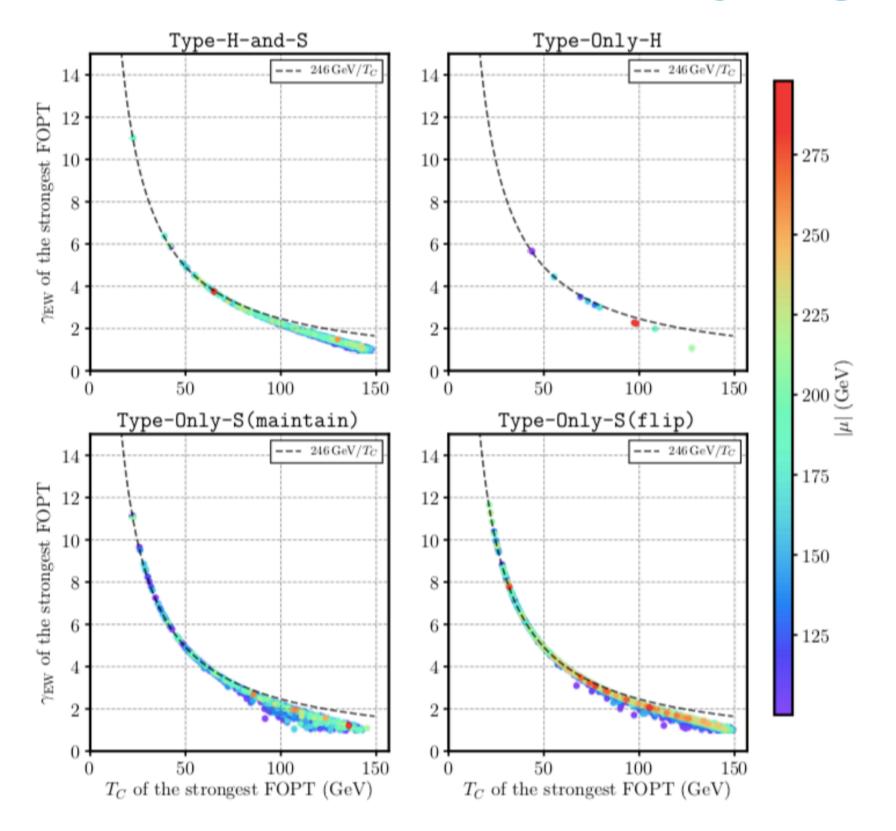
Appears to be viable with a tau source



1811.11104

Status of NMSSM baryogenesis

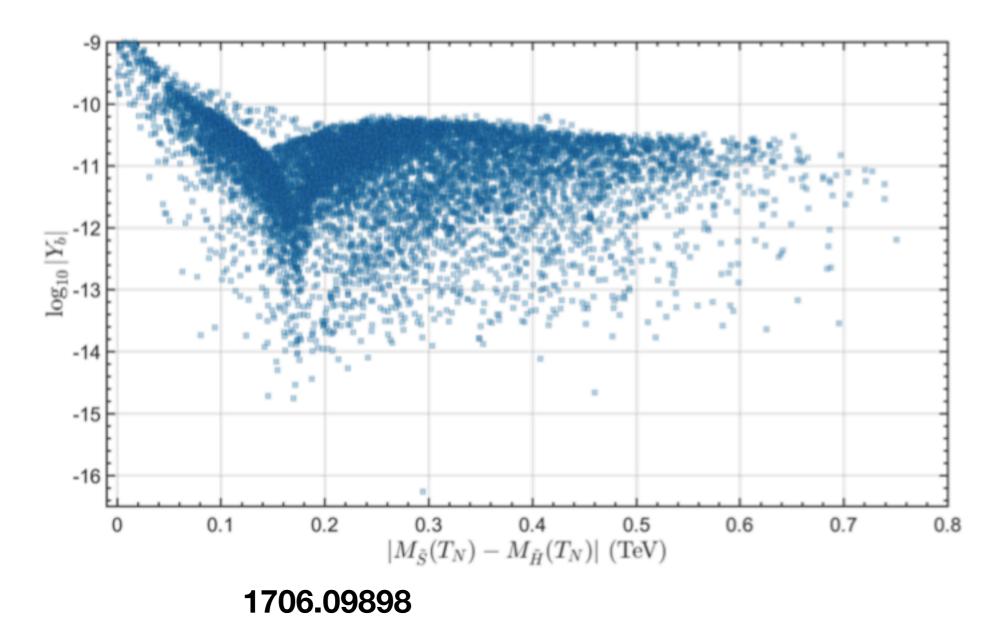
- Z3 invariant potential for simplicity
- Singlet (rather than stop) can catalyze 1st order transition
- Singlino-Higgsino interactions with Higgs can be a weakly constrained CPV source



