## Changes to paper since last update <br> (see presentation 26-3-18)

With comments from a.o. Jochen and Klaus

Lepcol meeting 23-4-18

## Changes repeated in paper

- Gridpix $\rightarrow$ GridPix
- Timepix1 $\rightarrow$ Timepix
- Removed many hyphens
- Marked changes in red (links to figures are also red)


## Abstract

- Text added


## Abstract

A GridPix readout for a TPC based on the Timepix3 chip is developed for future applications at a linear collider. The GridPix detector consists of a gaseous drift volume read out by a single Timepix3 chip with an integrated amplification grid. Its performance is studied in a test beam with 2.5 GeV electrons. The GridPix detector detects single ionization electrons with high efficiency. The Timepix3 chip allowed for high sample rates and time walk corrections. Diffusion is found to be the dominating error on the track position measurement both in the pixel plane and in the drift direction, and systematic distortions in the pixel plane are below $10 \mu \mathrm{~m}$. Using a truncated sum, an energy loss (dE/dx) resolution of $4.1 \%$ is found for an effective track length of 1 m .

## Introduction

## 1. Introduction

In the context of a Time Projection Chamber for a future linear collider a gaseous pixel detector is developed based on the Timepix3 chip. The GridPix single chip detector discussed here, allows for a detection of single electrons with a granularity of $256 \times 256$ pixels with a pitch of $55 \mu \mathrm{~m} \times 55 \mu \mathrm{~m}$. By counting the number of pixel hits, the number of clusters can be estimated allowing for a precise measurement of the energy loss $\mathrm{dE} / \mathrm{dx}$.

Since the invention of the device [1, 2], a series of developments have taken place that culminated in GridPix detectors using the Timepix chip [3]. In this paper the first results using a Timepix3 based Gridpix detector will be described. In the design of the detector, special attention has been given to minimize the distortions in the pixel plane and drift direction in order to meet the tracking precision needed for a TPC at a linear collider. The device can also be applied for medical imaging, proton radiotherapy or used in other particle physics experiments (4]. Here testbeam results taken at the ELSA test beam facility in Bonn will be presented. Some results using this device in a laser setup were presented at TIPP17 [5].

## 2. Description of the GridPix device

## Section 2

## - Changed text

A GridPix is a CMOS pixel readout chip with a gas amplification grid added by photolithographic postprocessing techniques [3. In this work, the GridPix consists of a Timepix3 chip [6] with a $4 \mu \mathrm{~m}$ thick Silicon-rich Nitride protective layer, and $50 \mu \mathrm{~m}$ high SU8 pillars that support the $1 \mu \mathrm{~m}$ thick Al grid that has $35 \mu \mathrm{~m}$ diameter circular holes aligned to the pixels. The growing of the protection layer of Timepix chips has been further optimized at the Fraunhofer Institute for Reliability and Microintegration (IZM) in Berlin, making the device more spark proof. An ionizing particle will liberate electrons in the TPC drift volume that will drift towards the grid and enter the avalanche region. The avalanche yields an electronic signal on the pixel. The Timepix3 chip has low noise ( $\approx 70 \mathrm{e}^{-}$) and allows per pixel for a precise measurement of the arrival time and the time over threshold using a TDC (clock frequency 640 MHz ). For the readout the SPIDR software is used [7].

In Figure 1 a cross section of the GridPix detector $(14.1 \mathrm{~mm} \times 14.1 \mathrm{~mm})$ located in a small drift volume is shown. The box has a length of 69 mm , a width (not shown) of 42 mm and a height of 28 mm with a maximum drift length of about 20 mm . The beam enters the drift volume through the 5 mm thick synthetic window from the right side. The electric drift field is defined by a series of parallel conductive strips in the cage and is set to about $280 \mathrm{~V} / \mathrm{cm}$. On the guard plane - located 1 mm above the grid - a voltage is applied that matches the local drift voltage.

