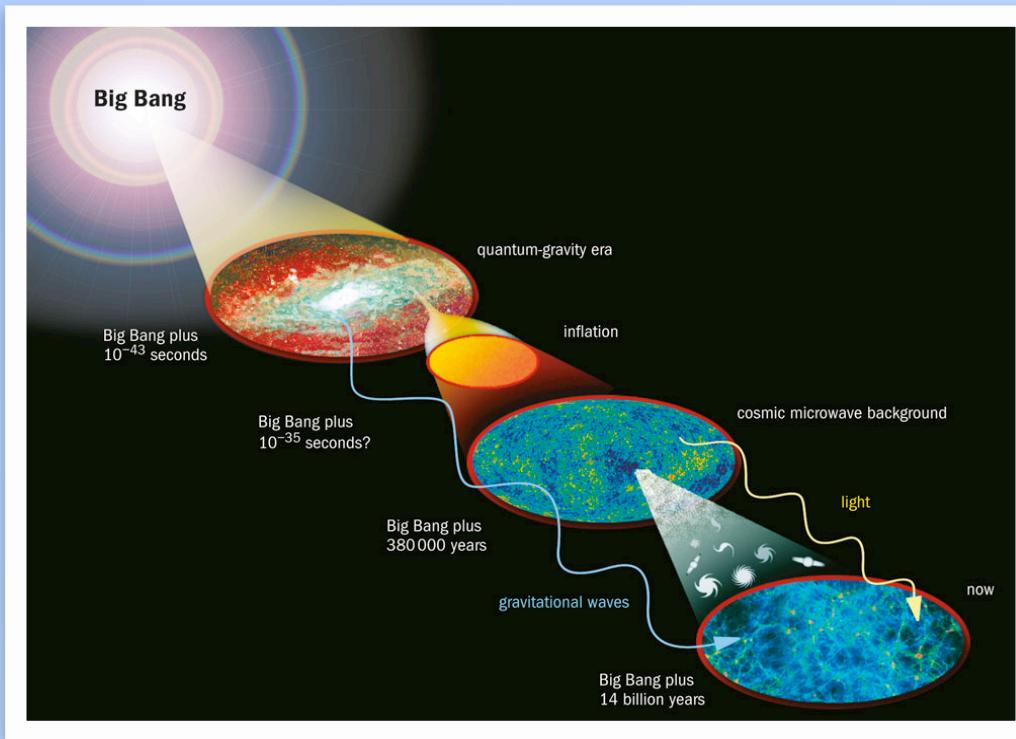


# Nikhef Topical Lectures

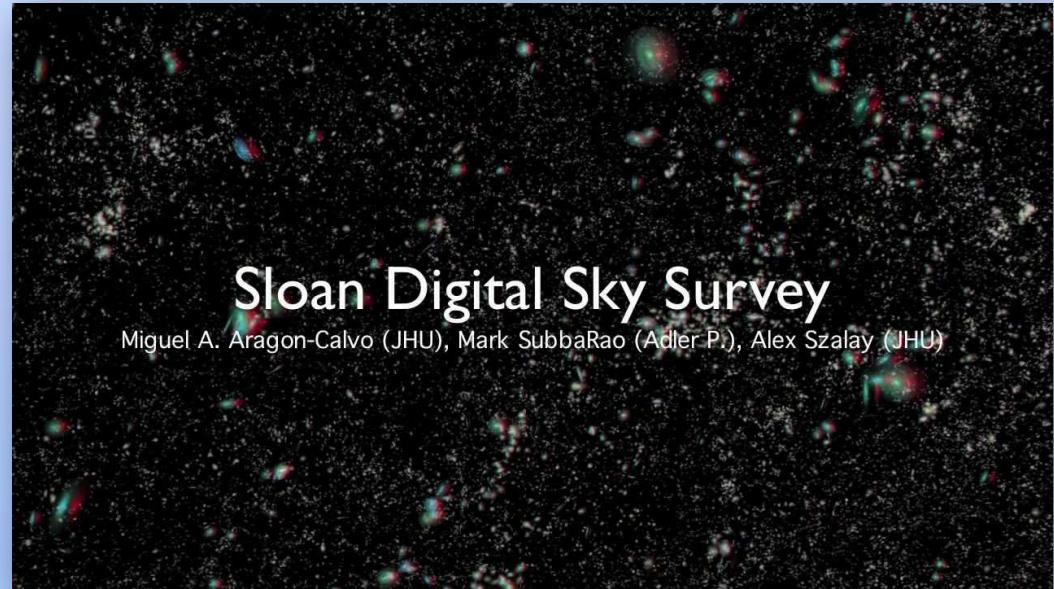
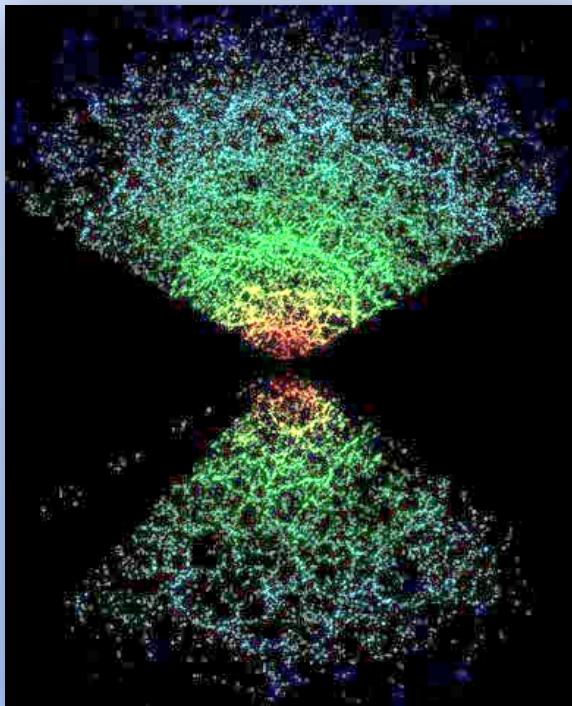
# Inflationary Cosmology



Jan Pieter van der Schaar  
Institute of Physics, Amsterdam

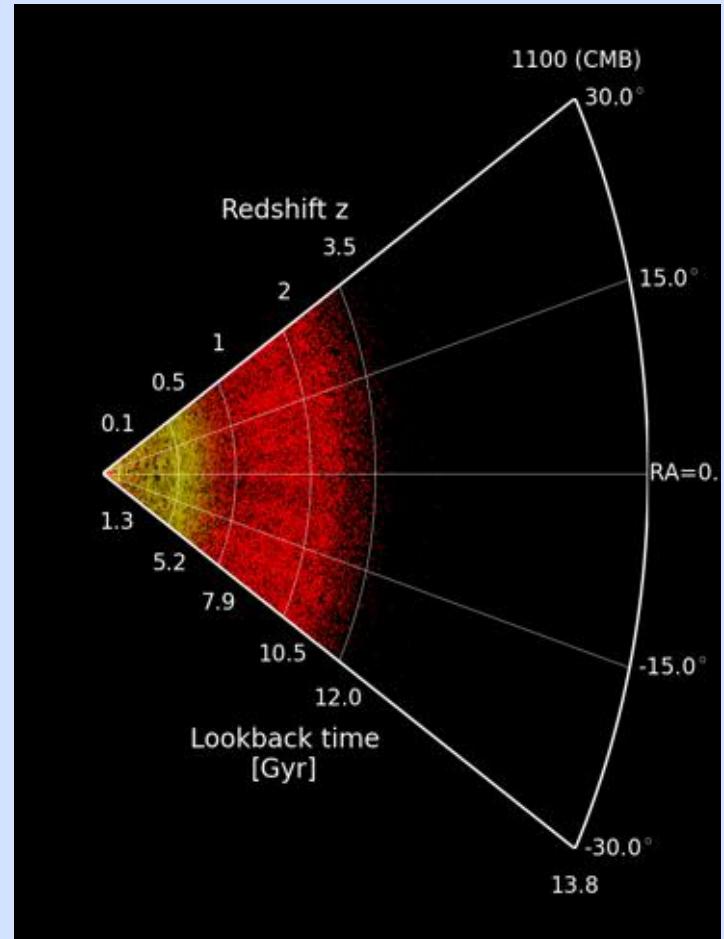
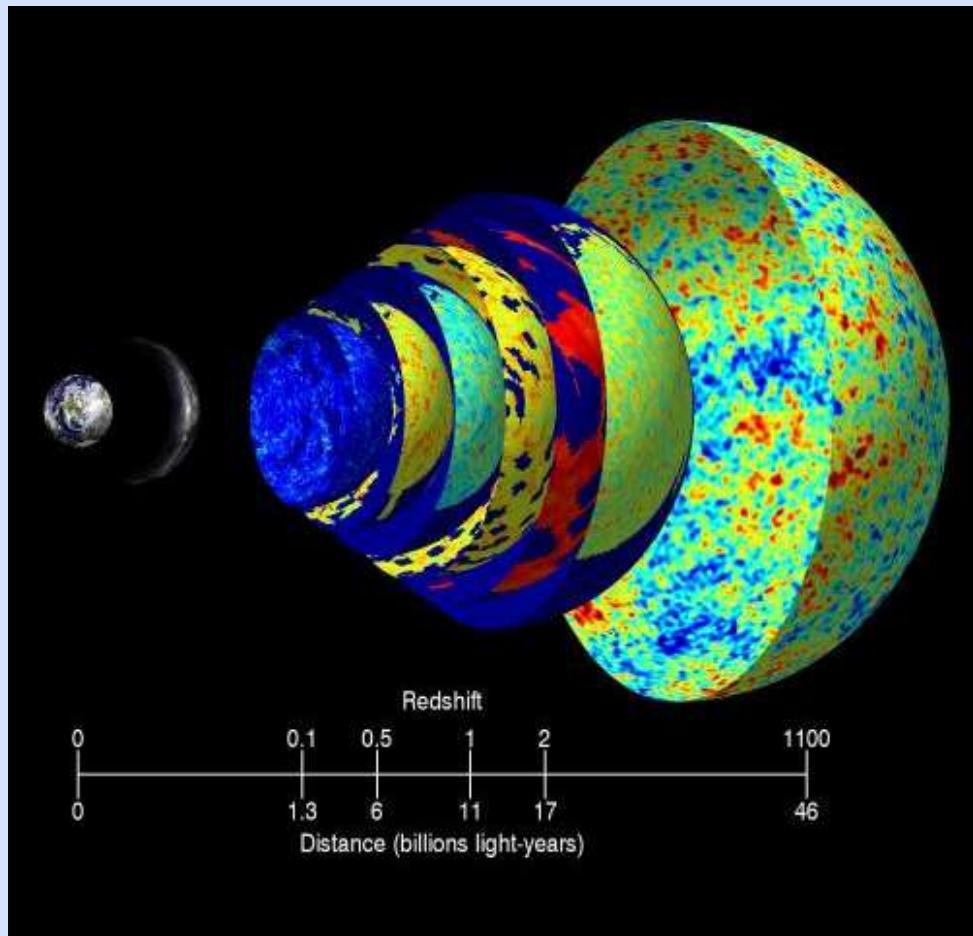
# SLOAN DIGITAL SKY SURVEY

## 3D journey

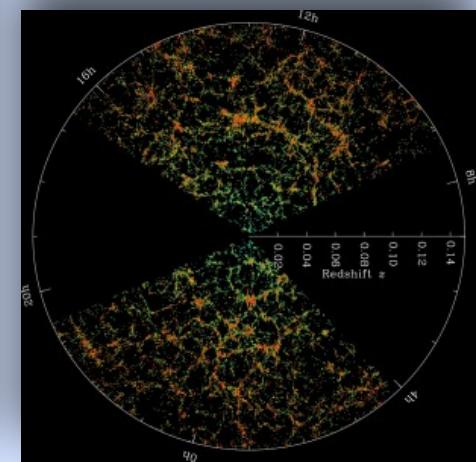
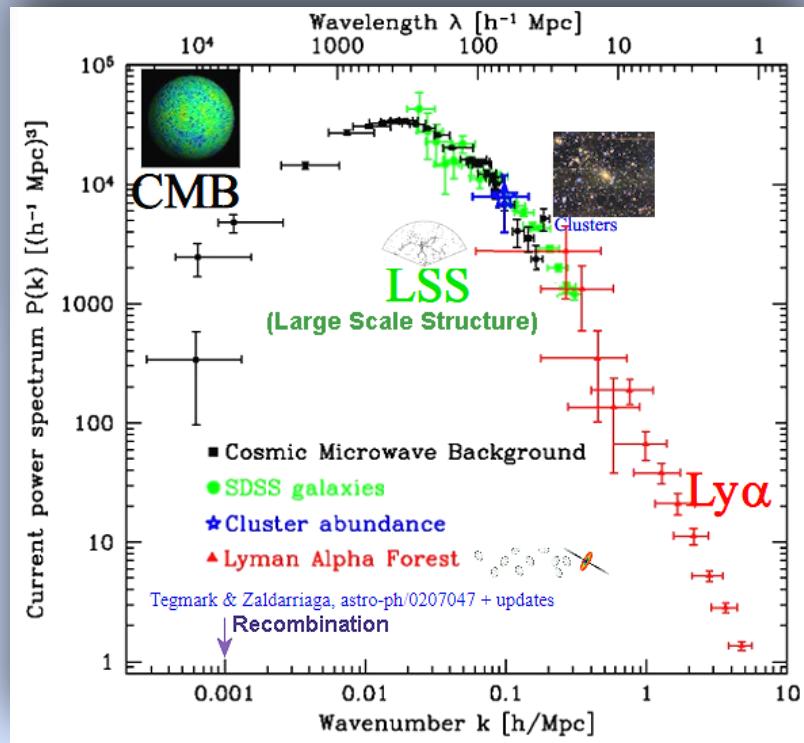
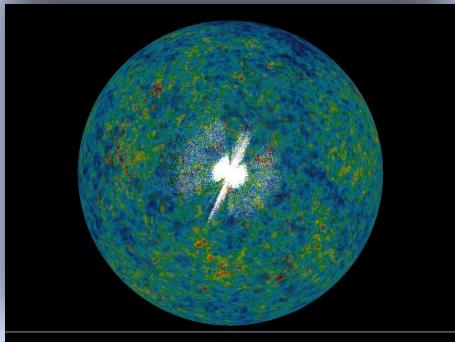