35

1908.11847

Type H and S is optimal Type only S means soft terms can be Large enough for Boltzmann suppression to matter Even still type only S is viable



How high can the CPV be?

A simple example

$$-\mathcal{L} \ni m_{A}^{2} |A|^{2} + m_{B}^{2} |B|^{2} + \left[\mu + \frac{\xi |H|^{2}}{\Lambda} \right] A^{\dagger} HB + (\kappa |A|^{2} + \kappa_{B} |B|^{2}) |H|^{2} + h \cdot c. + S \cdot B \cdot Ts$$

Quantum numbers

A = (1,2,1/2) B = (1,1,0)

How high can the CPV be?

Source estimation: tree level

 $-\mathcal{L} \ni A^\dagger \left[\mu_1 H_1 + \mu_2 H_2 \right] B$

 $\sim \mathrm{Im}[\mu_1 \mu_2] \beta'(x) v(x)^2 I(m_i, \Gamma_i)$

$$10^2 \lesssim {\rm Max} \left[\frac{Y_B}{Y_B^{\rm obs}} \right] \lesssim 10^3$$

Source estimation: loop

$$-\mathscr{L} \ni \left[\mu + \frac{\xi |H|^2}{\Lambda} \right] A^{\dagger} HB$$

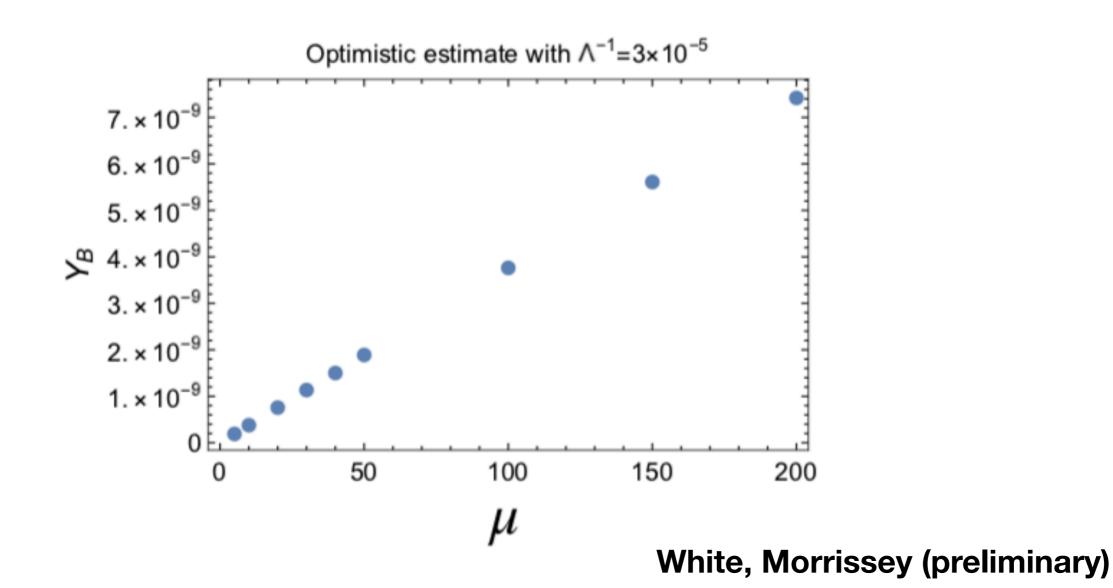
$$\sim \operatorname{Im}[\mu\xi] \frac{v(x)^2}{\Lambda} v(x) v'(x) I(m_i, \Gamma_i)$$