## Section 3 (1/2)

## 3. Testbeam measurement

In July 2017, measurements were performed at the ELSA test beam facility in Bonn. The ELSA storage ring delivered a beam of 2.5 GeV electrons at a rate set to a maximum of 10 kHz . Each beam cycle of approximately 6.5 s contained one spill of about 5 s . To acquire a precise reference track, a silicon pixel tracking telescope was introduced in the setup as shown in Figure 2. Electrons from the beam first passed through a scintillator that was used to provide a trigger signal. This was followed by the tracking Mimosa telescope, consisting of 6 silicon detection planes mounted on a slider stage with each $1152 \times 576$ pixels sized $18.4 \mu \mathrm{~m} \times 18.4 \mu \mathrm{~m}$. Finally, the beam crosses the gas volume of the GridPix detector. The whole GridPix detector was mounted on a remote controllable rotation stage. On the last telescope plane an inactive FEI4 chip [8] was present that caused multiple scattering of the beam corresponding to a r.m.s. of 0.25 mm at the GridPix detector.

## Section 3 (2/2)

Both the Mimosa telescope and the Timepix3 chip were operated in data driven mode. For synchronization, triggers were numbered by a Trigger Logic Unit (TLU) 9 and saved in the two data streams. The Mimosa chips were continuously read out with a rolling shutter taking ${ }^{1}$ $115.2 \mu \mathrm{~s}$, meaning that a single frame can contain multiple triggers. The Timepix3 hits are attributed to a single trigger by considering all hits within 400 ns of a trigger.

During data taking the gas volume of the GridPix detector was flushed with a premixed gas consisting of $95 \%^{1}$ Ar, $3 \% \mathrm{CF}_{4}$, and $2 \% \mathrm{iC}_{4} \mathrm{H}_{10}$. This gas - called T2K TPC gas - is suitable for a large TPC because of the low diffusion in a magnetic field and its relatively large drift velocity. The cathode and guard voltage of the GridPix were set such that the electric field was $280 \mathrm{~V} / \mathrm{cm}$, near the value ${ }^{1}$ at which the drift velocity is maximal for this gas. With Magboltz the drift velocity is predicted to be $75 \mu \mathrm{~m} / \mathrm{ns}$ [10]. To achieve a high efficiency, the grid voltage was set at 350 V . The threshold per pixel was put at $800 \mathrm{e}^{-}$to reduce the number of noise hits to a minimum. The temperature and pressure at time of data taking were stable at 301.6 K and 1034.20 mbar . The Oxygen concentration in the gas was 211 ppm . In Table 1 the parameters of the analyzed run are summarized.

From the measured time of arrival of the Timepix hits, the $z$-position is calculated using the predicted drift velocity of $75 \mu \mathrm{~m} / \mathrm{ns}$. This value is found to be consistent, but because of systematic uncertainties there was no attempt at a precise determination.

## Table 1

| Table 1: Parameters of the analyzed run. The error on temperature and pressure indicates the spread during the run. |  |
| :---: | :---: |
| Clock time length | 60 minutes |
| Triggers | 4733381 |
| $V_{\text {grid }}$ | 350 V |
| $E_{\text {drift }}$ | $280 \mathrm{~V} / \mathrm{cm}$ |
| Rotation (z-axis) | 17 degree |
| Rotation ( $y$-axis) | 0 degree |
| Threshold | $800 \mathrm{e}^{-}$ |
| Temperature | $(301.63 \pm 0.08) \mathrm{K}$ |
| Pressure | $(1034.20 \pm 0.05) \mathrm{mbar}$ |
| Oxygen concentration | 211 ppm |

## 4. Track reconstruction and event selection

### 4.1. Track fitting

To reconstruct a track, a straight line is fitted to the hits. The $x$-axis is chosen parallel to the beam, and the $y, z$-axes are perpendicular to the beam. The drift direction is parallel to the $z$-axis. Tracks are fitted using a linear regression fit in $y(x)$ and $z(x)$. Hits are assigned errors in the 2 directions perpendicular to the beam $\sigma_{y}$, $\sigma_{z}$. This will be discussed in detail in section 5.3 and 5.4 ,

To achieve an accurate reconstruction of the tracks, the telescope and the GridPix detector have to be aligned. In a first step, the positions of the 6 telescope planes are independently aligned. The planes are placed perpendicular to the beam, and their position along the beam is measured. The 5 rotations and $4 \times 2$ shifts are iteratively determined from data. In the second step, the GridPix detector is aligned to the beam by rotating it along 3 axes and measuring the shifts in the directions perpendicular to the beam.