**Sloan Digital Sky Survey**  
Miguel A. Aragon-Calvo (JHU), Mark SubbaRao (Adler P.), Alex Szalay (JHU)

# Explaining correlations



# Explaining correlations

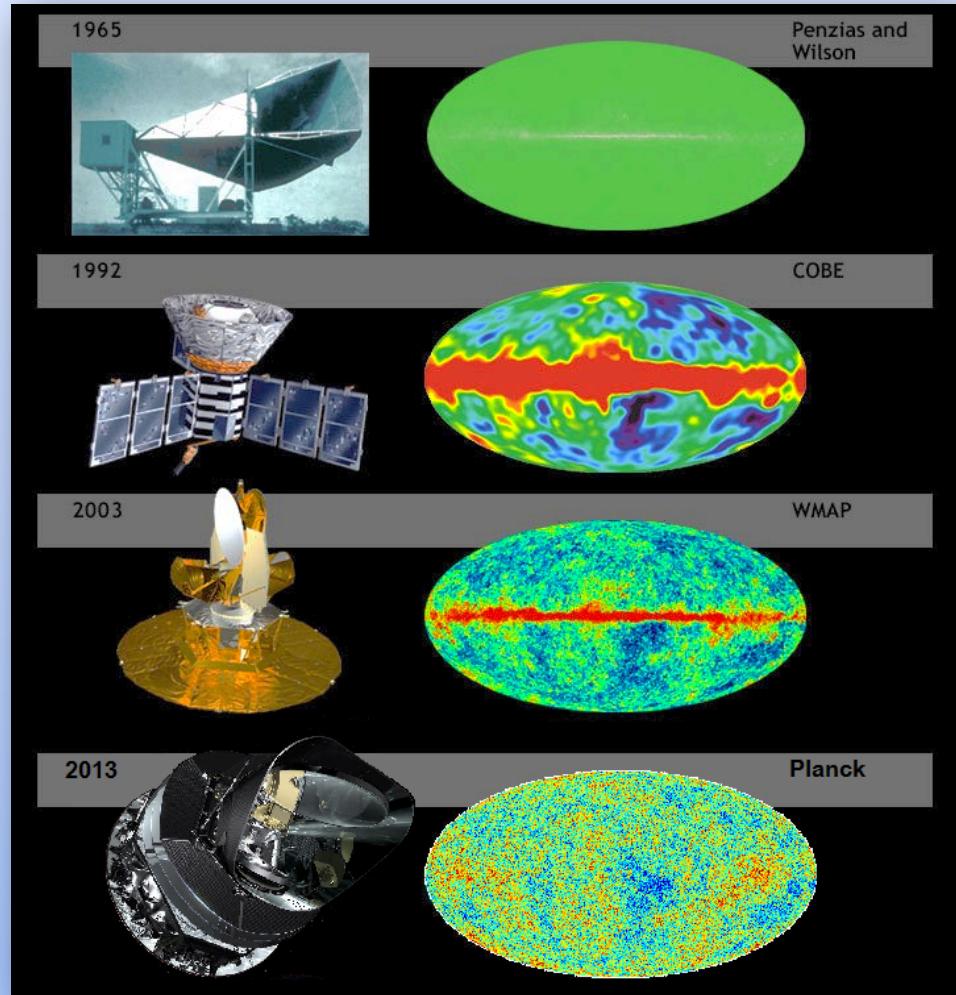
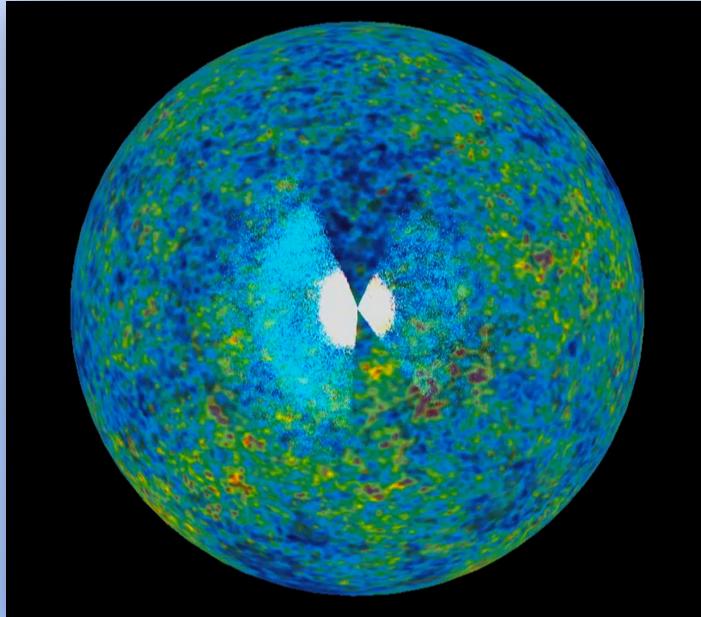
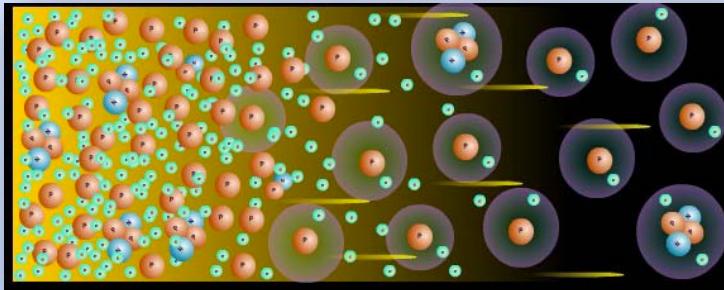


From CMB at  $z=1100$  to...

...large scale structure at  $z=0(1)$

**Not predicted by standard hot big bang**

# Big Bang afterglow



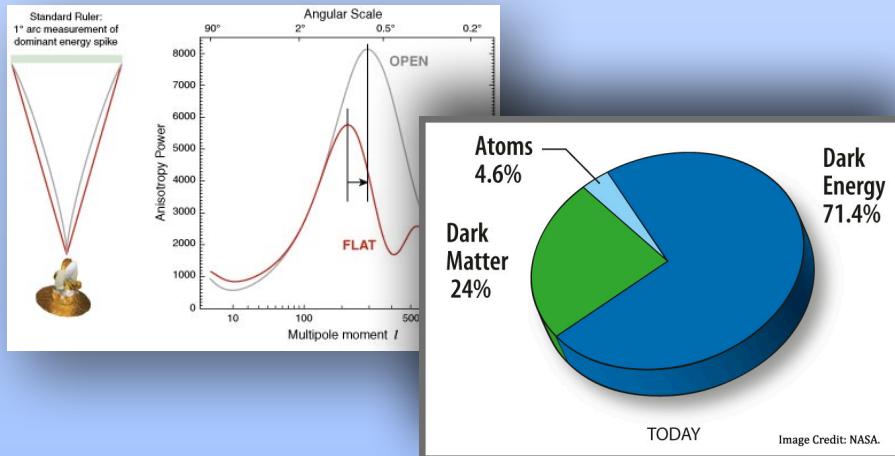
# Plan

## Today

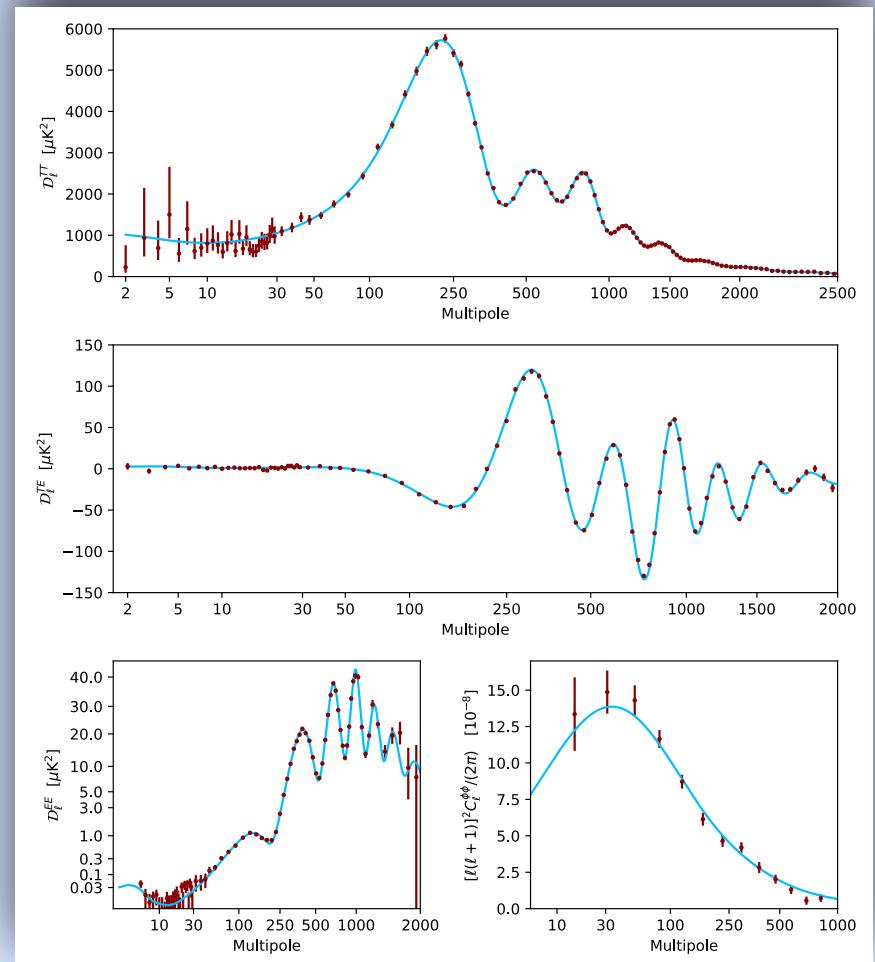
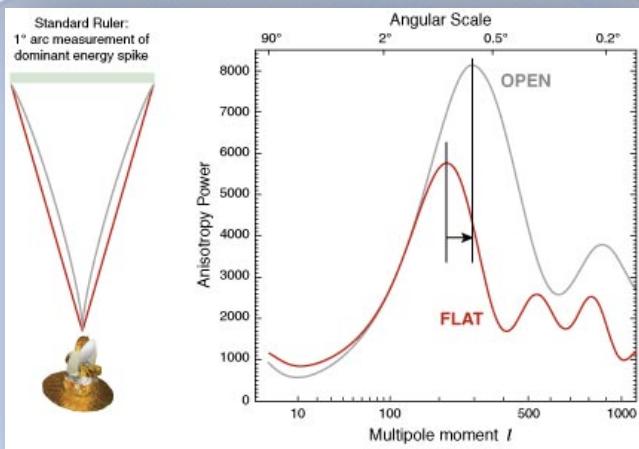
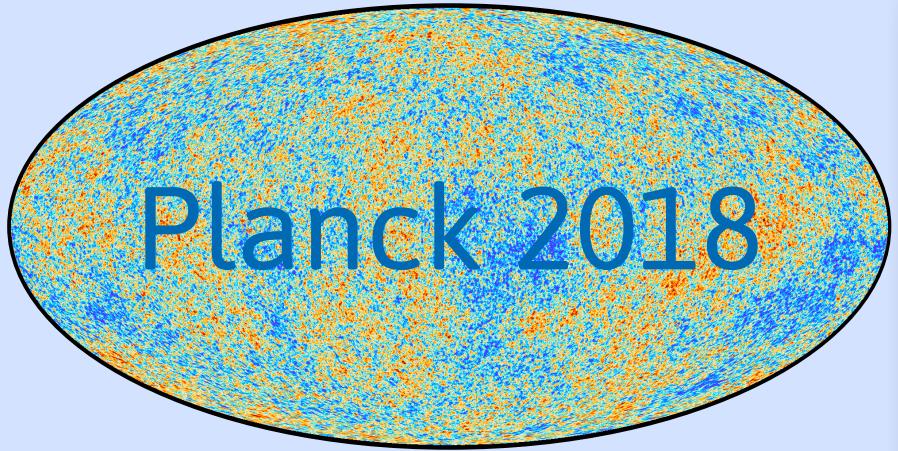
- Initial conditions of the hot Big Bang
- The paradigm of cosmological inflation