For $\mu \sim 200 \text{ GeV}$ Max[Λ] ~ $\mathcal{O}[1 - 10] \text{ PeV}$

29

38

How high can the CPV be?



39

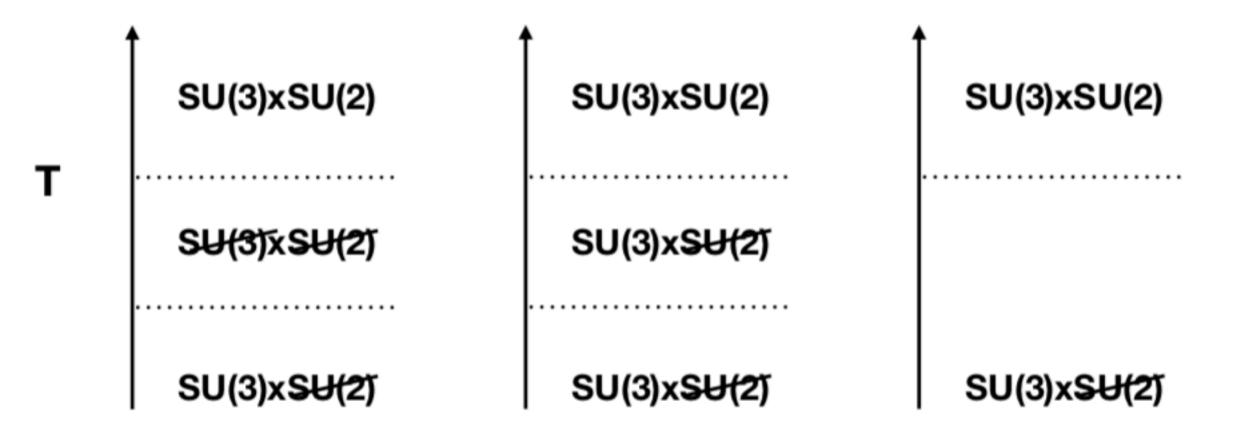
Summary of current status

- Many favourite models are still alive and very well
- EDMs usually but not necessarily the most stringent constraint
- Always have the caveat of theoretical uncertainties

Extensions to the minimal EWBG framework

- Extensions to ways of realizing CPV
- Modify B violation
- Modify departure from equilibrium

Extension 1 departure from equilibrium



Extension 1 colour breaking baryogenesis

Lagrangian and model for COB

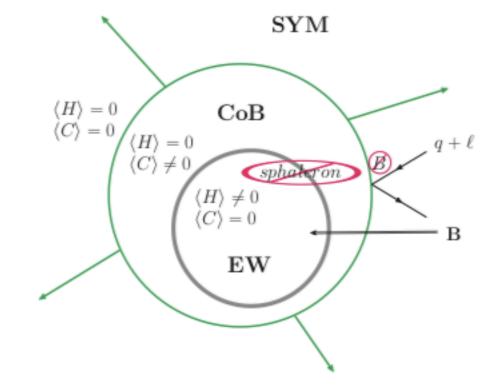
$$L = L_{\rm SM} + \lambda_i C_i \bar{b}_R L + \Delta V$$