Since the telescope track is affected by multiple scattering, the most precise track fit is obtained by fitting the hits from the GridPix detector with the combined hits in the telescope. The hits in the telescope planes are merged in one super-point with a $10 \mu \mathrm{~m}$ error at the center of the last telescope plane. An example of GridPix hits with a fitted track is shown in Figure 3.

## 5. Testbeam results

### 5.1. Number of hits on track

In Figure 4 the number of GridPix hits associated to the track in the fiducial volume ( 216 pixels) is shown for a grid voltage of 350 V . The most probable number of hits is 91 and the mean is 114 for an effective track length of 12 mm . This is in agreement with the expected 100 electron-ion pairs/cm [11]. The shape of the distribution is Landau-like with a long tail due to delta electrons.

In Figure 5 the most probable number of hits is shown as a function of the grid voltage. One expects that the efficiency of the GridPix detector increases with grid voltage until it reaches a plateau at an efficiency of almost $100 \%$. Increasing further the grid voltage will induce crosstalk and far above 400 V sparks would be produced. The analyzed run was taken at a voltage of 350 V , close to the expected plateau and at a high efficiency. A search for double hits did not yield any indication for crosstalk.

### 5.2. Time walk correction

A pixel is hit when the collected charge is above the threshold. Because it takes longer to reach the threshold for a small signal than it does for a large signal, the measured Time of Arrival (ToA) depends on the magnitude of the signal. This is called the time walk. The capability to record simultaneously both ToA and Time over Threshold (ToT) per pixel is one of the improvements of

- Removed text the Timepix3 chip over its predecessor the Timepix chip. The Timepix3 chip allows to correct for the time walk by using the ToT as a measure of the signal magnitude.

First, the ToT was found to vary as a function of the column and therefore per column a correction factor for the ToT was introduced. The uncorrected mean of $z$-residuals - defined as the difference of the track fit prediction and the $z$-position of the hit - is shown as a function of the corrected ToT in Figure 7. The relation can be parametrised using the time walk $z_{\text {tw }}$ as a function of the corrected ToT $t_{\mathrm{ToT}}: \frac{\text { The following formula }}{\delta z_{\mathrm{tw}}=\frac{c_{1}}{t_{\mathrm{ToT}}+t_{0}}}$,
where $c_{1}$ and $t_{0}$ are constants determined from a fit. The distribution of $z$-residuals before and after applying the time walk correction, is shown in Figure 6. Functions with more parameters were also tried, but did not improve the results.

## New figure



Figure 7: Mean $z$-residual without time walk correction as function of ToT, fitted with equation 1

### 5.3. Hit resolution in the pixel plane

The momentum resolution of a TPC depends on the hit resolution in the pixel plane. There are two important contributions to the hit resolution in the pixel plane: a constant contribution caused by the pixel size $d_{\text {pixel }}$ and a transverse drift component that scales with drift distance and the diffusion coefficient $D_{T}$. The resolution $\sigma_{y}$ is given by:

$$
\begin{equation*}
\sigma_{y}^{2}=\frac{d_{\mathrm{pixel}}^{2}}{12}+D_{T}^{2}\left(z-z_{0}\right), \tag{2}
\end{equation*}
$$

where $z_{0}$ is the position of the grid. The measured hit resolution as a function of $z$-position is shown in Figure 8 The diffusion gives the largest contribution to the error in most of the detector volume. The measured diffusion coefficient $D_{T}=289 \mu \mathrm{~m} / \sqrt{\mathrm{cm}}$ is lower than the expected $D_{T}=310 \mu \mathrm{~m} / \sqrt{\mathrm{cm}}$. The $\chi^{2}$ is too high for the degrees of freedom, because no systematic uncertainties were taken into account.