## Tomorrow

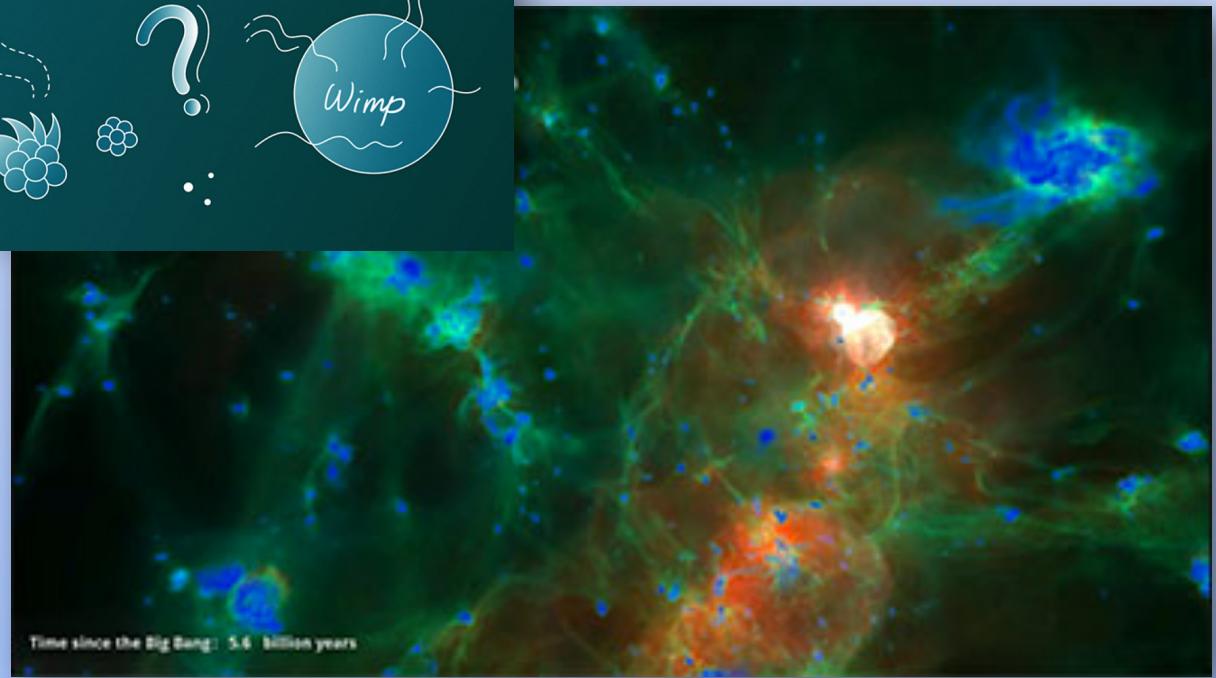
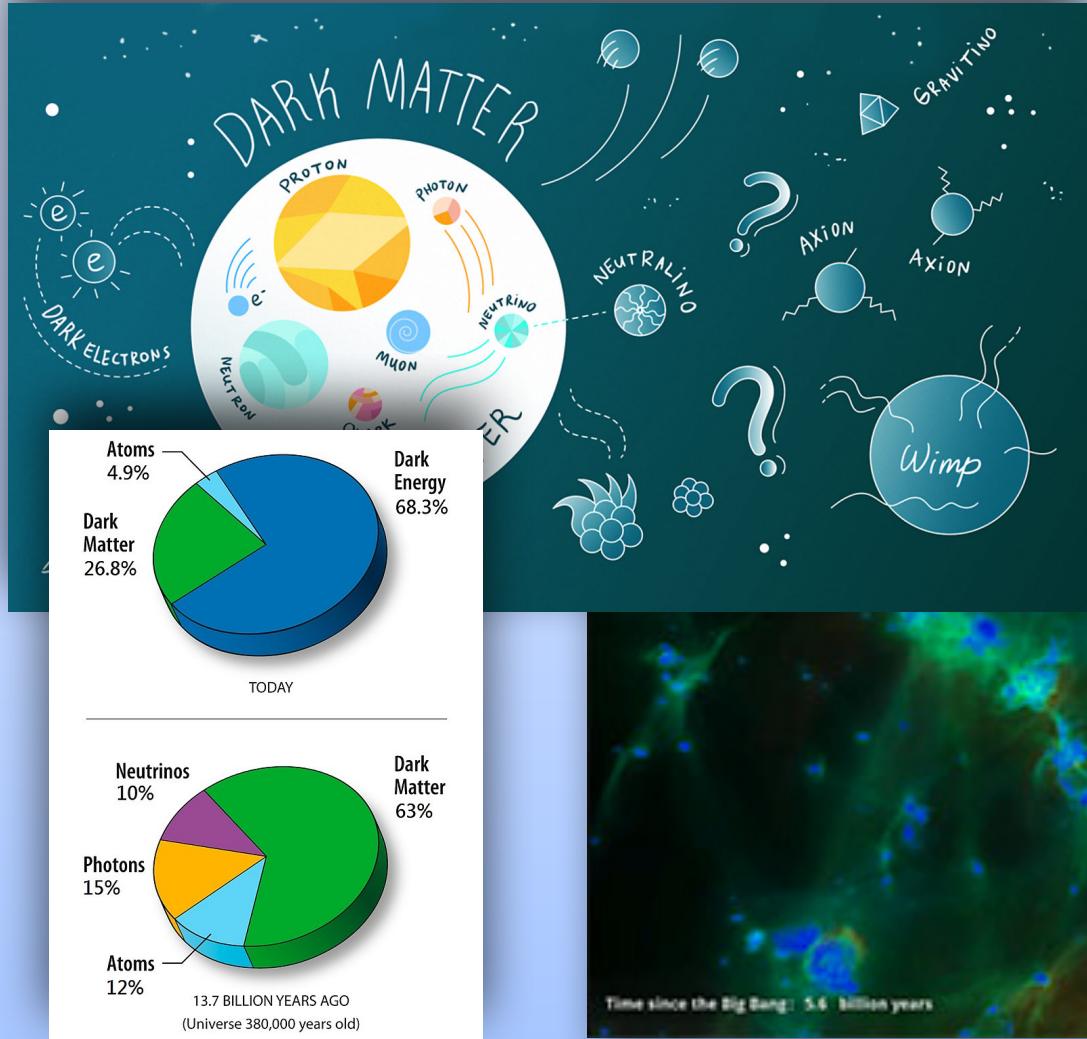
- Inflation and the origin of structure
- The cosmological accelerator: constraining inflation



# Part I: Initial conditions



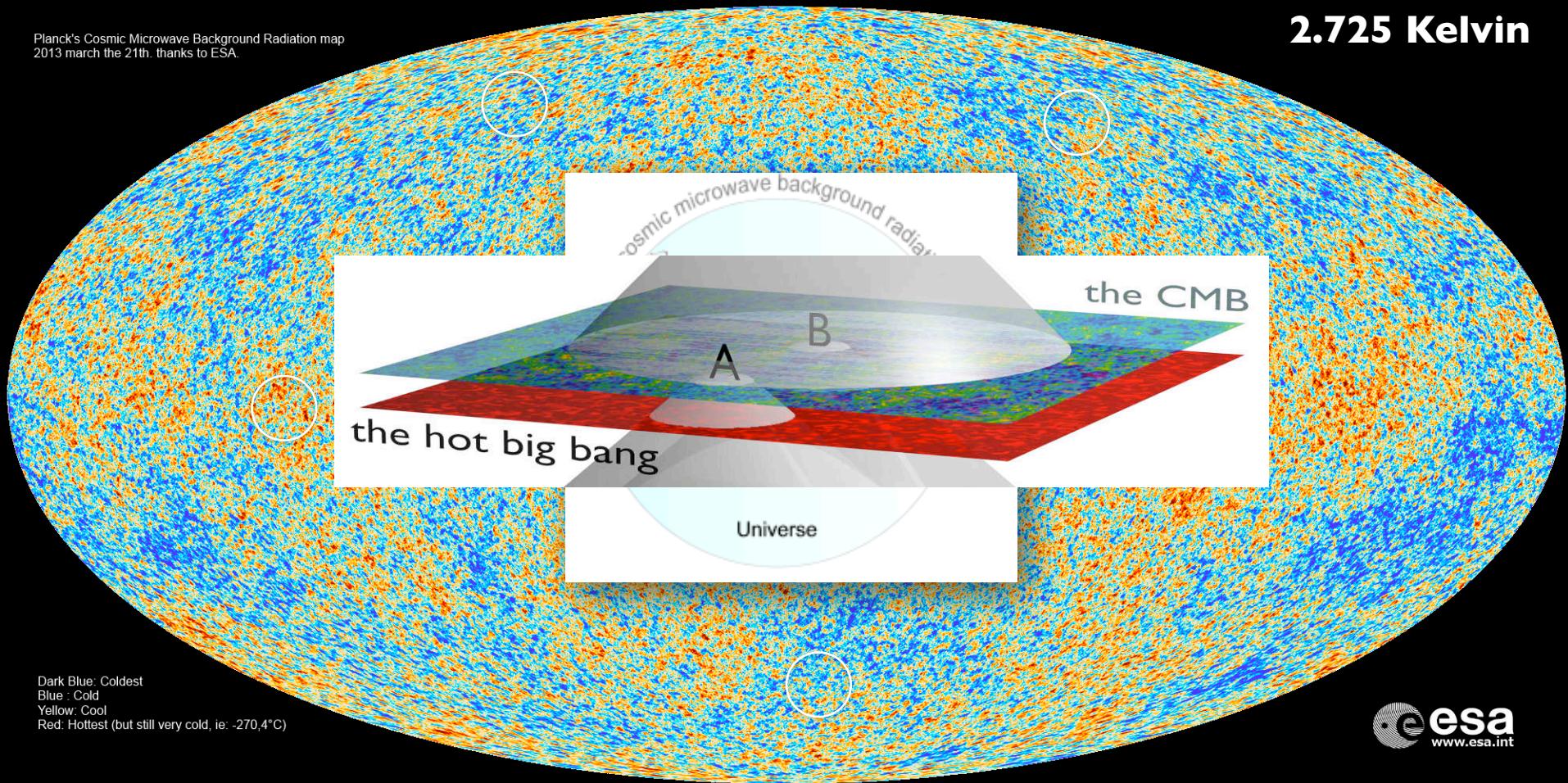
# Evolution



# Smooth and flat

Planck's Cosmic Microwave Background Radiation map  
2013 march the 21th. thanks to ESA.

**2.725 Kelvin**



# Smooth and flat

Hot Big Bang: matter and radiation only  
Equation of state parameter  $w=0$  or  $1/3$   
Ignore dark energy (small correction)

- Size of Last Scattering Surface (LSS)
- Size of horizon at LSS
- *Spatial curvature evolution*

$$ds^2 = -dt^2 + a(t)^2(dr^2 + r^2d\Omega^2)$$

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \sum_i \rho_i - \frac{k}{a^2}$$

$$\dot{\rho} + 3H\rho(1+w) = 0, \quad w \equiv p/\rho$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p) \quad \dot{\Omega}_k = \Omega_m H(1+3w)\Omega_k$$

$$|\Omega_k|_0 \lesssim 10^{-3} \Rightarrow |\Omega_k|_{BBN} \lesssim 10^{-18}$$

$$\Omega_m \equiv \frac{8\pi G}{3H^2}\rho = \frac{\rho}{\rho_{crit}}$$
$$\Omega_k \equiv -\frac{k}{(aH)^2}$$

# Smooth and flat

Hot Big Bang: matter and radiation only  
Equation of state parameter  $w=0$  or  $1/3$   
Ignore dark energy (small correction)

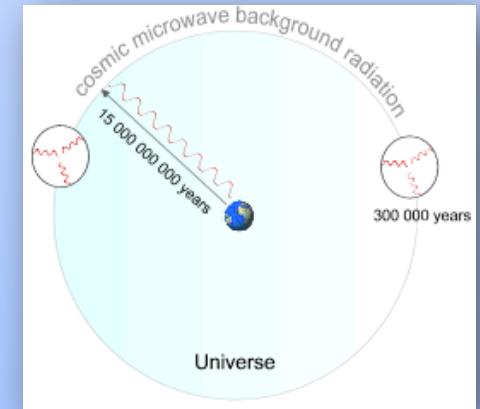
- *Size of Last Scattering Surface (LSS)*
- *Size of horizon at LSS*
- Spatial curvature evolution

$$a(t) \approx \left( \frac{t}{t_0} \right)^{\frac{2}{3(1+w)}} \quad \frac{dr}{dt} = \pm \frac{1}{a(t)} \quad (a_0 = 1)$$

$$d_{LSS} = a_0 \int_{t_{LSS}}^{t_0} dt \left( \frac{t}{t_0} \right)^{-2/3} \approx 3t_0$$

$$d_H = a_{LSS} \int_0^{t_{LSS}} \left( \frac{t}{t_0} \right)^{-1/2} dt = 2t_{LSS}$$

$$\frac{d_H}{d_{LSS}} \sim 10^{-5}$$



# Smooth and flat

*Hot Big Bang: matter and radiation only*

*Equation of state parameter w=0 or 1/3*

*Ignore dark energy (small correction)*

The early universe (LSS) is

- Ridiculously flat
- Ridiculously uniform
- Absence of exotic (beyond SM) particles?
- Bang of the Big Bang?
- Origin of small anisotropies?

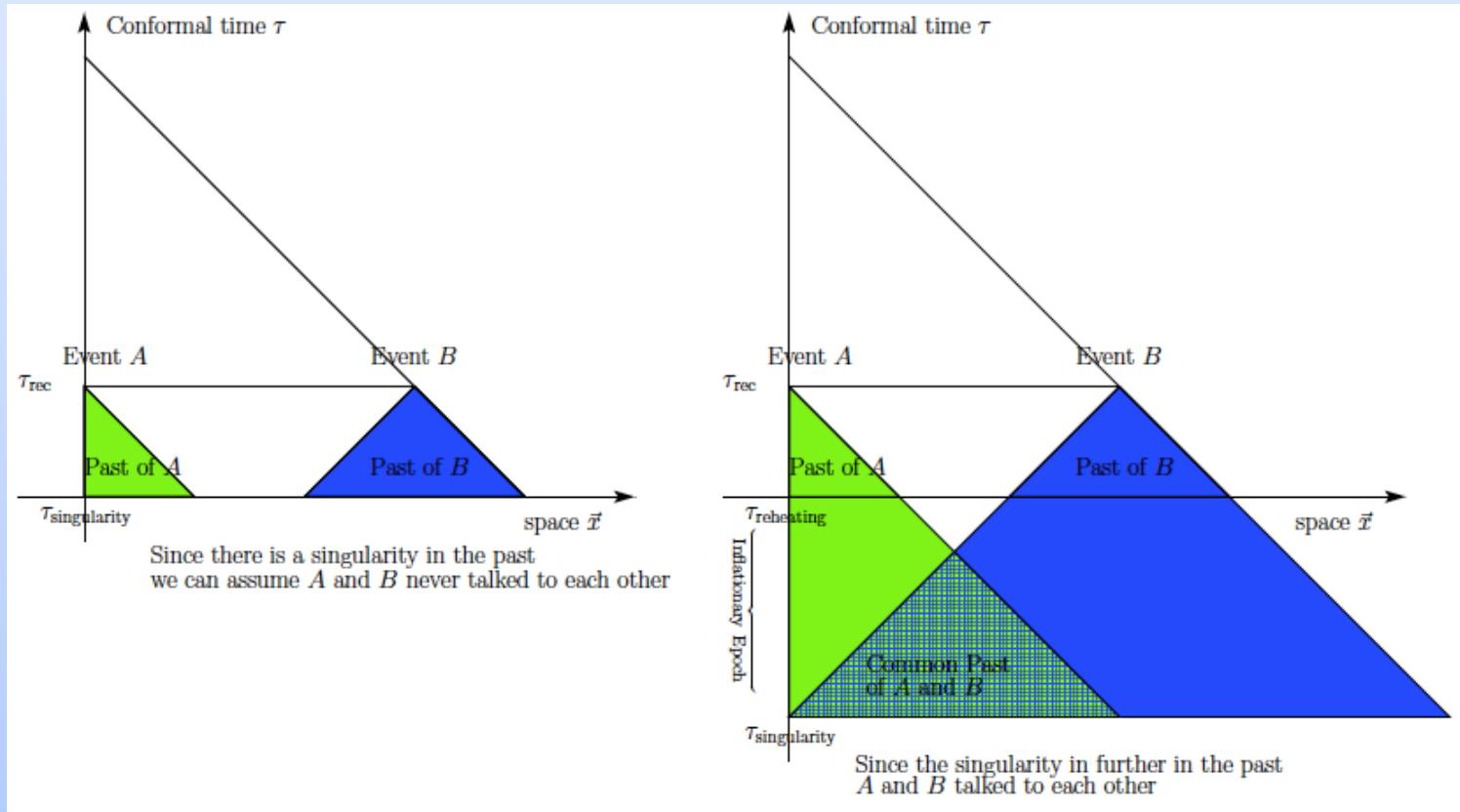
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho(1+3w)$$