 $C_i = (3, 2, 1/6)$

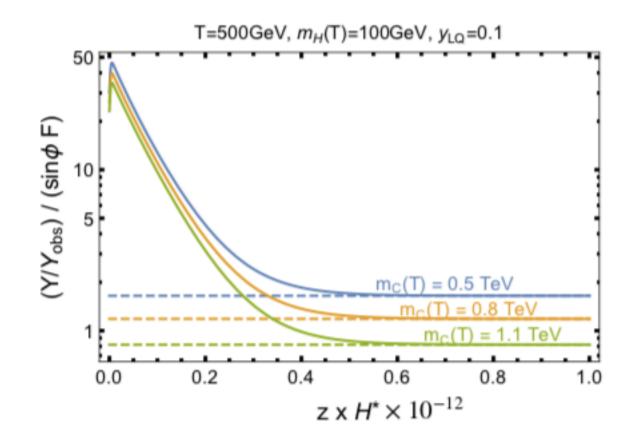
 $\langle B \rangle \rightarrow$ Spontaneous and Sphalerons

 $\mathcal{CP} \leftrightarrow \lambda_i$ Restrict to 3rd generation

Departure from Equilibrium – Colour breaking phase. Can happen at multi TeV scale

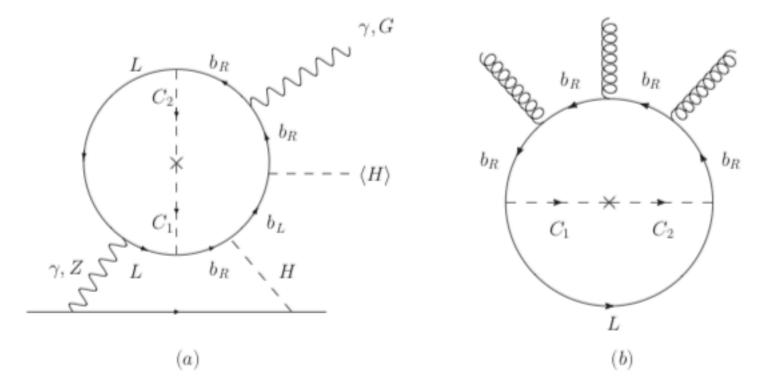


Extension 1 colour breaking baryogenesis



Extension 1 colour breaking baryogenesis

Experimental constraints



Experimental signal: flavour anomalies, Gravitational waves, neutron EDM, leptoquark production at upgraded LHC?

Extension 2 Changing sphalerons

Case 1: Modifying couplings through a PT $\mathscr{L} \supset -\frac{1}{4} \left(\frac{1}{g_3^2} + a_{\phi} \frac{\phi}{\Lambda_{\phi}} \right) G^{a \ \mu\nu} G^a_{\mu\nu} = -\frac{1}{4} \frac{1}{g_3^2_{\text{eff}}} G^{a \ \mu\nu} G^a_{\mu\nu} \quad g_{3_{\text{eff}}} = g_3 \left(\frac{\Lambda_{\phi}}{a_{\phi} g_3^2 \phi + \Lambda_{\phi}} \right)^{1/2}$ UV completion: 1. triangle diagram2. Dilaton-like field $T_{\text{st}} = \frac{1}{2} \frac{1}{g_3^2} G^{a \ \mu\nu} G^a_{\mu\nu} = \frac{1}{2} \frac{1}{g_3^2_{\text{eff}}} G^{a \ \mu\nu} G^a_{\mu\nu} = \frac{1}{2} \frac{1}{g_3^2} \frac{1}{g_3^2} G^{a \ \mu\nu} G^a_{\mu\nu} = \frac{1}{2} \frac{1}{g_3^2} \frac{1}{g_3^2} G^{a \ \mu\nu} G^a_{\mu\nu} = \frac{1}{2} \frac{1}{g_3^2} \frac{1}{g_3^2}$