Figure 8: Measured hit resolution in pixel plane (blue points) fitted with resolution according to equation $\sqrt{2}$ (red line).

### 5.5. Deformations

For a large TPC with pixel readout it is important that systematic deviations of position measurements are small and stay well below typically $20 \mu \mathrm{~m}$. Here we study deformations in the pixel plane and the drift direction, caused by for example distortions in the drift field or geometric uncertainties. The chip is divided in $64 \times 64$ bins of $4 \times 4$ pixels each for which the mean deformation is calculated. For every hit, the expected originating position on the track is traced and the residual is filled at that bin. In Figure 10 and Figure 11 the mean residual in the $x y$-plane and the mean $z$-residual are shown respectively. In the diagram only bins with more than 1000 entries are shown. Only bins in the selected fiducial area were used to make the distribution of the mean residuals shown in Figure 12. The r.m.s. of the deviations is $8 \mu \mathrm{~m}$ in the pixel plane and $31 \mu \mathrm{~m}(0.4 \mathrm{~ns})$ in the drift direction. This means that the systematics on the position measurement in the pixel plane - the bending plane of the TPC - are less than $10 \mu \mathrm{~m}$.

### 5.6. Energy loss measurement

A TPC can identify particles using their characteristic energy loss. The GridPix detector measures the energy loss $\mathrm{dE} / \mathrm{dx}$ by counting the number of detected electrons. Because of the large fluctuation in energy loss, the mean is dominated by a few high energy deposits. To retrieve a better estimate, the truncated sum of electrons is calculated.

Along the track, the number of electrons is counted for 20 pixel intervals. A fixed fraction of intervals with the highest number of electrons is rejected. Optimally, the top $10 \%$ of intervals with the highest number of electrons are rejected and from the other $90 \%$ of the intervals a truncated sum is obtained. In Figure 13 the truncated sum is shown for an effective track length of 1 m or 83 events. The: resolution, expressed as the standard deviation divided by the mean, is $4.1 \%$.

In order to estimate the separation power, the energy loss distribution for a Minimum Ionizing Particle (MIP): is estimated, see also Figure 13. The hit positions of the electron data are scaled track by track by a factor 0.7 to acquire the estimated ionization for a MIP, i.e. 0.7 m of electron data is scaled to 1 m of expected MIP data. The expected resolution for a 1 m MIP track is $4.9 \%$ and the separation between a MIP and a 2.5 GeV electron would be $6.2 \sigma$ for a 1 m track.

The truncated sum using slices of 20 pixels, does not: make full use of the fine granularity of the GridPix detector. We expect that particle identification can be improved by employing the full resolution to resolve primary ionization clusters 12 .
[12] M. Hauschild, 2D and 3D cluster counting with GEMs and small pads: The digital TPC?, in: Linear colliders. Proceedings, International Workshop on physics and experiments with future electron-positron linear colliders, LCWS 2002, Seogwipo, Jeju Island, Korea, August 26-30, 2002, pp. 464-469.

## 6. Conclusions and Outlook

A GridPix detector based on the Timepix3 chip was operated reliably in a testbeam setup for the first time. The resolution of the detector in the pixel planes and in the drift direction is limited primarily by diffusion. The additional systematic uncertainties are smaller than $10 \mu \mathrm{~m}$ in the pixel plane. Furthermore, by counting the ionization electrons, the energy loss $\mathrm{dE} / \mathrm{dx}$ can be measured with a precision of $4.1 \%$ for an effective track length of 1 m .

The next step towards a TPC for future applications, is the construction of larger size prototype detectors. R\&D has started to build a 4-chip module based on the Timepix3 that can be used to cover larger areas. With these developments, a pixelised readout is on its way to become a mature technology option for a large TPC at a future linear collider.

## Acknowledgements

This research was funded by the Netherlands Organisation for Scientific Research NWO. The authors want to thank the support of the mechanical and electronics departments at Nikhef and the accelerator group at the ELSA facility in Bonn. Their gratitude is extended to the Bonn SiLab group for providing the beam telescope.