$$\dot{\Omega}_k = \Omega_m H(1+3w)\Omega_k$$

**Enter cosmic inflation!**

**Primordial phase with w~1  
Adding extra ‘time’**

# Part II: Inflation

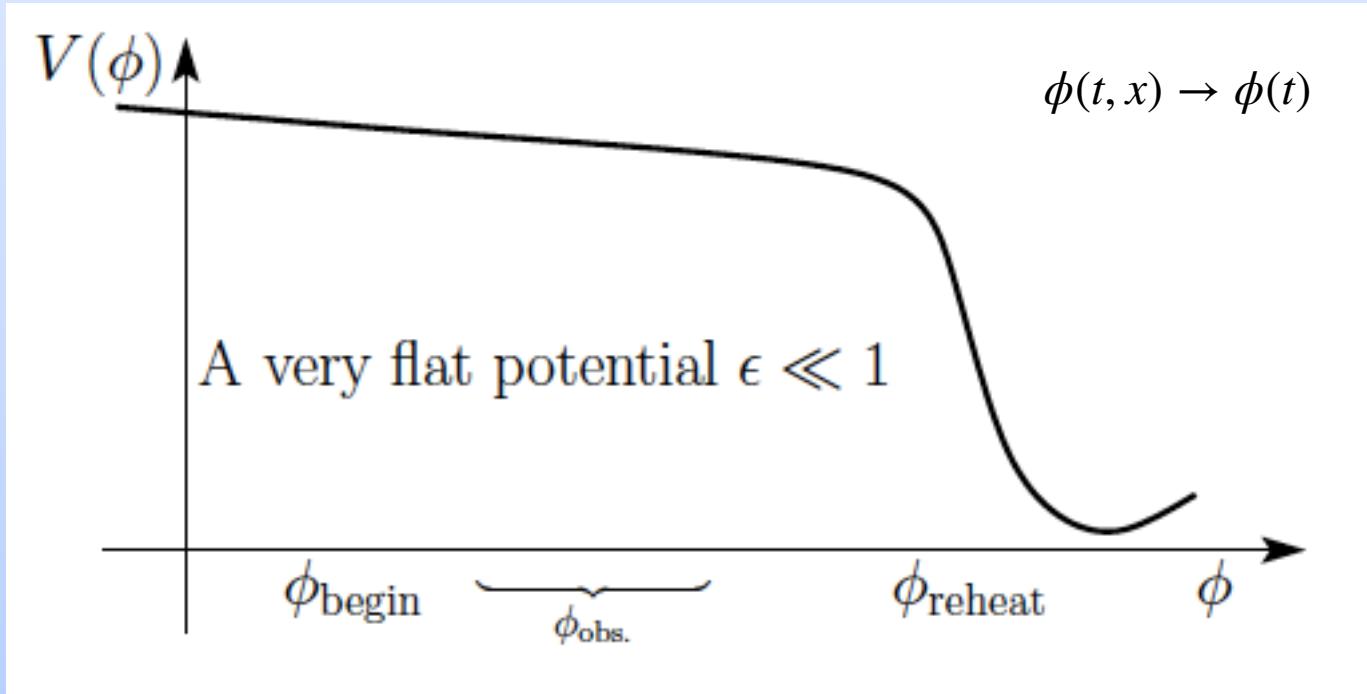


$$w \approx -1 \Rightarrow \dot{\rho} \approx 0 \Rightarrow \dot{H} \approx 0 \Rightarrow a(t) \approx e^{H_I t}$$

$$\tau \equiv \int dt \frac{1}{a} \propto a^{\frac{1}{2}(1+3w)} \rightarrow a(\tau) = -\frac{1}{H_I \tau} \quad \tau \in [-\infty, -\tau_r]$$

**Exponential expansion  
Approximately de Sitter  
Special fluid (scalar field)**

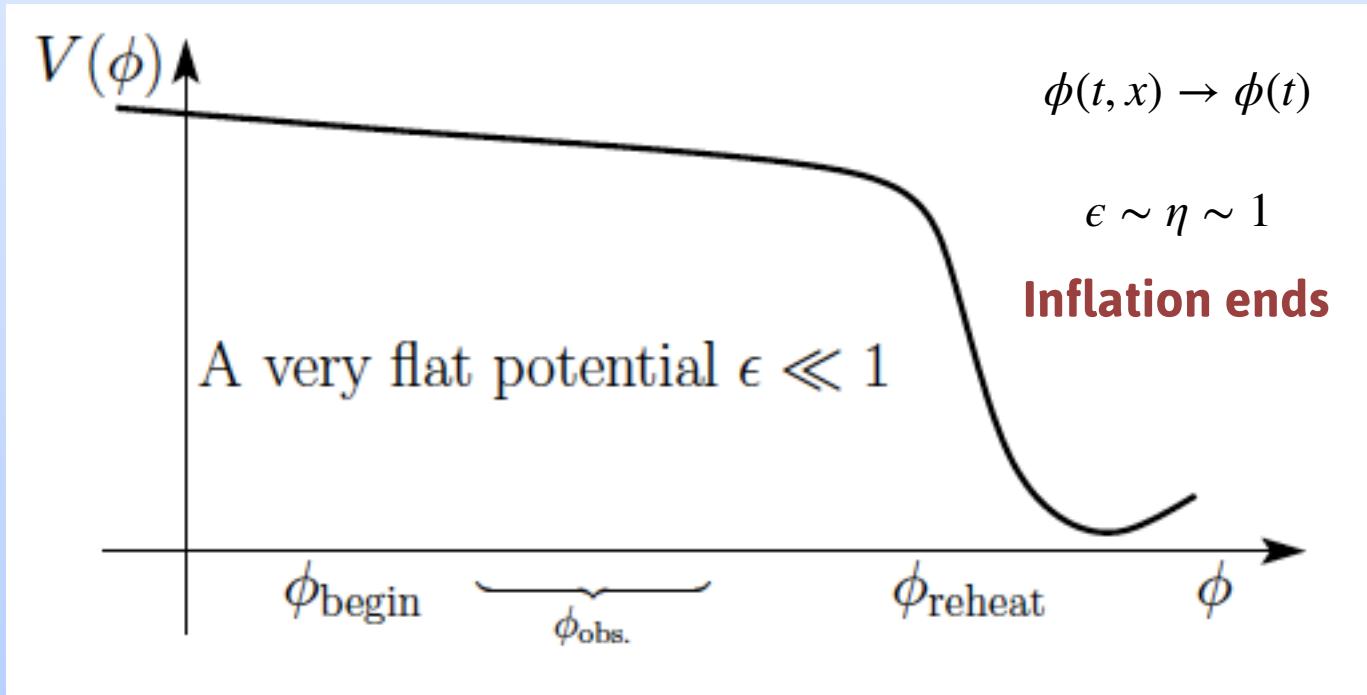
# Inflation and the inflaton



$$\rho_\phi = \frac{1}{2}\dot{\phi}^2 + V(\phi) \quad \frac{1}{2}\dot{\phi}^2 \ll V \Rightarrow w \approx -1 \quad \epsilon \equiv -\frac{\dot{H}}{H^2} \sim \frac{\dot{\phi}^2}{V} \ll 1$$

$$p_\phi = \frac{1}{2}\dot{\phi}^2 - V(\phi) \quad \ddot{\phi} + 3H\dot{\phi} + V' = 0 \quad \dot{\phi} \approx -\frac{V'}{3H} \Rightarrow \eta \equiv -\frac{\ddot{\phi}}{H\dot{\phi}} \ll 1$$

# Slow-roll inflation



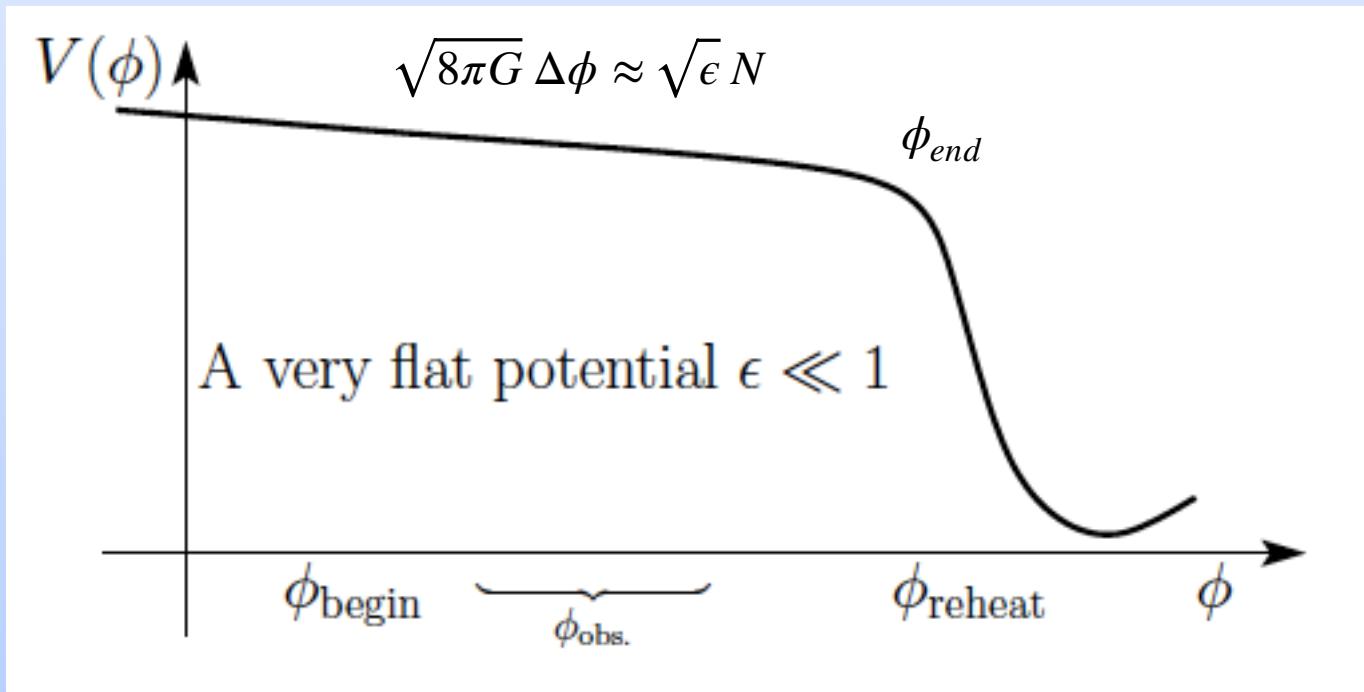
$$\epsilon \equiv -\frac{\dot{H}}{H^2} \ll 1 \qquad \epsilon \approx \frac{1}{16\pi G} \left( \frac{V'}{V} \right)^2$$

$$\eta \equiv -\frac{\ddot{\phi}}{H\dot{\phi}} \ll 1 \qquad \eta \approx \frac{1}{8\pi G} \frac{V''}{V} - \epsilon$$

**Slow-roll conditions: flat potential**

**Quasi-de Sitter  
Prolonged attractor solution**

# Duration and field excursion



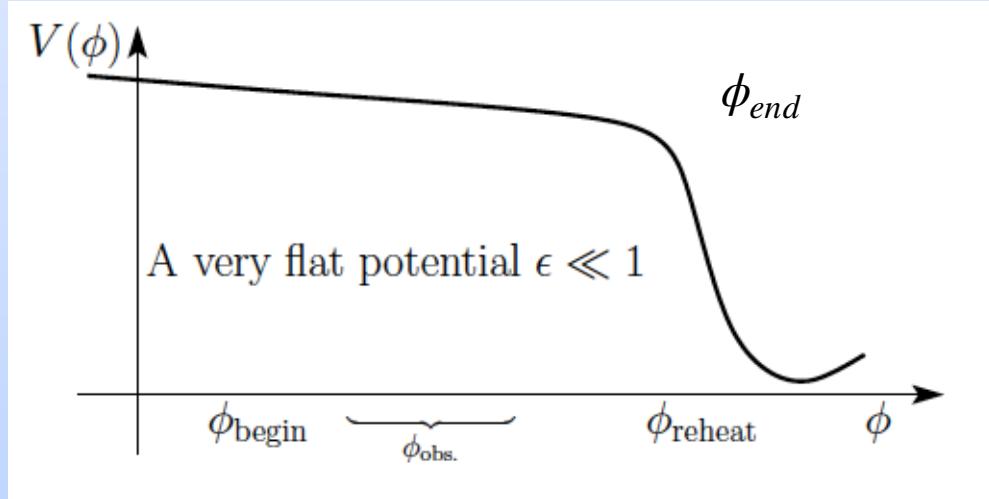
$$N \equiv \ln \frac{a_{end}}{a} \approx \int_t^{t_{end}} H dt' = \int_{\phi}^{\phi_{end}} \frac{H}{\dot{\phi}} d\phi \approx 8\pi G \int_{\phi}^{\phi_{end}} \frac{V}{V'} d\phi$$

$$\Omega_k(a_{end}) = \Omega_k(a_{in}) \frac{a_{in}^2}{a_{end}^2} = \Omega_k(a_{in}) e^{-2N} = \Omega_k(a_0) \frac{H_0^2}{a_{end}^2 H_I^2}$$