 $\Gamma_{ws} \sim g_2^5$

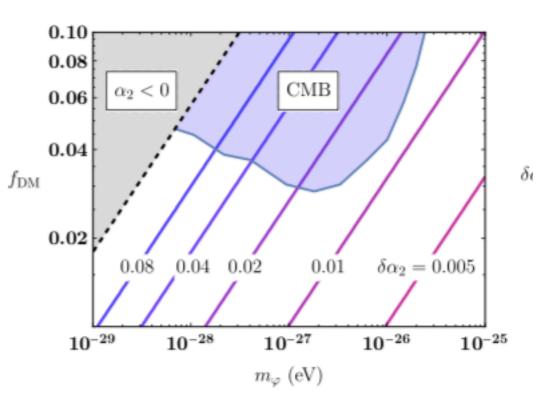
 $\Gamma_{ss} \sim g_2^5$

Extension 2 Changing sphalerons

Case 2: oscillating dilaton $\alpha_w \rightarrow \alpha_w(T)$

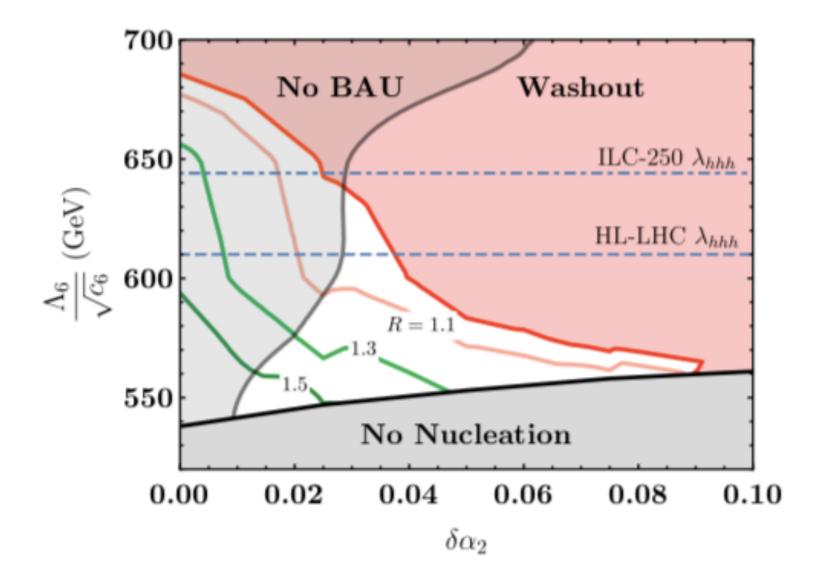
$$\begin{split} \delta L &= -\frac{1}{4} \left(1 - \frac{d_{gY}}{g_Y^2} \kappa \phi \right) A^{\mu\nu} A_{\mu\nu} - \frac{1}{4} \left(1 - \frac{d_{g2}}{g_2^2} \kappa \phi \right) W^{a\mu\nu} W^a_{\mu\nu} \\ \phi &\approx \frac{\sqrt{2f_{DM}\rho_{DM}}}{m_{\phi}} \cos(m_{\phi}(t - v \cdot x + \cdots) \propto \left(\frac{T}{T_{CMB}} \right)^{3/2} \end{split}$$