**Number of e-folds  
Tied to field excursion**

**Depends on scale of  
inflation, roughly 60 e-  
folds are needed**

# Starting the hot big bang

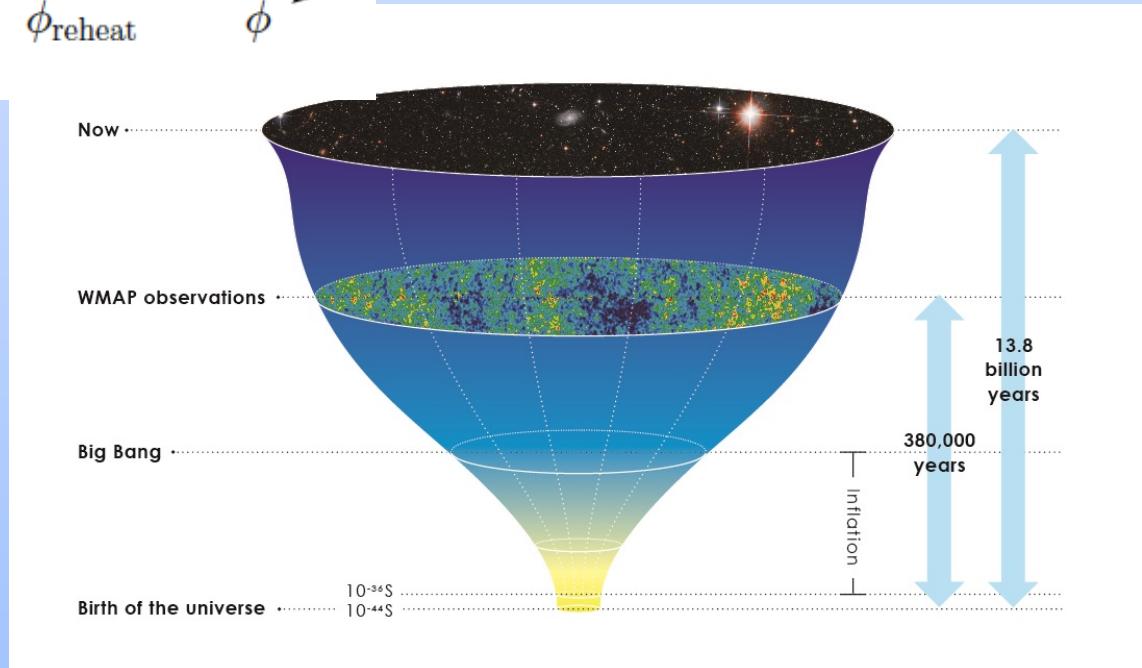


$$T_{BBN} \sim 100 \text{ MeV} \ll H_I \lesssim 10^{14} \text{ GeV}$$

Could be as high as GUT scale

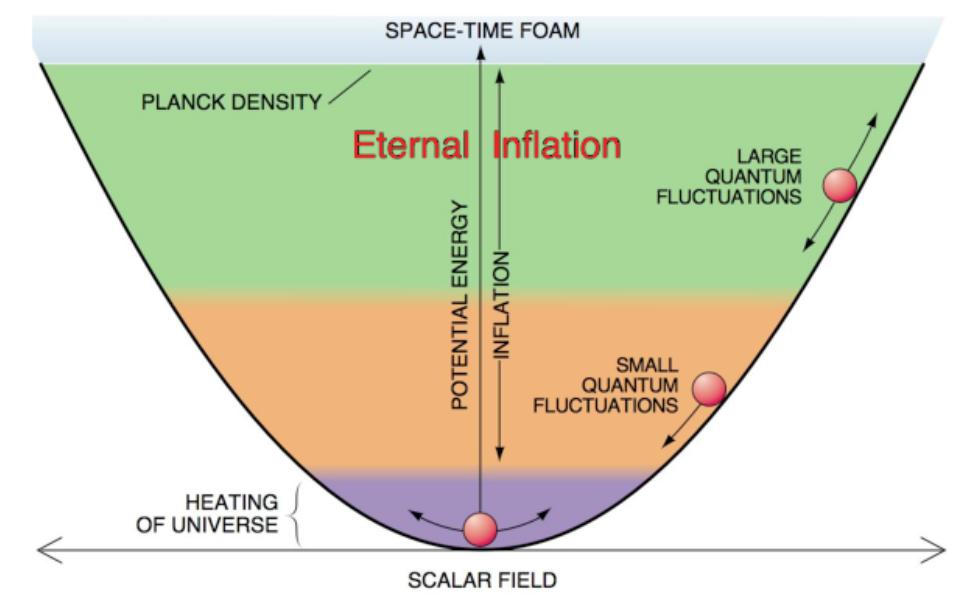
$$\frac{d}{dt}\rho_\phi + (3H + \Gamma)\rho_\phi = 0$$

Non-relativistic matter coupled to SM matter



# Simple example

$$V(\phi) = \frac{m^2}{2}\phi^2$$



$$\epsilon = \eta = \frac{2M_p^2}{\phi^2} \Rightarrow \phi \gg \sqrt{2}M_p$$

$$V \ll M_p^4 \Rightarrow \sqrt{2}M_p \ll \phi \lesssim \left(\frac{M_p}{m}\right) M_p$$

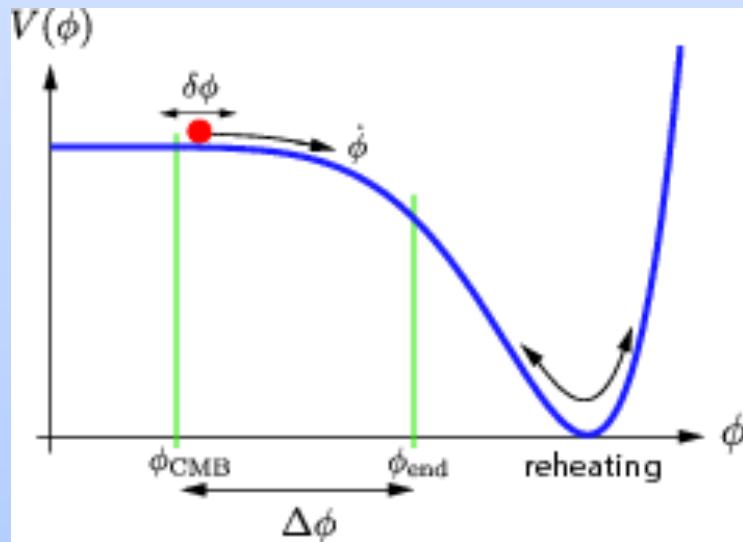
$$N \sim \frac{1}{\epsilon} \Rightarrow \frac{\Delta\phi}{M_p} \sim \sqrt{N} \gg 1$$

**Chaotic inflation**  
**Large field excursion**  
**Ruled out by observations**  
**(stay tuned!)**

# Small field inflation

$$\frac{\Delta\phi}{M_p} \approx \sqrt{\epsilon} N \Rightarrow \epsilon \lesssim N^{-2}$$

**Small field inflation**



$$V(\phi) = V_0 \left( 1 - \left( \frac{\phi}{\mu} \right)^2 \right)$$

$$\epsilon = \frac{M_p^2 \phi^2}{\mu^4} \quad \mu \gtrsim M_p N$$

**Guarantees enough e-folds**

**Inflationary paradigm: small/large field, convex/concave potentials. As long as the slow-roll conditions are satisfied it produces a large and smooth universe.**  
**Theoretical consistency, predictions?**

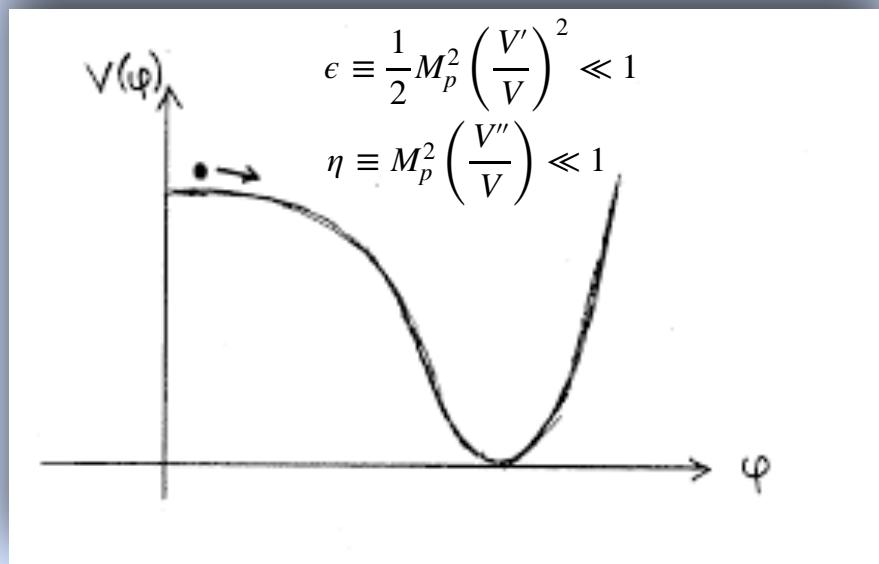
# Challenge and opportunity: UV sensitivity

$$\frac{H_I}{M_p} \lesssim 10^{-4} \sqrt{\epsilon} \lesssim 10^{14} \text{ GeV}$$

$$\frac{\Delta\phi}{M_p} \approx N\sqrt{\epsilon}$$

- Scale of inflation could be GUT
- Large field excursion probes UV
- Flat potential is fine-tuned

$$m \ll H_I \ll \Lambda_c$$



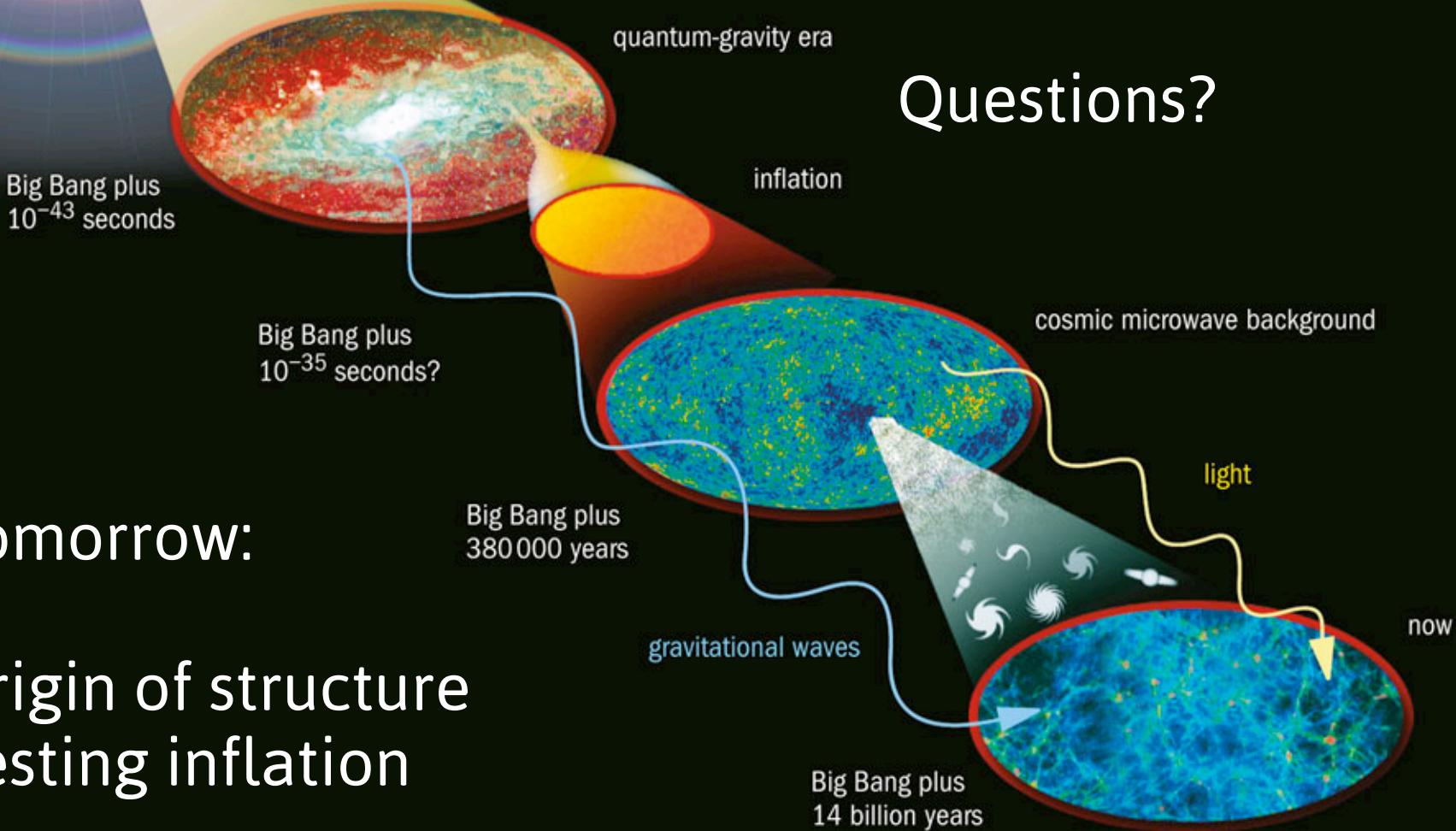
Inflation explains  
large, smooth,  
expanding hot big  
bang universe

Questions?

Tomorrow:

Origin of structure  
Testing inflation

?



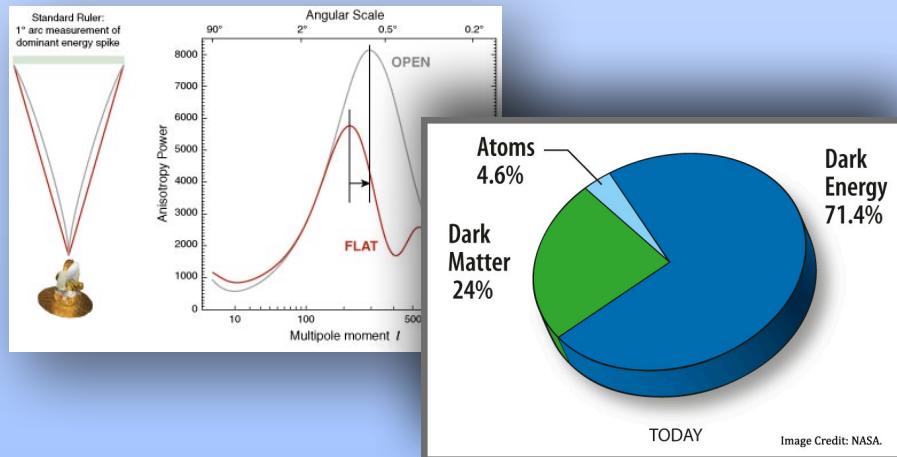
# Plan

## Yesterday

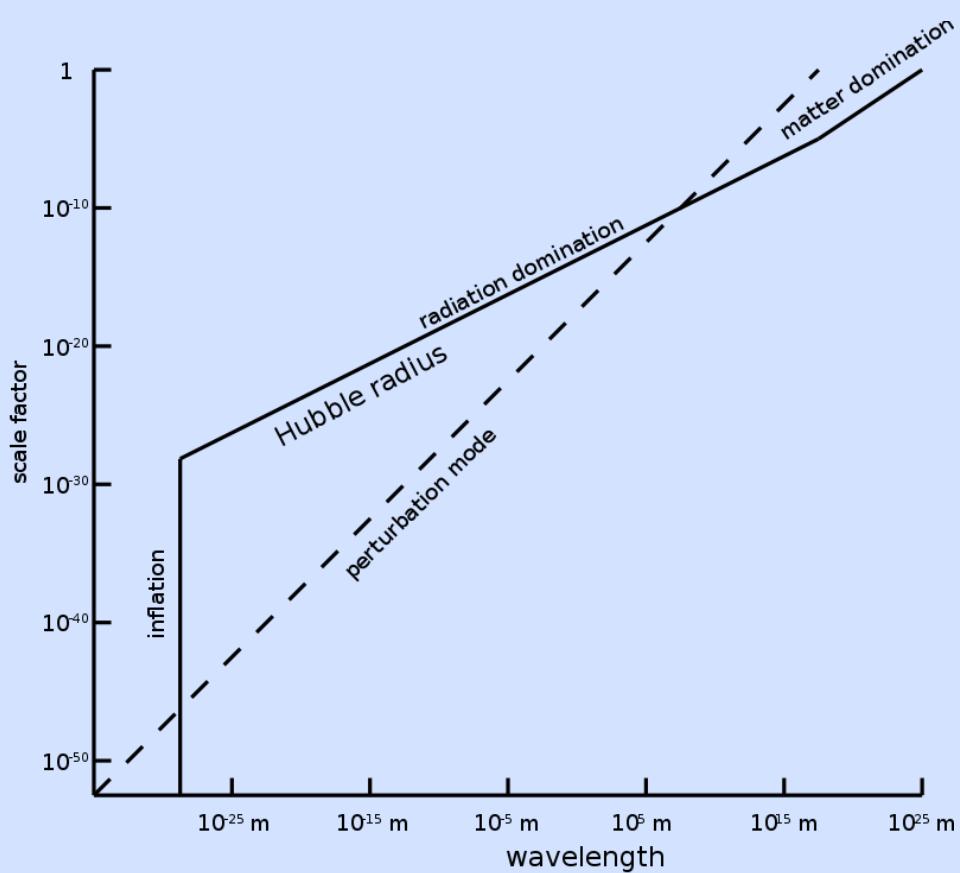
- Initial conditions of the hot Big Bang
- The paradigm of cosmological inflation

## Today

- Inflation and the origin of structure
- The cosmological accelerator: constraining inflation



# The origin of structure

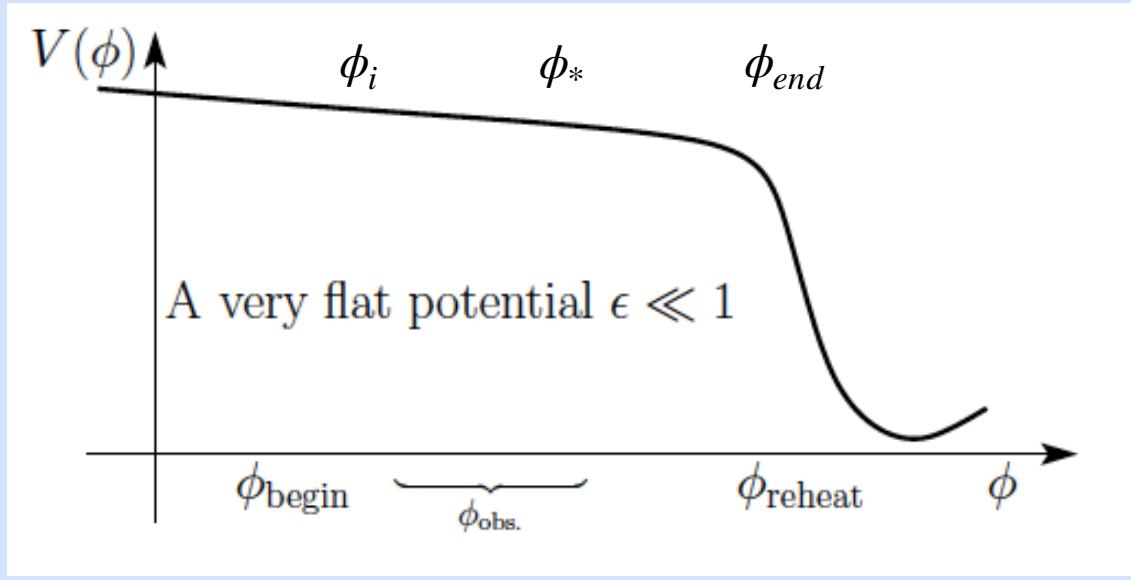


Primordial phase of exponential expansion, caused by dominating vacuum energy of a scalar field

Predicts almost scale-invariant spectrum of primordial density fluctuations (scalar and tensor)



# Probing inflation



$$\frac{\delta\rho}{\rho} = \frac{\delta a}{a} \propto \frac{V}{V'} \delta\phi = \frac{H}{M_p \sqrt{\epsilon}}$$

$$P_S(k) \propto \frac{H_*^2}{M_p^2 \epsilon_*} \left( \frac{k}{k_*} \right)^{n_s - 1}$$

$$P_T(k) \propto \frac{H_*^2}{M_p^2} \left( \frac{k}{k_*} \right)^{n_T}$$

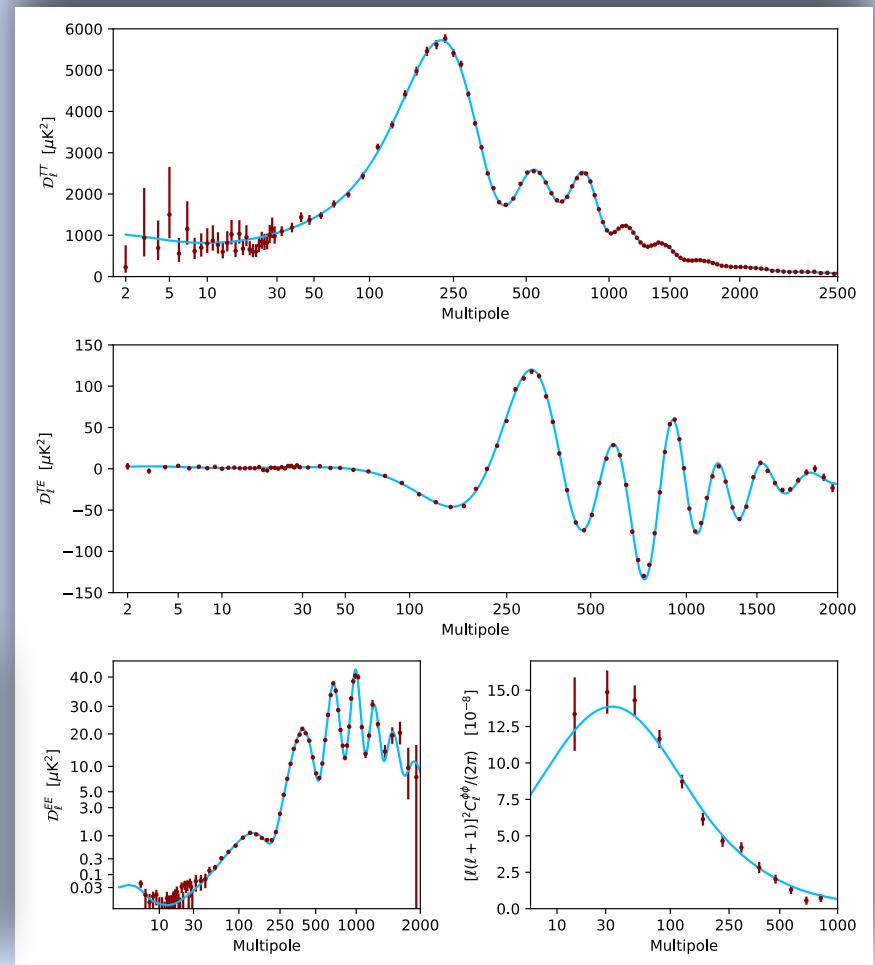
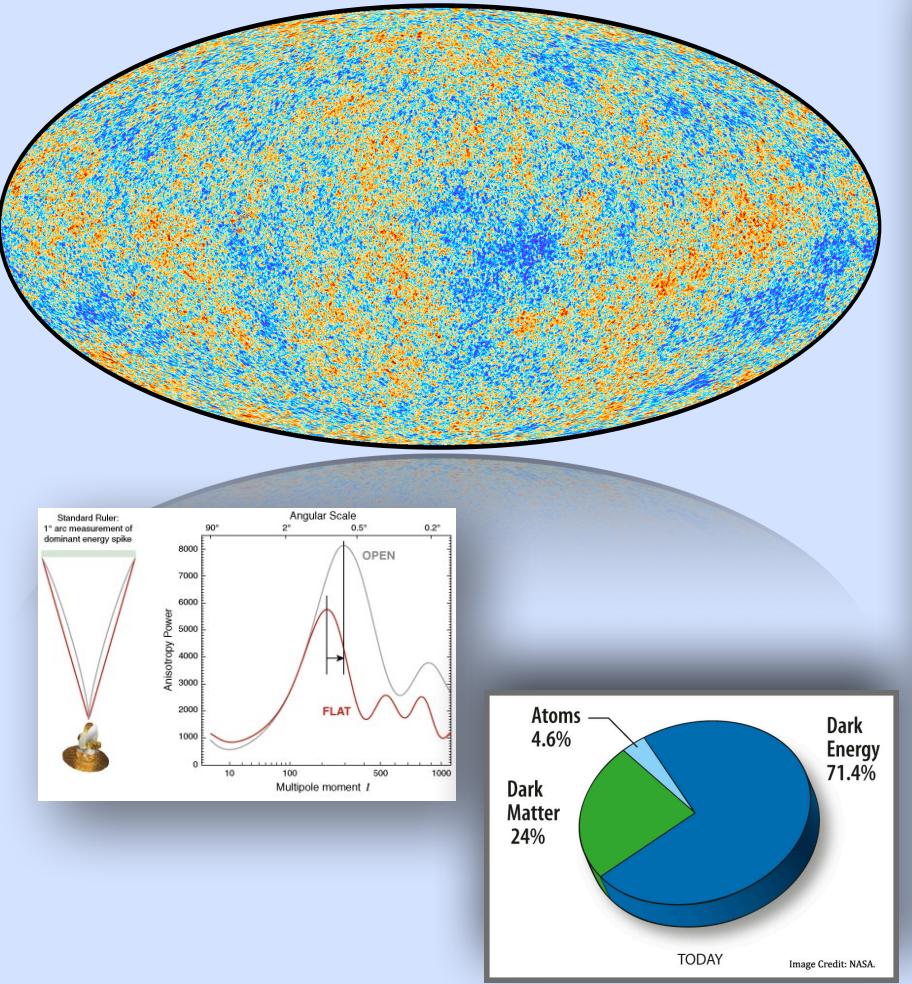
$$n_s - 1 = 4\epsilon - 2\eta , n_T = 1 - 2\epsilon$$

## Inflationary predictions:

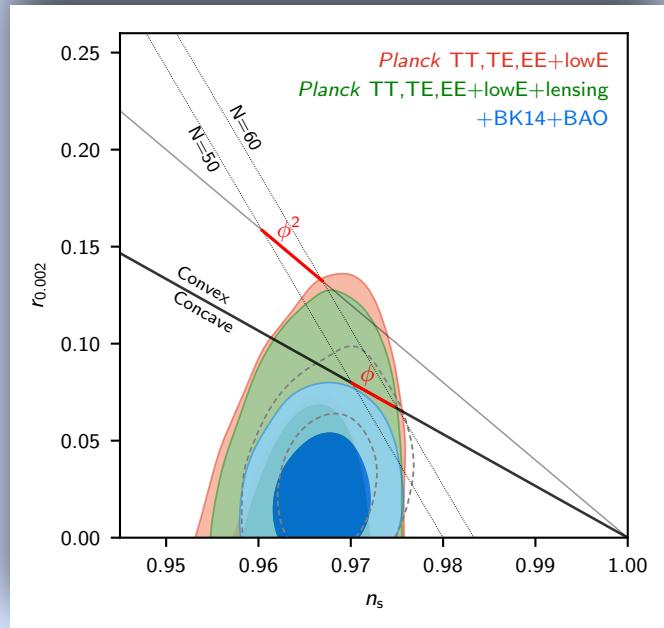
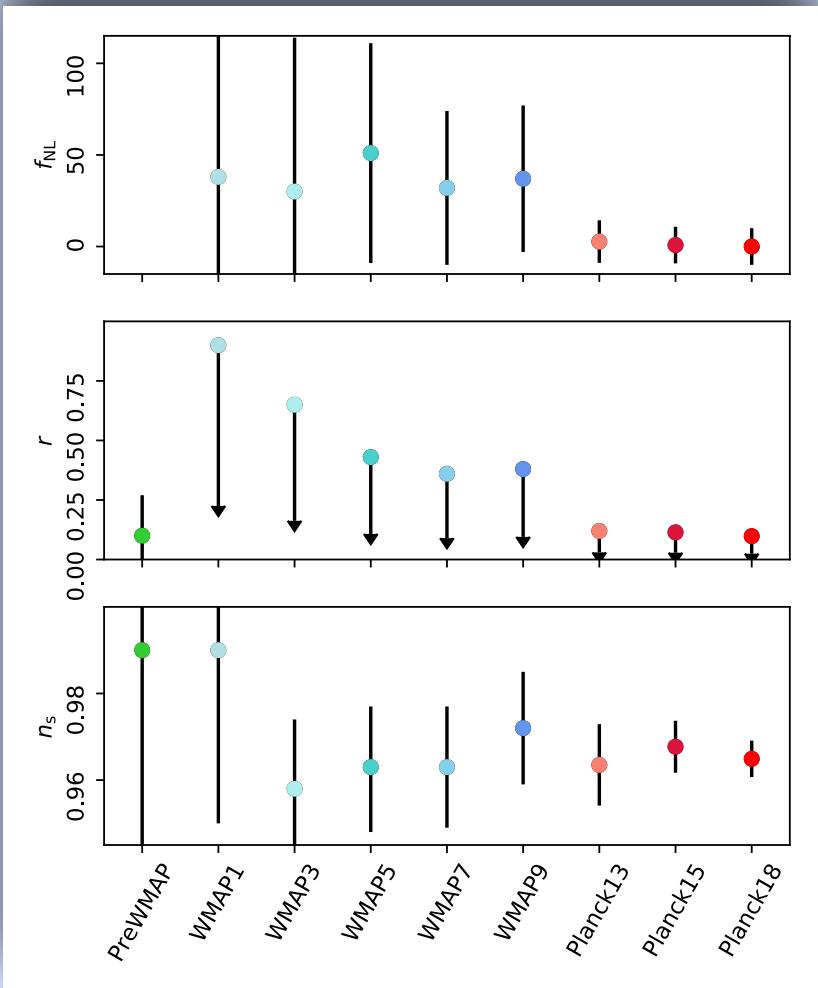
- **Nearly Gaussian spectrum**
- **Almost scale invariant**
- **Details (amplitudes and spectral indices) depend on potential!**

$$\epsilon \approx \frac{1}{16\pi G} \left( \frac{V'}{V} \right)^2 \quad \eta \approx \frac{1}{8\pi G} \frac{V''}{V} - \epsilon$$