Planck Suppressed coupling Can be large in the early universe

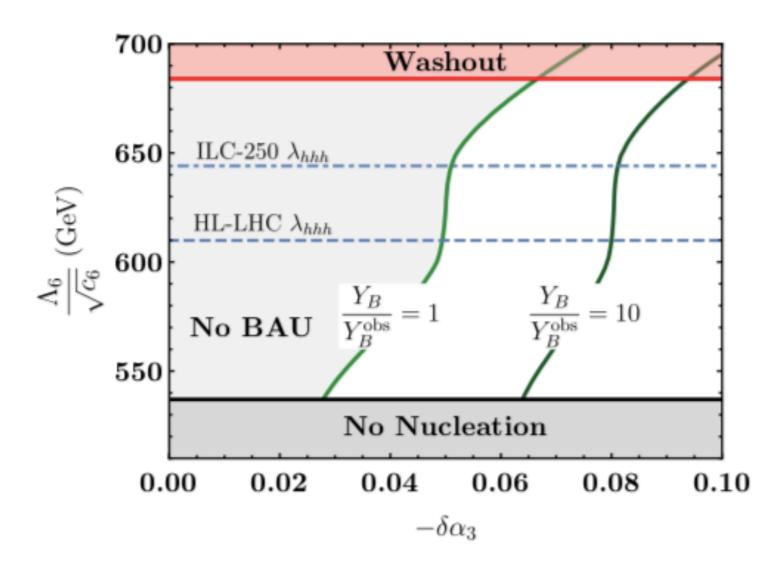


Extension 2 changing sphalerons

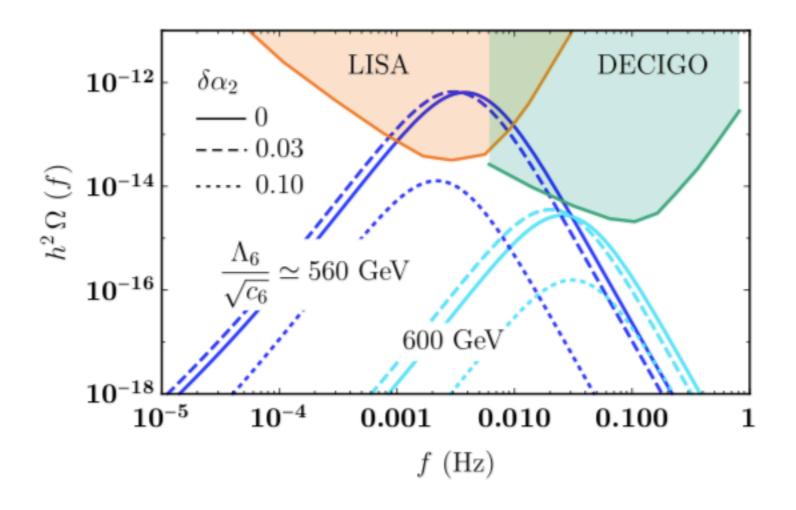
Apply paradigm to quark SMEFT



Extension 2 changing sphalerons



Extension 2 changing sphalerons



Extension 3 modified CPV

What if CPV was different in the early Universe compared to today?

Example: CKM can be less suppressed in early Universe if you have dynamical Yukawas

$$L_Y = \left(\frac{S}{M}\right)^{n_i + n_j} \bar{Q}_L^i H q_R^j$$

Extension 3 modified CPV

Original idea: Jarlskog invariant $\sim \left(\frac{S}{M}\right)^{20}$

If S was close to M in the early Universe there would be no suppression

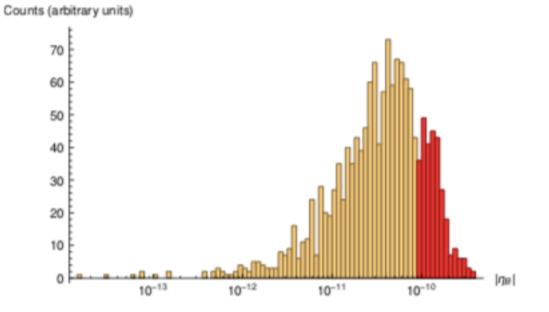
Recent idea: if Yukawas vary during the phase transition, CPV is no longer Loop suppressed

Extension 3 modified CPV

$$y(y_0, y_1, \phi, n) = (y_0 - y_1) \left[1 - \left(\frac{\phi}{v}\right)^n \right] + y_1$$

Varying the up charm yukawas Can produce enough BAU

1706.08534



Extension 3 modified CPV

Lots to be done on this front:

Both the cases where the Yukawas vary or are simply larger need testing with other methods of calculating the CP asymmetry

Conclusions

Electroweak baryogenesis is a testable and minimal explanation for why we are here The phenomenology is rich - from colliders to GWs to EDM searches Much theoretical uncertainty remains Many of the most popular models are still works in progress Extensions to the basic paradigm are phenomenologically rich and underexplored