# Planck 2018



# Planck 2018

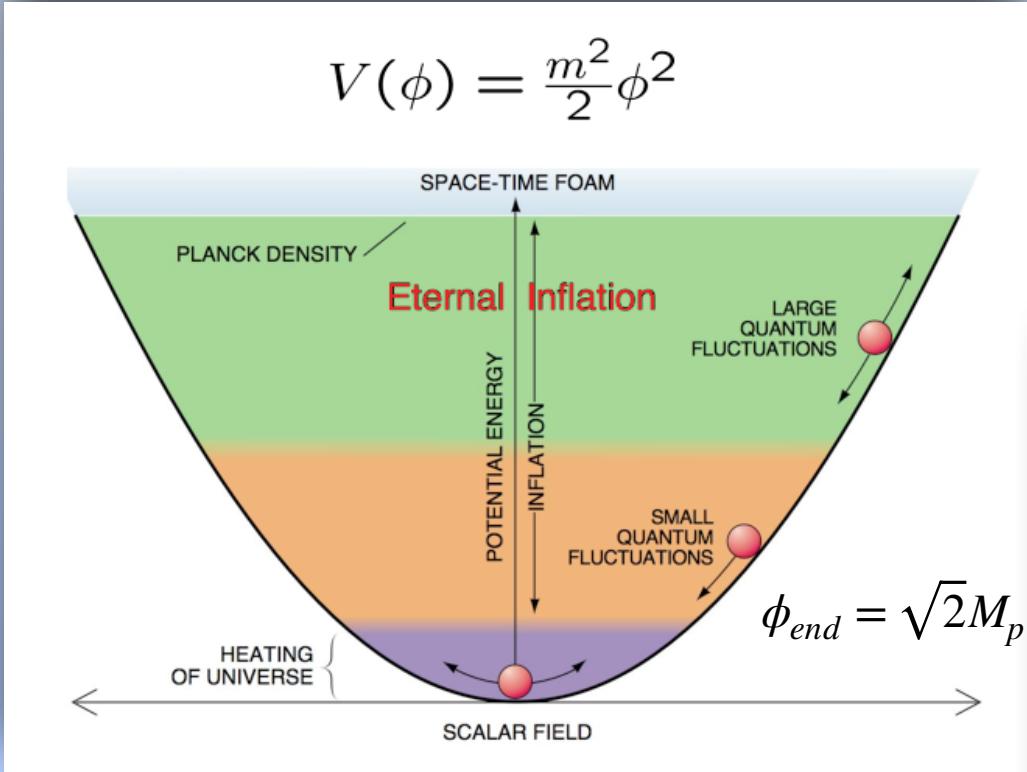


Gaussian almost scale invariant  
No tensors ( $r < 0.1$ ), no running,  
no features, concave potential:

Plain vanilla single field

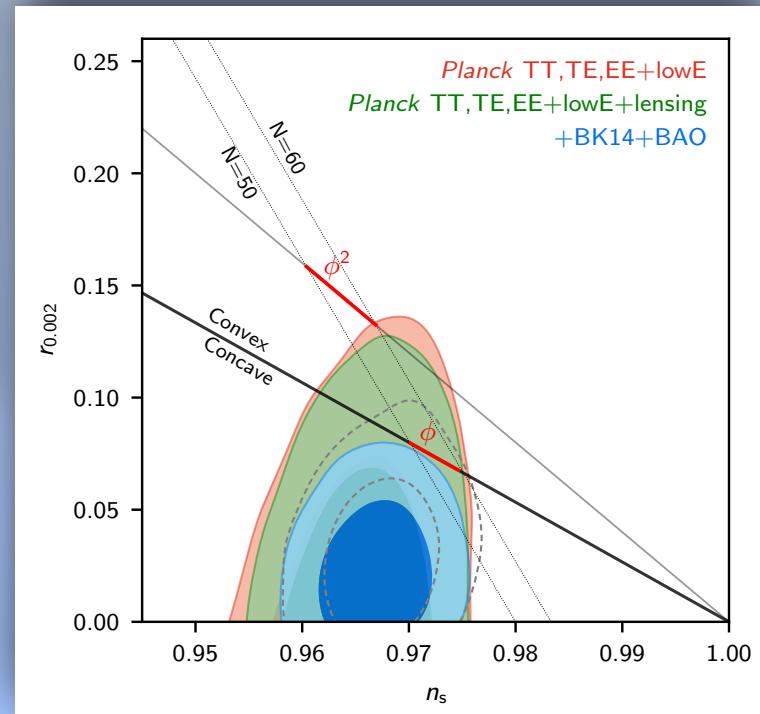
# Back to simplest example

$$V(\phi) = \frac{m^2}{2}\phi^2$$



$$\epsilon = \eta = \frac{2M_p^2}{\phi^2} \Rightarrow \phi \gg \sqrt{2}M_p$$

$$V \ll M_p^4 \Rightarrow \sqrt{2}M_p \ll \phi \lesssim \left(\frac{M_p}{m}\right) M_p$$

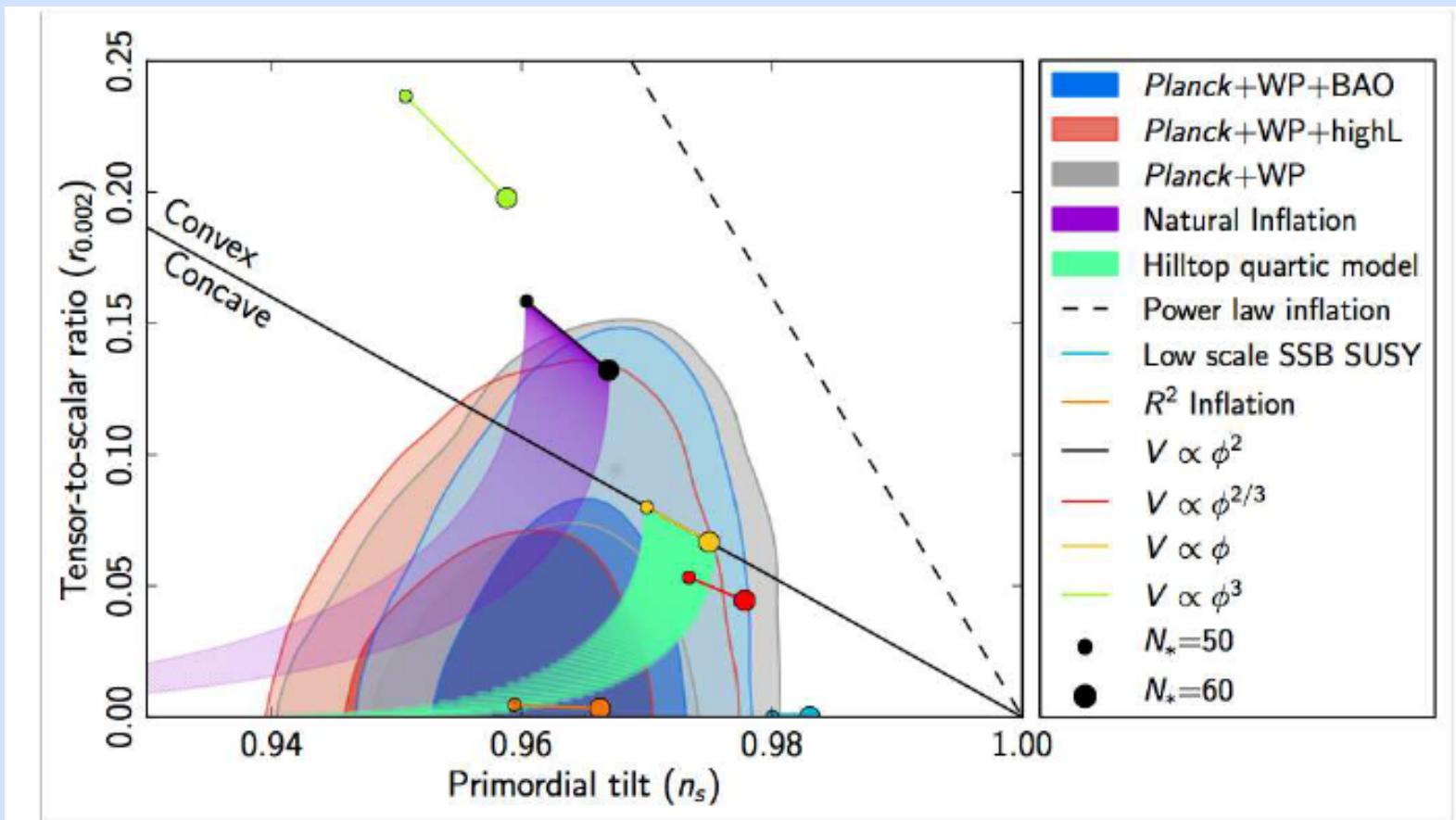


$$n_s \approx \frac{N-2}{N}$$

$$r \equiv 16\epsilon = \frac{8}{N}$$

**Almost ruled out!**

# Planck 2018



Gaussian almost scale invariant. No tensors ( $r < 0.1$ ), no running, no features, concave potential. Plain vanilla single field

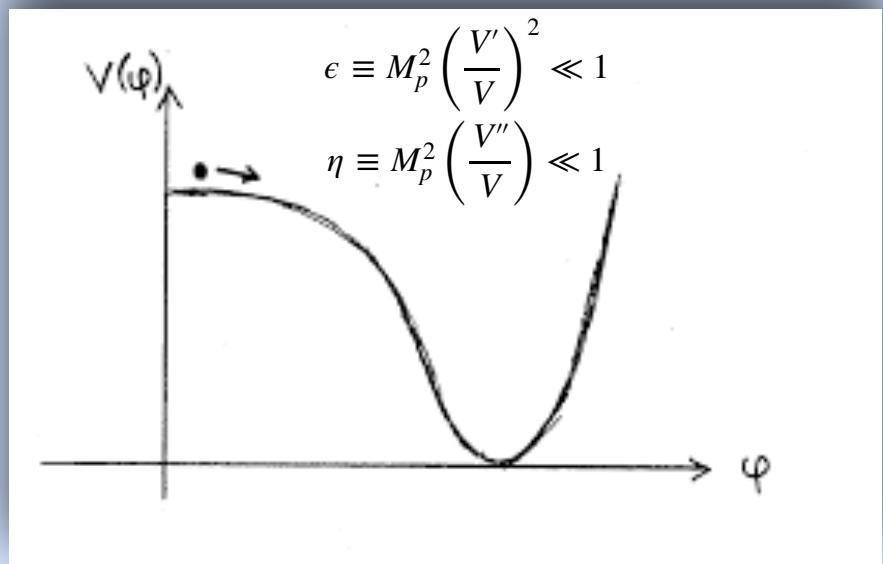
# Reminder: UV sensitivity

$$\frac{H_I}{M_p} \sim 10^{-4} \sqrt{r/16}$$

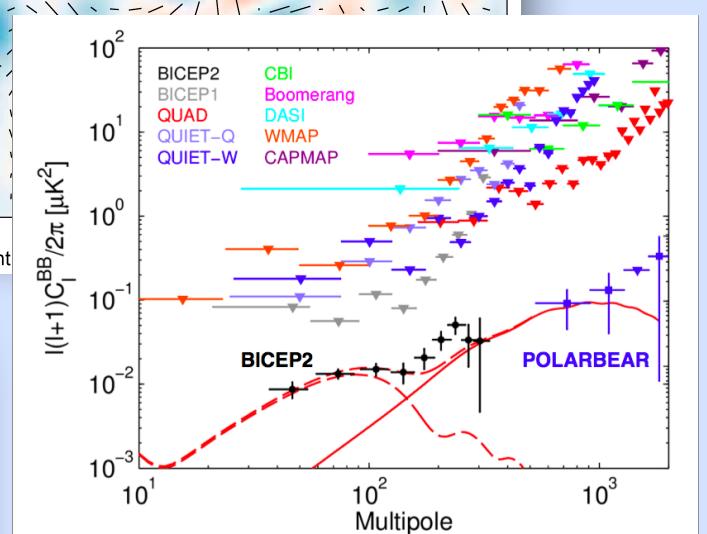
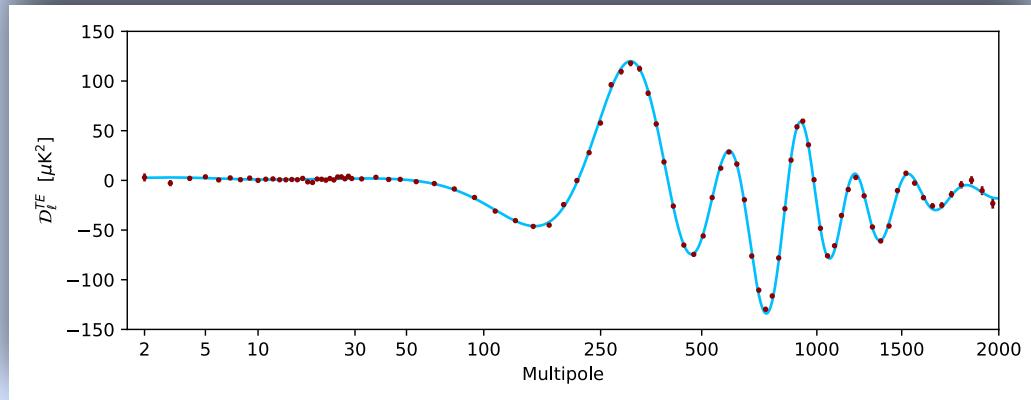
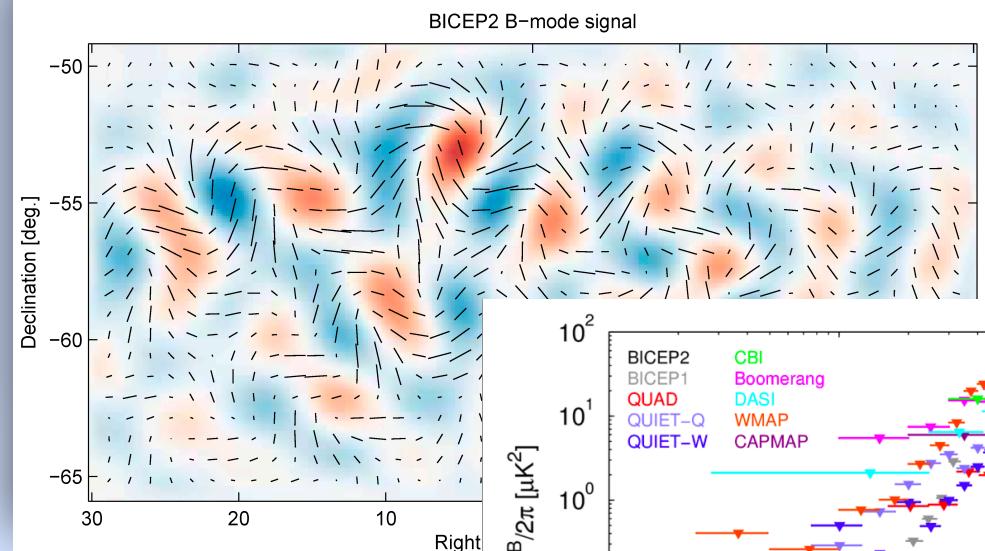
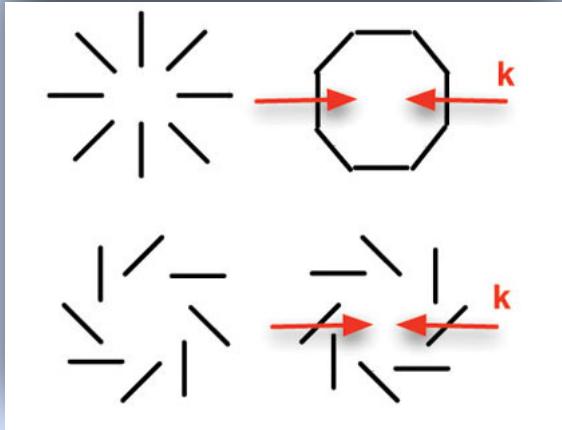
$$\frac{\Delta\varphi}{M_p} \approx N\sqrt{r/8} \approx \sqrt{r/0.01}$$

- Scale of inflation could be GUT
- Large field excursion probes UV
- Flat potential is fine-tuned

$$m \ll H_I \ll \Lambda_c$$



# Polarisation and tensors



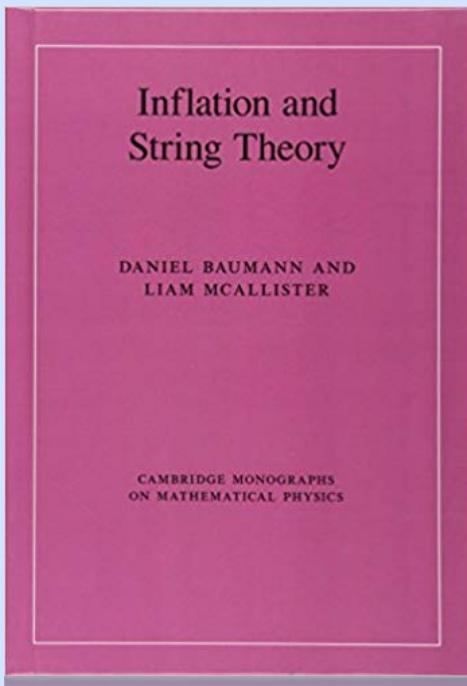
**Gravitational waves: B-mode**

$$r \gtrsim 10^{-3}$$

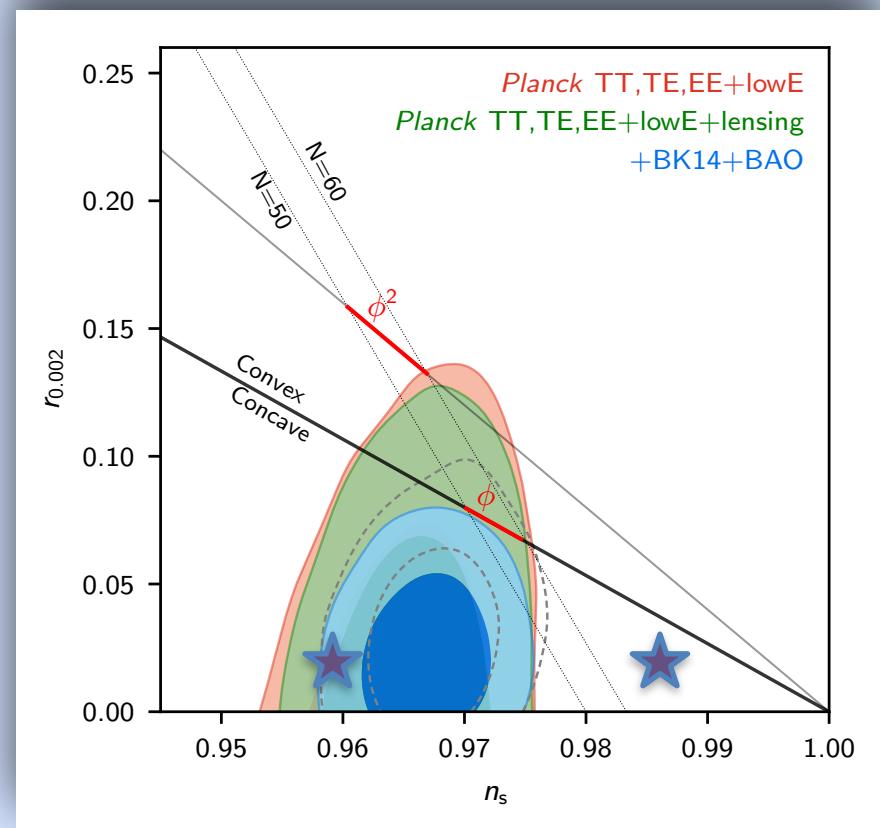
# Inflation in theory

Beyond phenomenology: consistent UV embedding of inflation  
Requires understanding of GUT and Planck scale physics

**Goal: more predictive power**



**Hard! No fully understood case**



# String landscape



Specifically: de Sitter vacua

- No supersymmetry
- NO-GO theorems
- Exotic ingredients
- Claims of success, but...

De Sitter swampland conjectures:

$$M_p \left| \frac{V'}{V} \right| \gtrsim \mathcal{O}(1) \text{ or } -M_p^2 \frac{V''}{V} \gtrsim \mathcal{O}(1)$$

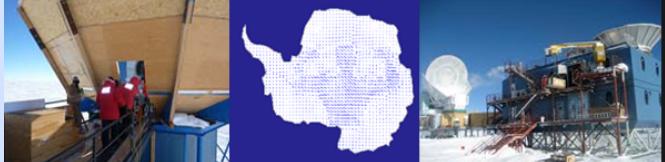
Plethora of 4d flat vacua with particle physics similar to SM

Also true for cosmology?  
Weak gravity/distance conjectures



**Anthropic multiverse  
or de Sitter swamp?**

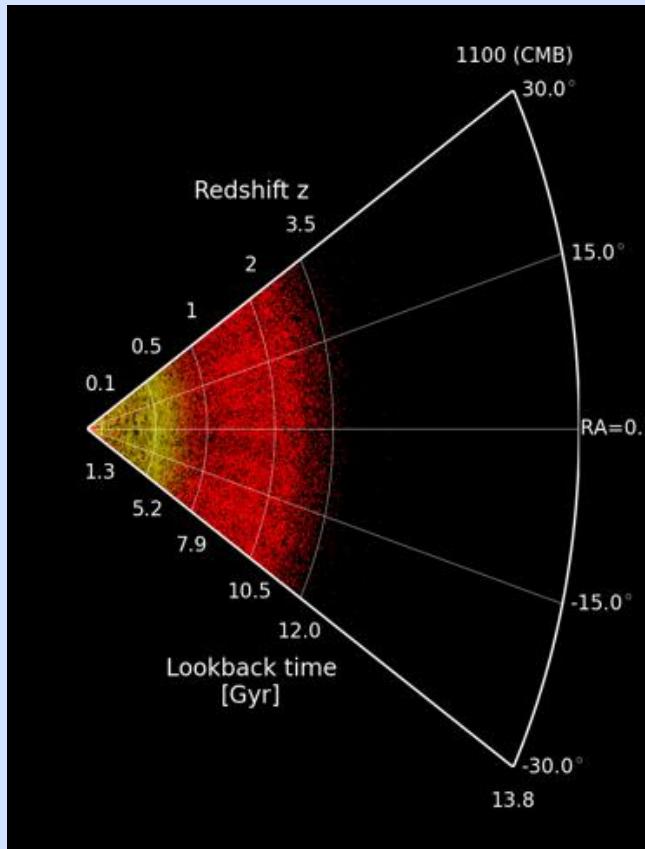
# Observational prospects



BICEP2



Keck Array



CMB

- Tensors
- NonGaussian
- Precision



21cm



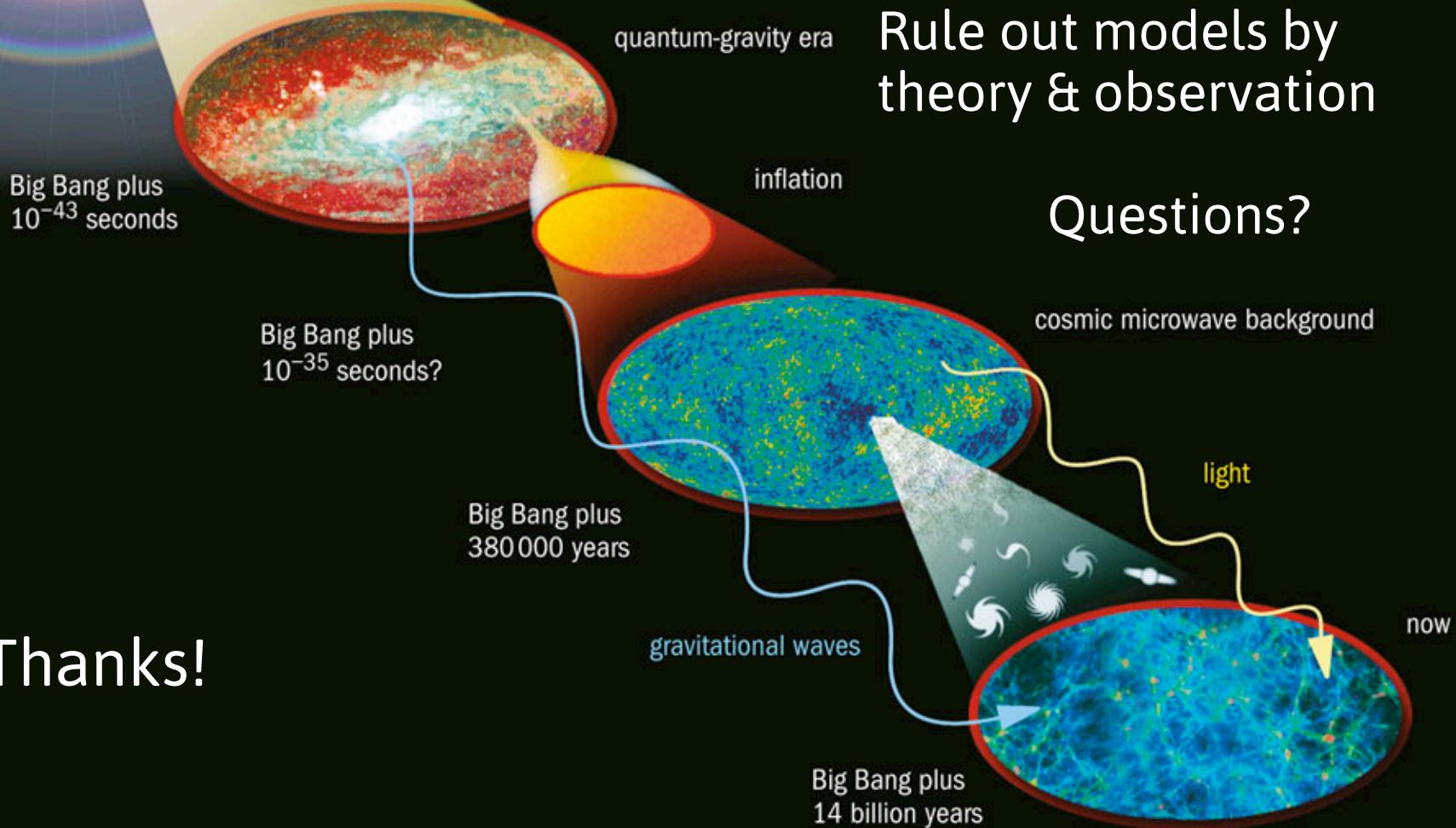
LiteBIRD



LSS



?



Inflation explains large, smooth, expanding universe, with small inhomogeneities.

Rule out models by theory & observation

Questions?

Thanks!